



Dark Photon and Stellar Cooling

Po-Yan Tseng (Yonsei U.)

References:

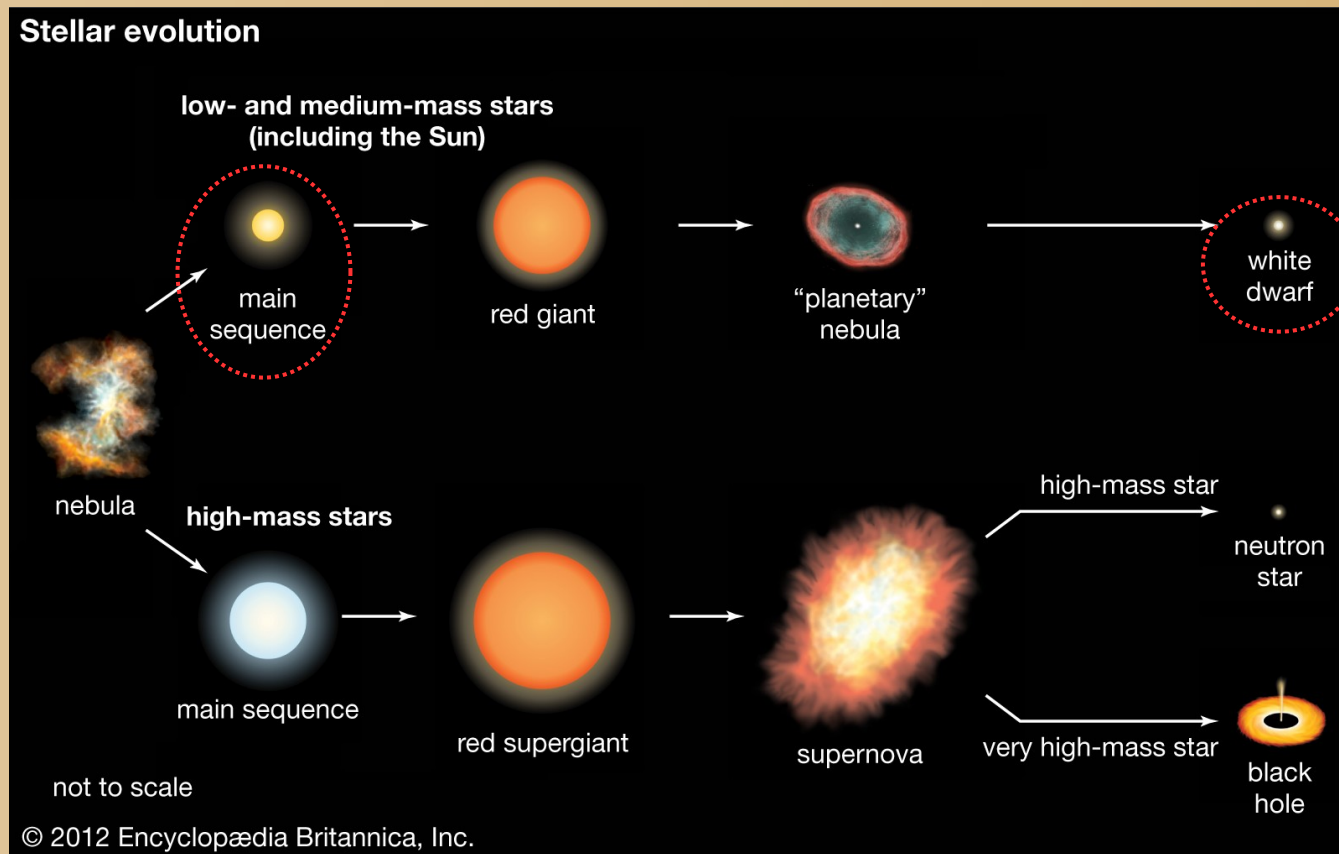
Wai-Yee Keung, Danny Marfatia: [arXiv-2009.04444](https://arxiv.org/abs/2009.04444)

**The 17th Saga-Yonsei partnership program of HEP,
7-8, 21-22, Jan. 2021**

Life of Stellar

Stellar evolution:

Stellar evolution: Encycloaedia Britannica, Inc.



Outline

- ◆ Stellar cooling & White Dwarf luminosity function.
- ◆ Dark Photon production from stellar
- ◆ XENON1T
- ◆ Inelastic DM
- ◆ Summary

Stellar Cooling

- Neutrino emission dominates the stellar cooling, if $T \gtrsim 5 \times 10^7 \text{ K} \sim \text{keV}$.

$$\epsilon_\nu = 1.40 \times 10^{15} \lambda^9 \gamma^6 e^{-\gamma} (f_T + f_L) \text{ erg g}^{-1} \text{ s}^{-1}$$

$$\lambda = 1.69 \times 10^{-3} T_7, \quad \gamma = \frac{28}{T_7},$$
$$f_T = 2.4 + 0.6\gamma^{1/2} + 0.51\gamma + 1.25\gamma^{3/2}, \quad f_L = \frac{8.6\gamma^2 + 1.35\gamma^{7/2}}{225 - 17\gamma + \gamma^2}$$

- Photon emission dominates, when $T \lesssim 10^7 \text{ K}$.

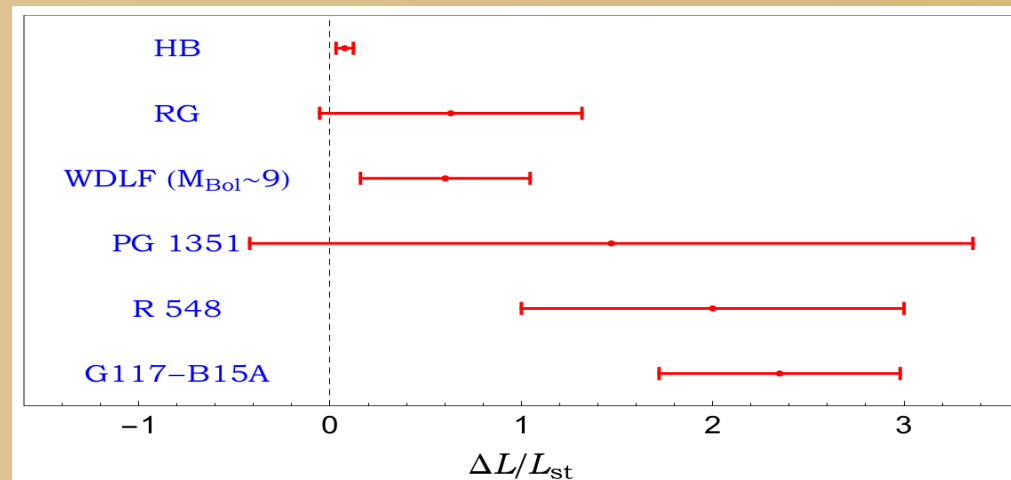
$$\epsilon_\gamma = 3.29 \times 10^{-3} T_7^{7/2} \text{ erg g}^{-1} \text{ s}^{-1}$$

L.Ubaldi: 1310.5073

Stellar Cooling

- Neutrino emission dominates the stellar cooling, if $T \gtrsim 5 \times 10^7 \text{ K} \sim \text{keV}$.
- Photon emission dominates, when $T \lesssim 10^7 \text{ K}$.
- Deviation from SM expectation where found in diverse stellar systems:

M.Giannotti, I.Irastorza, J.Redondo, A.Ringwald: 1512.08108



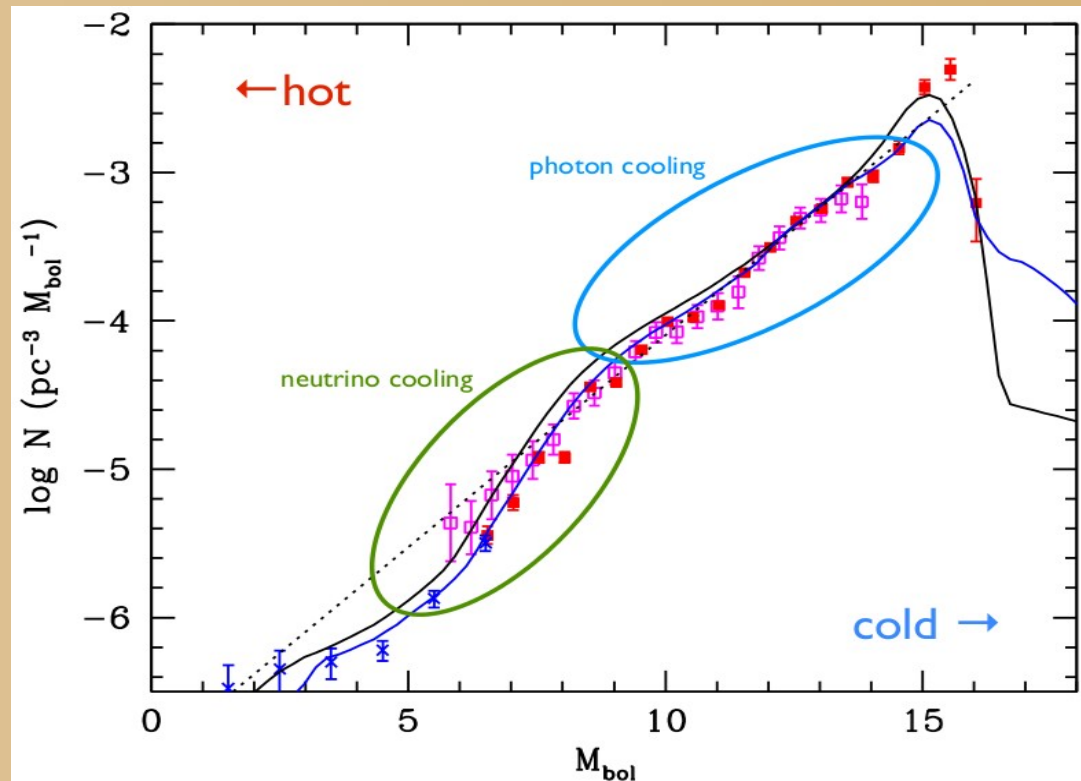
WDLF

- White Dwarf Luminosity Function(WDLF): number density distribution of WD per luminosity or brightness(bolometric magnitude M_{bol}) bin.

$$\frac{dN}{dM_{\text{bol}}} = -B \frac{dU}{dM_{\text{bol}}} \frac{1}{L_{\gamma} + L_{\nu} + L_x}$$

$$M_{\text{bol}} = -2.5 \log_{10}(L/L_{\odot}) + 4.74$$

L.Ubaldi: 1310.5073



Stellar Cooling

- Massive **Dark photon(DP)** with mass lighter than few keV.
- **DP** has transverse and longitudinal modes, and photon can oscillate into DP w/o an external magnetic or electric field. Turn out more efficient cooling.

M.Giannotti, I.Irastorza, J.Redondo, A.Ringwald: 1512.08108

- The **DP** *thermal production rate* from *thermal plasma* is related to imaginary part of *self-energy* in the medium.

$$\Gamma_{\text{prod}} = -\frac{\text{Im } \Pi}{\omega (e^{\omega/T} - 1)}$$

J.Redondo, G.Raffelt: 1305.2920,
H.An, M.Pospelov, J.Pradler: 1302.3884.

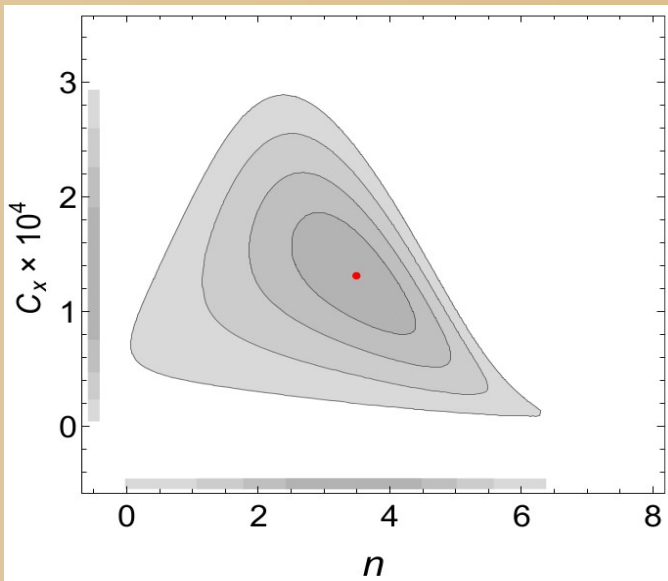
Stellar Cooling

- ▶ Massive **Dark photon(DP)** with mass lighter than few keV.
- ▶ For colder WD $7.75 \lesssim M_{\text{bol}} \lesssim 14.25$, where neutrino cooling is negligible Used the *model-independent* exotic cooling rate $L_x = C_x L_{\odot} T_7^n$. They found **4-sigma** significance of additional cooling.

M.Giannotti, I.Irastorza, J.Redondo, A.Ringwald: 1512.08108

Stellar Cooling

- Massive **Dark photon(DP)** with mass lighter than few keV.



$$L_X \simeq 3.13 \times 10^{32} \text{ GeV/s}$$

$$L_{\odot} = 2.39 \times 10^{36} \text{ GeV/s}$$

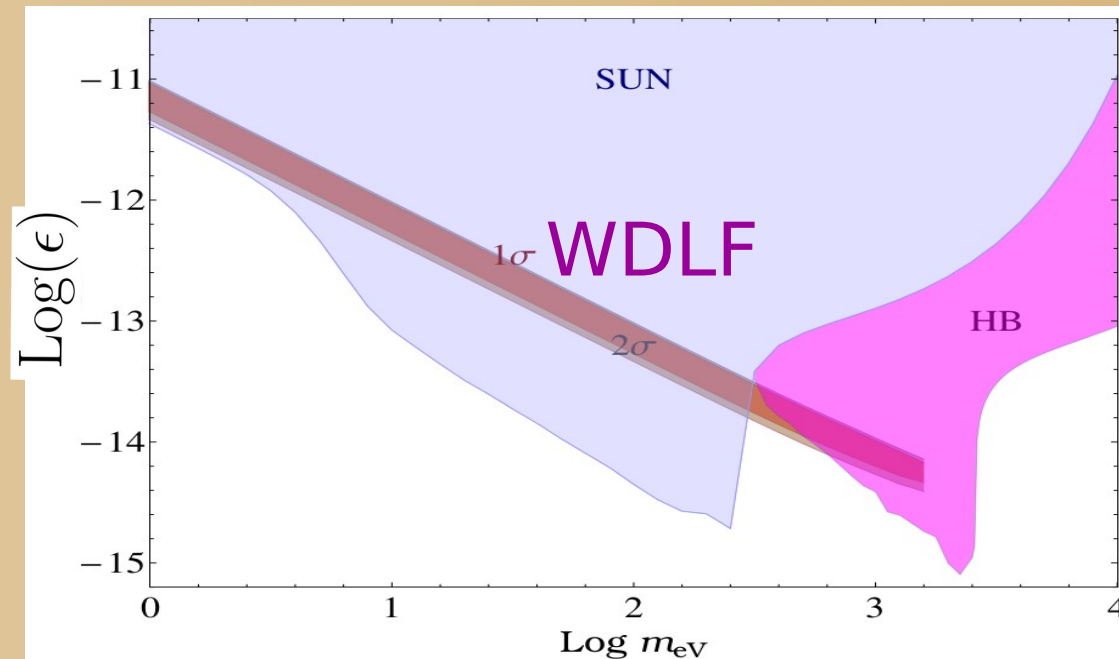
Figure 2. Contours of the 1, 2, 3 and 4 σ C.L. in the n and C_x parameter space, defined in Eq. (3.2). The red dot refers to the best fit value: $n_{\text{best}} = 3.49$, $C_{x,\text{best}} = 1.31 \times 10^{-4}$. The analysis is based on the data in [9] for $7.75 \lesssim M_{\text{bol}} \lesssim 14.25$.

M.Giannotti, I.Irastorza, J.Redondo, A.Ringwald: 1512.08108

Stellar Cooling

- Massive **Dark photon(DP)** with mass lighter than few keV.

$$\mathcal{L}_{A'} \supset (\epsilon e)(\bar{\ell}\gamma^\mu\ell)A'_\mu + \frac{\epsilon}{2}F_{\mu\nu}F'^{\mu\nu}$$



M.Giannotti, I.Irastorza, J.Redondo, A.Ringwald: 1512.08108

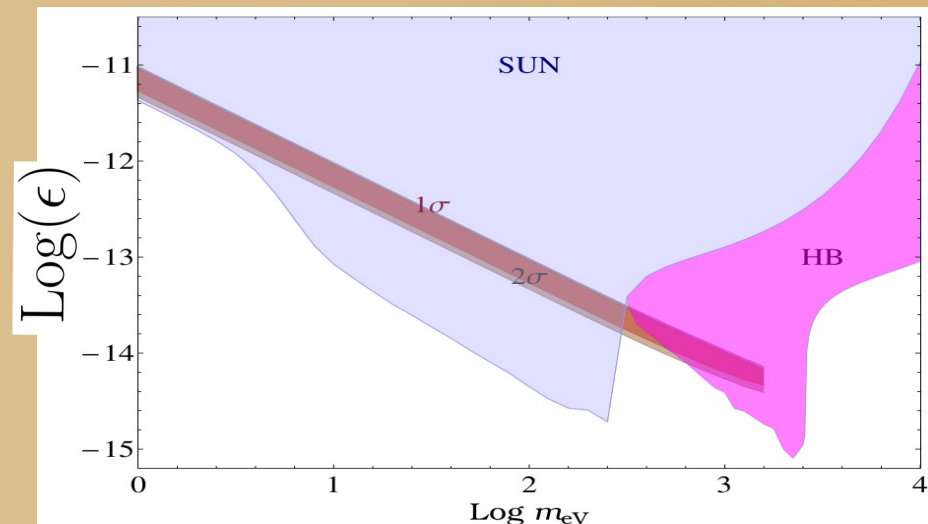
Dark Photon production from Stellar

- **DP** can be produced from **photon conversion**, **bremsstrahlung** and **Compton scattering** from **thermal plasma** inside stellar.
- The plasma frequency $\omega_p = (n_e e^2 / m_e)^{1/2}$

Sun: $\omega_p \simeq 0.3$ keV

HB: $\omega_p \simeq 2$ keV

WD: $\omega_p \simeq 30$ keV



Dark Photon production from Stellar

- Emission rate of **DP** longitudinal mode from stellar:

$$\frac{d^2\Gamma_{A'}^\odot}{dV d\omega} \Big|_L = \frac{1}{\pi^2} \frac{\omega \sqrt{\omega^2 - m_{A'}^2}}{e^{\frac{\omega}{T_\odot}} - 1} \frac{\epsilon^2 m_{A'}^2 \omega^2 \Gamma_L}{(\omega^2 - \omega_p^2)^2 + (\omega \Gamma_L)^2},$$

$$\Gamma_L = \frac{8\pi\alpha^3 n_e \sum_{i=H,He} Z_i^2 n_{Z_i}}{3\sqrt{2\pi T_\odot} m_e^{3/2} \omega^3} F\left(\frac{\omega}{T_\odot}\right) + \frac{8\pi\alpha^2 n_e}{3m_e^2} \sqrt{1 - \frac{\omega_P^2}{\omega^2}},$$

$$F(\omega) = (1 - e^{-\omega}) \int_0^\infty dx x e^{-x^2} \int_{\sqrt{x^2 + \omega - x}}^{\sqrt{x^2 + \omega + x}} \frac{t^3 dt}{(t^2 + y^2)^2}, \quad y = k_s / \sqrt{2m_e T}, \quad k_s \text{ a screening scale}$$

H.An, M.Pospelov, J.Pradler,1302.3884. J.Redondo,G.Raffelt:1305.2920

- Longitudinal mode dominates if $m_{A'} \ll \omega_p$.

Dark Photon production from Stellar

- Emission rate of **DP** transverse mode from stellar:

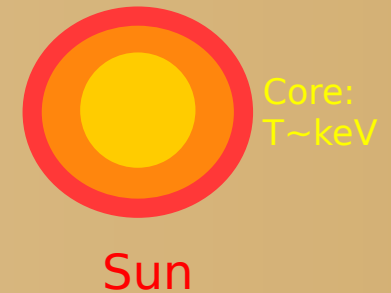
$$\frac{d^2\Gamma_{A'}^\odot}{dV d\omega}\Big|_T = \frac{1}{\pi^2} \frac{\omega \sqrt{\omega^2 - m_{A'}^2}}{e^{\frac{\omega}{T_\odot}} - 1} \frac{\epsilon^2 m_{A'}^4}{(\omega_p^2 - m_{A'}^2)^2 + (\omega \Gamma_T)^2} \Gamma_T,$$
$$\Gamma_T = \frac{16\pi^2 \alpha^3}{3m_e^2 \omega^3} \sqrt{\frac{2\pi m_e}{3T_\odot}} n_e \sum_{i=H, He} Z_i^2 n_{Z_i} \bar{g}_i (1 - e^{-\frac{\omega}{T_\odot}}) + \frac{8\pi \alpha^2}{3m_e^2} n_e$$

J.Redondo, JCAP 0807,008(2008): 0801.1527

- Transverse mode dominates if $m_{A'} \simeq \omega_p$ and $m_{A'} \gtrsim \omega_p$.
- The \bar{g}_i is a Boltzmann averaged Gaunt factor.

Dark Photon production from Stellar

- **KeV DP** flux produced from thermal plasma of *core* our **Sun** could be one interpretation.
- We did **NOT** assume **DP** is DM relic.
- The **DP** flux from the Sun,
- For $T_{\odot} = 1.15 \text{ keV}$ and $R_{\text{core}}/R_{\odot} \simeq 0.18$:



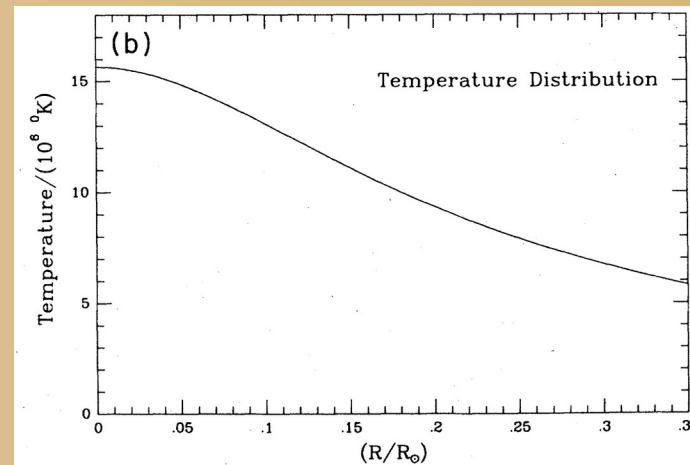
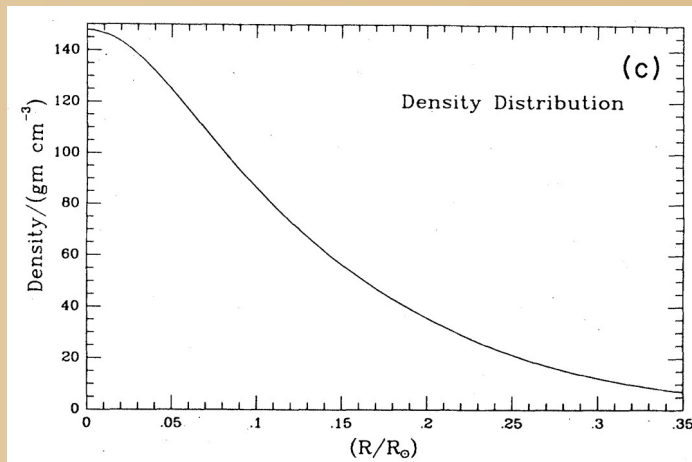
$$\frac{d\Phi_{A'}^{\odot}}{d\omega} = \frac{1}{4\pi D_{\text{SE}}^2} \int dV \frac{d^2\Gamma_{A'}^{\odot}}{dV d\omega}$$

$$\frac{dR}{d\omega} \simeq N_T \sigma_{A'} \frac{d\Phi_{A'}^{\odot}}{d\omega}$$

Dark Photon production from Stellar

- ◆ **KeV DP** flux produced from thermal plasma of *core* our **Sun** could be one interpretation.
- ◆ We did **NOT** assume **DP** is DM relic.
- ◆ The **DP** flux from the Sun,
- ◆ For $T_{\odot} = 1.15 \text{ keV}$ and $R_{\text{core}}/R_{\odot} \simeq 0.18$:

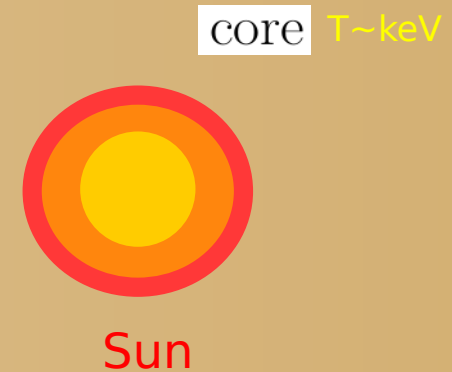
core $T \sim \text{keV}$



John Bahcall, Roger Ulrich, Rev.Mod.Phy 60.297

Dark Photon production from Stellar

- **KeV DP** flux produced from thermal plasma of *core* our **Sun** could be one interpretation.
- We did **NOT** assume **DP** is DM relic.
- The **DP** flux from the Sun,
- For $T_{\odot} = 1.15 \text{ keV}$ and $R_{\text{core}}/R_{\odot} \simeq 0.18$:



$$\frac{d\Phi_{A'}^{\odot}}{d\omega} = \frac{1}{4\pi D_{\text{SE}}^2} \int dV \frac{d^2\Gamma_{A'}^{\odot}}{dV d\omega}$$

$$\frac{dR}{d\omega} \simeq N_T \sigma_{A'} \frac{d\Phi_{A'}^{\odot}}{d\omega}$$

Events at XENON1T

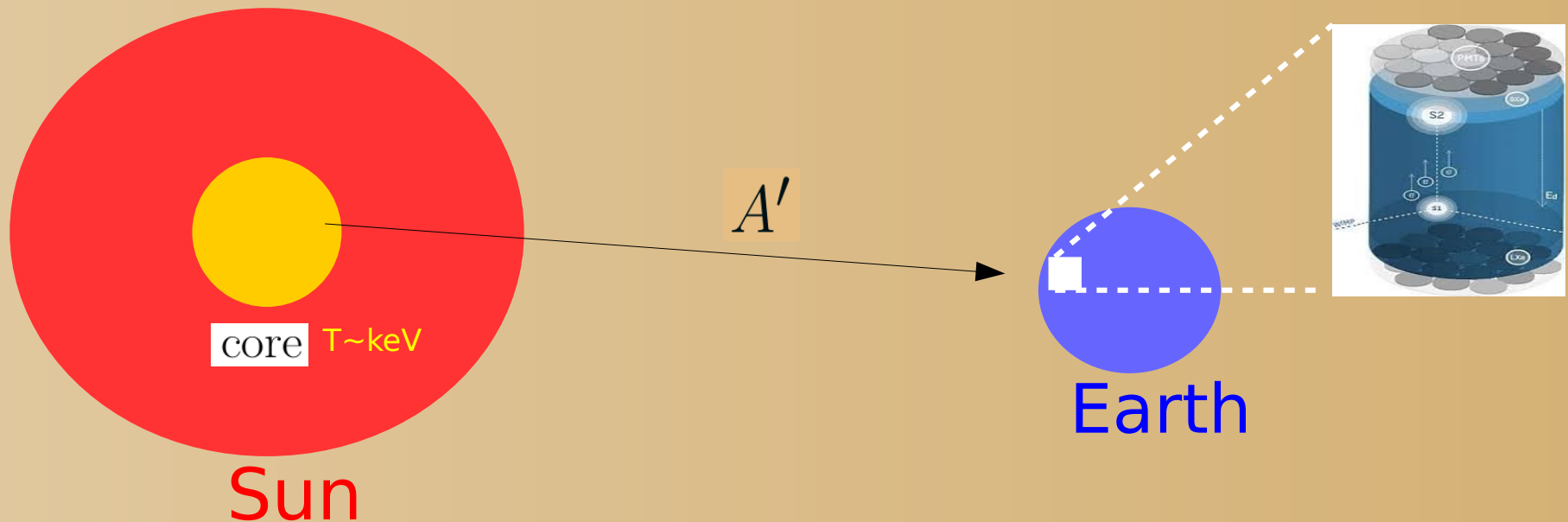
- **KeV DP** flux produced from thermal plasma of *core* our **Sun** could be one interpretation.
- We did **NOT** assume **DP** is DM relic.
- Then **DP** contributes to **XENON1T** event through absorption by bounded electron: $A'e \rightarrow e$

$$\sigma_{A'} \begin{cases} \simeq \left(\frac{m_{A'}}{\omega}\right)^2 \epsilon^2 \sigma_\gamma \left(\frac{c}{v_{A'}}\right) & \text{for longitudinal } A', \\ = \epsilon^2 \sigma_\gamma \left(\frac{c}{v_{A'}}\right) & \text{for transverse } A' \text{ [25]}, \end{cases}$$

Events at XENON1T

- KeV DP flux produced from thermal plasma of *core* our Sun could be one interpretation.

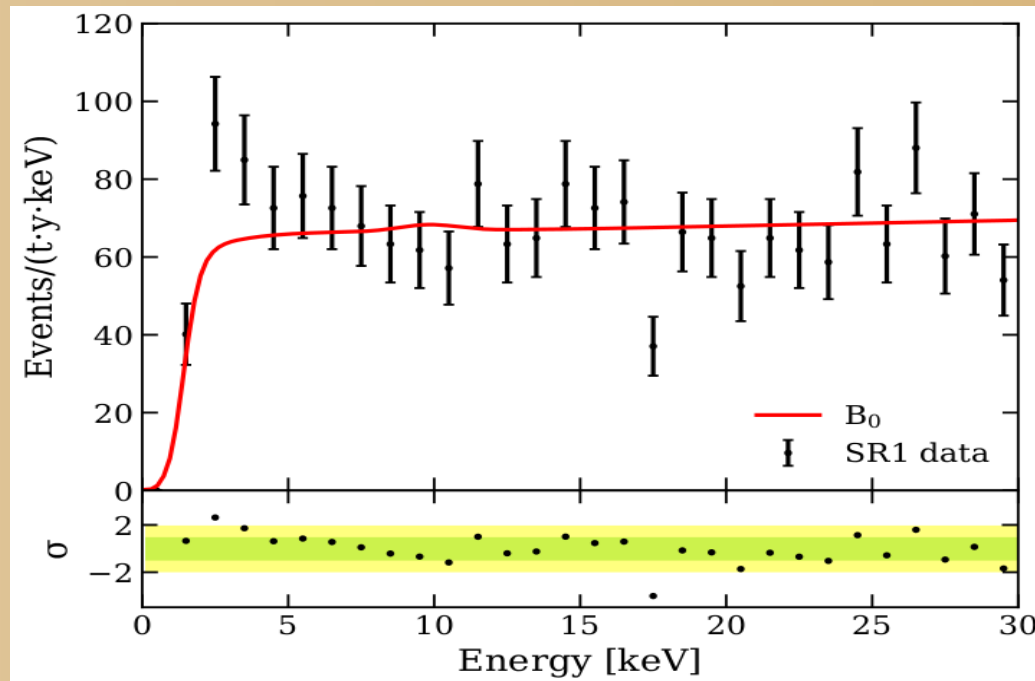
$$\mathcal{L}_{A'} \supset (\epsilon e)(\bar{\ell}\gamma^\mu\ell)A'_\mu + \frac{\epsilon}{2}F_{\mu\nu}F'^{\mu\nu}$$



W.Y.Keung,D.Marfatia, P.Y.Tseng, arXiv:2009.04444

Events at XENON1T

- XENON1T experiment reported **3-sigma** excess in the electron recoil spectrum between 2-3 keV.



XENON Collaboration, PRD 102.072004, arXiv:2006.09721

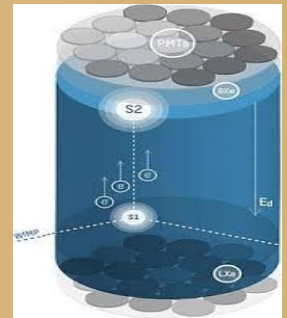
Events at XENON1T

- ◆ **KeV DP** flux produced from thermal plasma of *core* our **Sun** could be one interpretation.
- ◆ The event rate at XENON1T detector :

$$\frac{dR}{d\omega} \simeq N_T \sigma_{A'} \frac{d\Phi_{A'}^{\odot}}{d\omega}$$

- ◆ The $N_T \simeq 4.52 \times 10^{27}$ is the number of Xe atoms per ton.

XENON1T



Events at XENON1T

- ◆ **KeV DP** flux produced from thermal plasma of *core* our **Sun** could be one interpretation.
- ◆ The event rate at XENON1T detector :

$$\frac{dR}{d\omega} \simeq N_T \sigma_{A'} \frac{d\Phi_{A'}^{\odot}}{d\omega}$$

- ◆ The absorption cross section $A'e \rightarrow e$:

$$\sigma_{A'} \begin{cases} \simeq \left(\frac{m_{A'}}{\omega}\right)^2 \epsilon^2 \sigma_{\gamma} \left(\frac{c}{v_{A'}}\right) & \text{for longitudinal } A', \\ = \epsilon^2 \sigma_{\gamma} \left(\frac{c}{v_{A'}}\right) & \text{for transverse } A' \text{ [25]}, \end{cases}$$

XENON1T

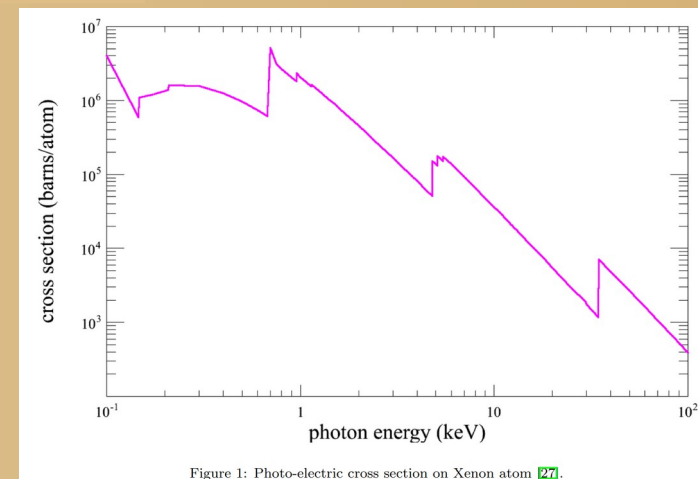
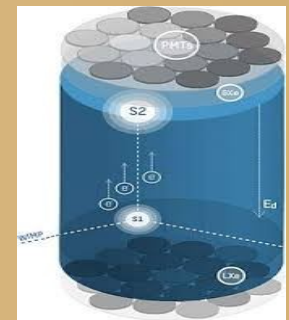


Figure 1: Photo-electric cross section on Xenon atom [27].

K. Arosala et al., 1209.3810

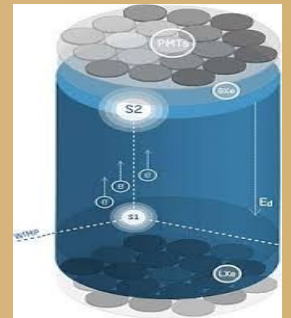
Events at XENON1T

- ◆ **KeV DP** flux produced from thermal plasma of *core* our **Sun** could be one interpretation.
- ◆ The event rate at XENON1T detector :

$$\frac{dR}{d\omega} \simeq N_T \sigma_{A'} \frac{d\Phi_{A'}^{\odot}}{d\omega}$$

$$\Rightarrow \frac{dT}{dE} = N_T \int \left(\sigma_{A'} \frac{d\Phi_{A'}^{\odot}}{d\omega} \right) Res(E, \omega) d\omega$$

XENON1T



- ◆ The energy resolution by Gaussian smearing:

$$\frac{\sigma_{\text{det}}}{E_R} = \frac{a}{\sqrt{E_R/\text{keV}}} + b$$

$$a = 0.3171 \pm 0.0065 \text{ and } b = 0.0015 \pm 0.0002$$

$$Res(E, E_R) = \frac{1}{\sqrt{2\pi\sigma_{\text{det}}^2}} e^{-\frac{(E-E_R)^2}{\sigma_{\text{det}}^2}} \alpha(E)$$

J.Bramante, N.Song, 2006.14089

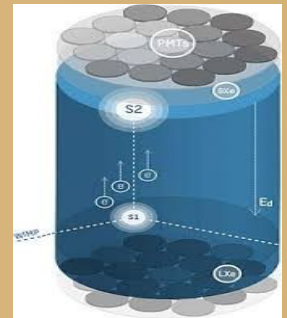
Events at XENON1T

- ◆ **KeV DP** flux produced from thermal plasma of *core* our **Sun** could be one interpretation.
- ◆ The event rate at XENON1T detector :

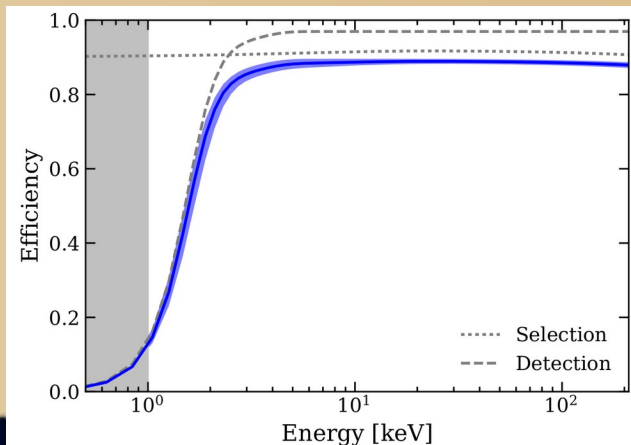
$$\frac{dR}{d\omega} \simeq N_T \sigma_{A'} \frac{d\Phi_{A'}^{\odot}}{d\omega}$$

$$\Rightarrow \frac{dT}{dE} = N_T \int \left(\sigma_{A'} \frac{d\Phi_{A'}^{\odot}}{d\omega} \right) Res(E, \omega) d\omega$$

XENON1T



- ◆ The energy resolution by Gaussian smearing:

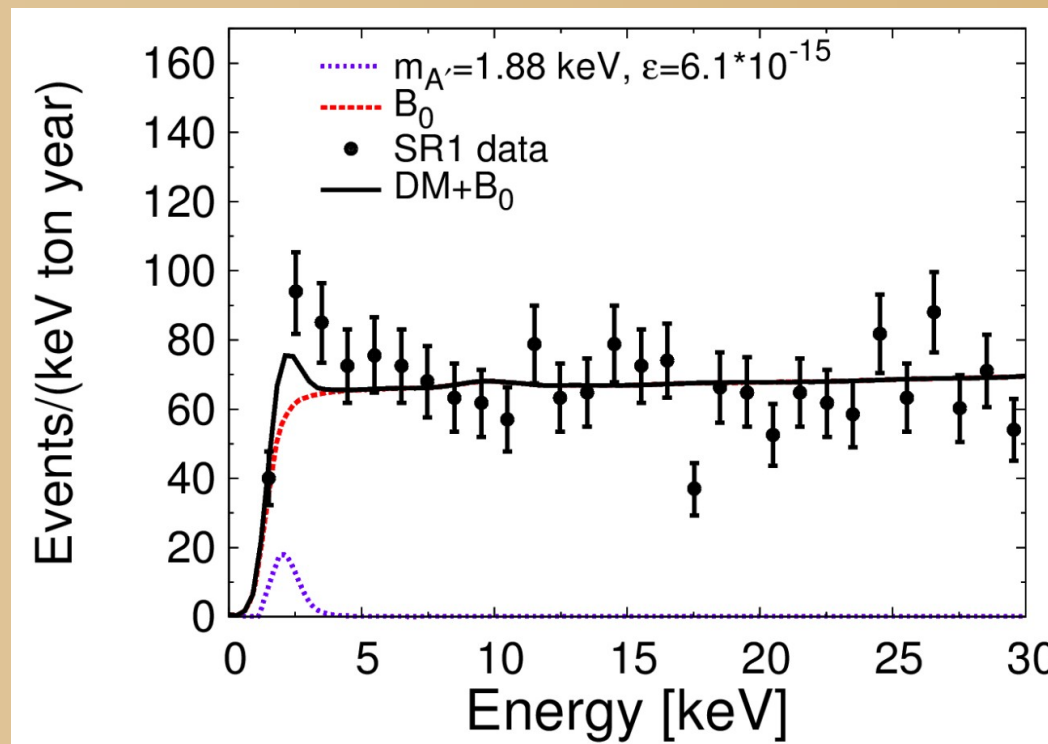


$$Res(E, E_R) = \frac{1}{\sqrt{2\pi\sigma_{det}^2}} e^{-\frac{(E-E_R)^2}{\sigma_{det}^2}} \alpha(E)$$

XENON1T collaboration, 2006.09721

Events at XENON1T

- **KeV DP** flux produced from thermal plasma of *core* our **Sun** could be one interpretation.
- The **DP** flux from the Sun and **XENON1T** event rate:



WDLF Anomaly

- **KeV DP** contribute additionally emission for WD cooling.
- Using the approximation in the limit $m_{A'} \ll \omega_p$

$$\frac{d^2\Gamma_{A'}^{\text{WD}}}{dV d\omega} = \frac{1}{4\pi} \frac{\epsilon^2 m_{A'}^2 \omega^2}{e^{\omega/T} - 1} \delta(\omega - \omega_p)$$

- Additional luminosity from **KeV DP**:

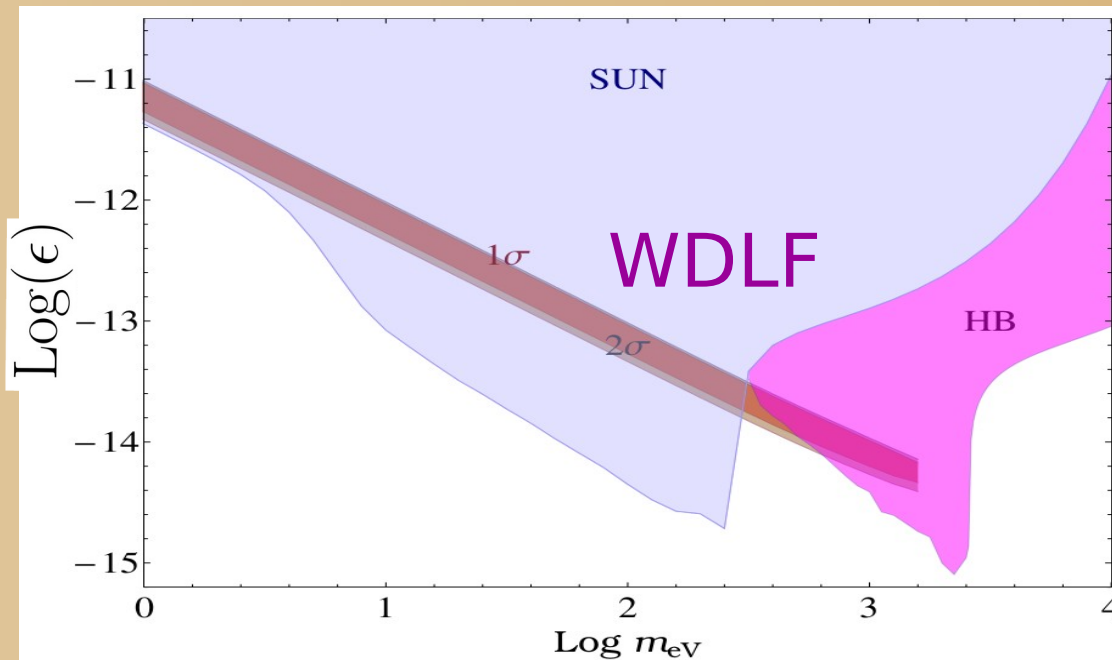
$$L_X = \int dV \int d\omega \omega \frac{d^2\Gamma_{A'}^{\text{WD}}}{dV d\omega} = \left(\frac{4}{3} \pi R_{\text{WD}}^3 \right) \frac{1}{4\pi} \frac{\epsilon^2 m_{A'}^2 \omega_p^3}{e^{\omega_p/T_{\text{WD}}} - 1},$$

i.e., $\epsilon m_{A'} \simeq 10^{-14} \text{ keV}$.

$$L_X \simeq 3.13 \times 10^{32} \text{ GeV/s}$$

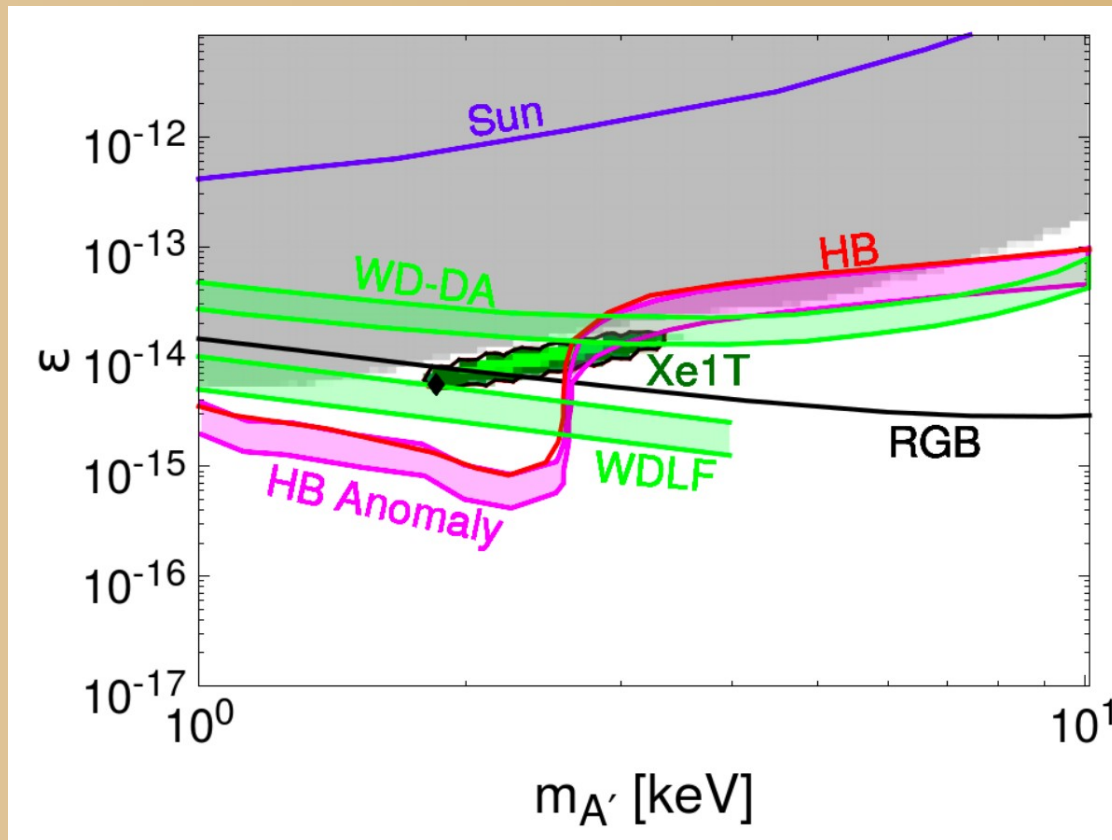
WDLF Anomaly

- **KeV DP** contribute additionally emission for WD cooling.



XENON1T and WDLF

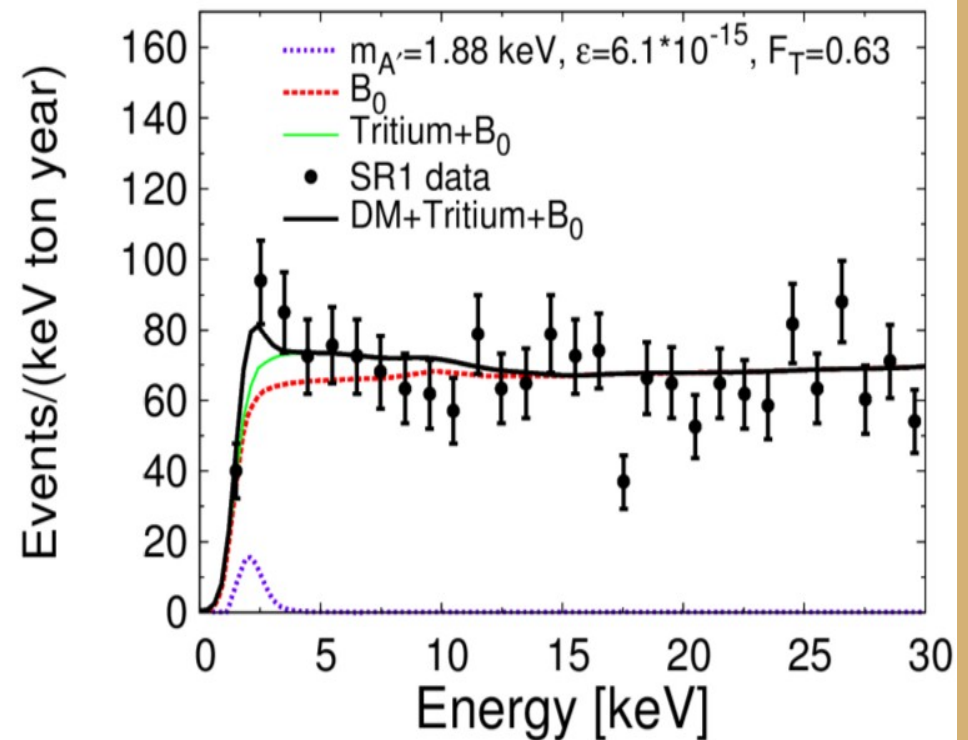
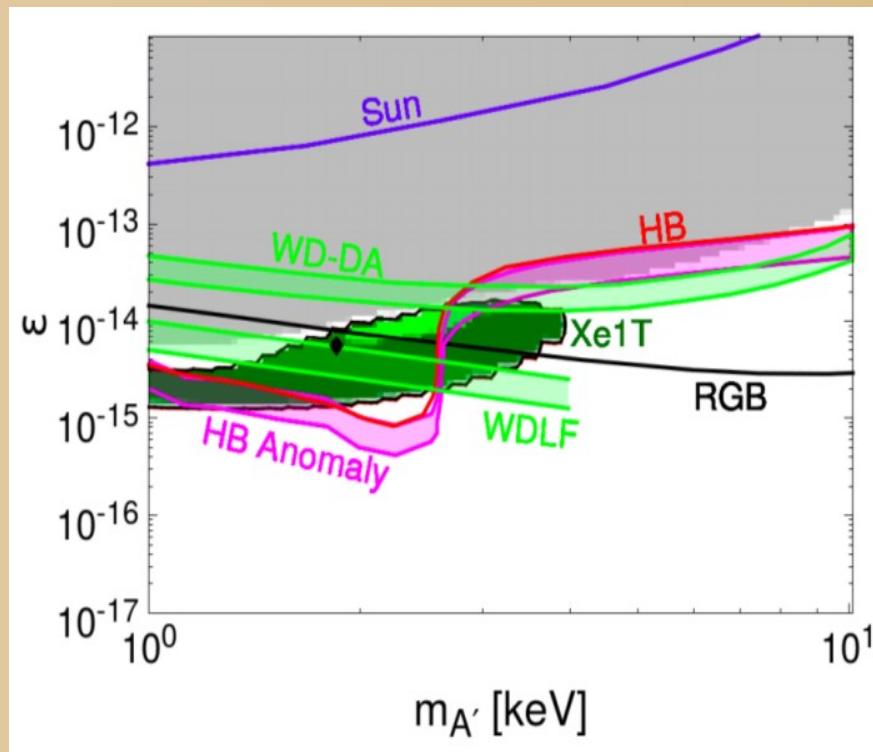
- Comparing to stellar cooling constraints:



W.Y.Keung, D.Marfatia, P.Y.Tseng, arXiv:2009.04444

XENON1T and WDLF

- Including **Tritium** contribution to **XENON1T**:



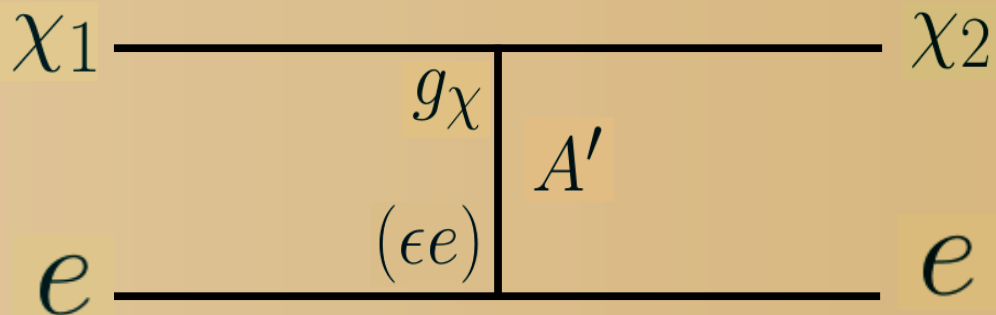
W.Y.Keung, D.Marfatia, P.Y.Tseng, arXiv:2009.04444

Inelastic DM

- If **DP** is much heavier than keV, it cannot be thermally produced from the Sun.
- We consider **DP** as a mediator between *MeV Inelastic DM* and electron.

$$\mathcal{L} \supset (\epsilon e) A'_\mu (\bar{e} \gamma^\mu e) + \left(\frac{i g_\chi}{2} A'_\mu (\bar{\chi}_2 \gamma^\mu \chi_1) + \text{h.c.} \right)$$

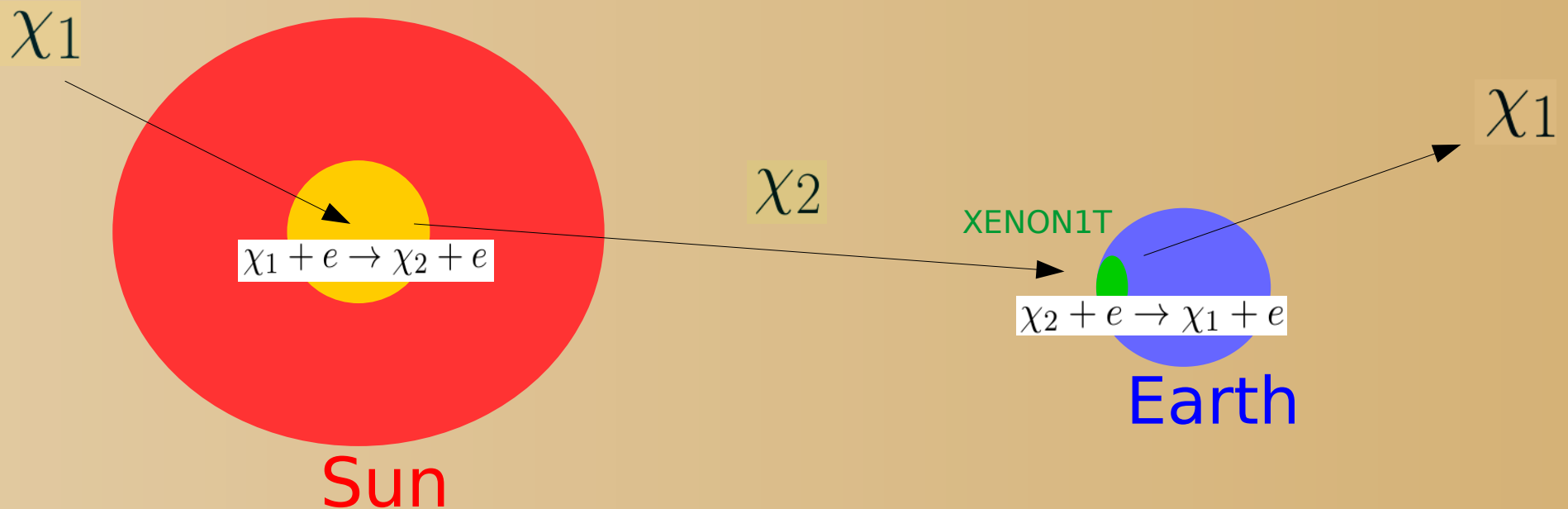
$$\Delta m_\chi \equiv m_{\chi_2} - m_{\chi_1} \approx \text{keV}$$



W.Y.Keung, D.Marfatia, P.Y.Tseng, arXiv:2009.04444

Inelastic DM

- If DM relic is mainly composed by χ_1 . When DM wind cross the Sun, χ_1 is excited into χ_2 . Then χ_2 propagate to Earth, down scattering at XENON1T.



W.Y.Keung,D.Marfatia, P.Y.Tseng, arXiv:2009.04444

Inelastic DM

- If DM relic is mainly composed by χ_1 . When DM wind cross the Sun, χ_1 is excited into χ_2 . Then χ_2 propagate to Earth.

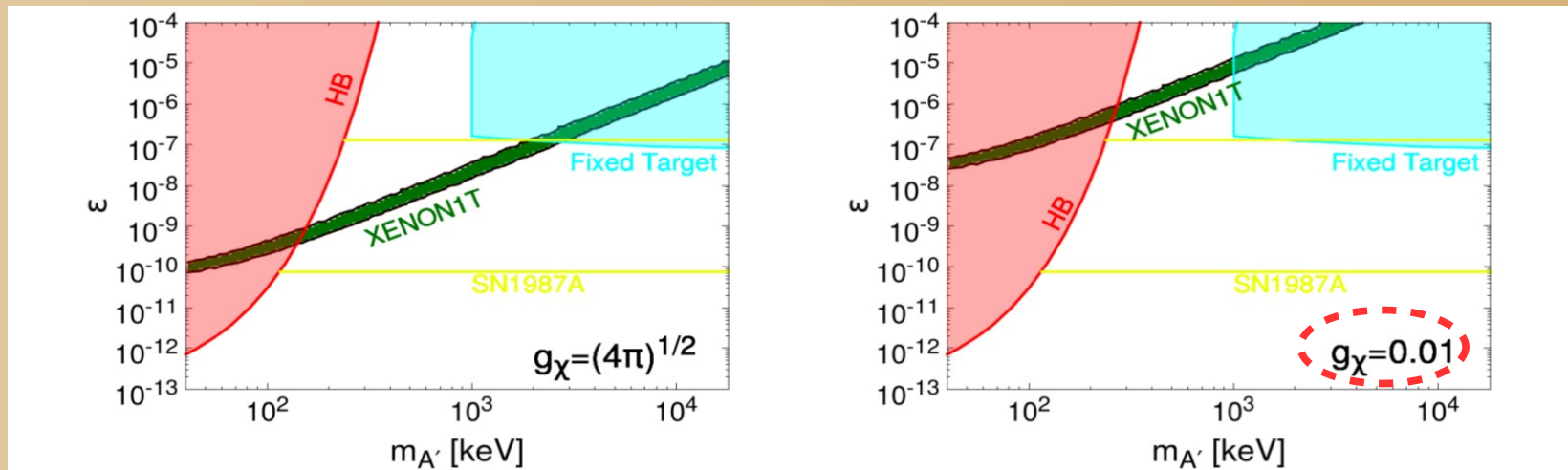


Figure 2. $m_{A'} > \Delta m_{\chi}$. The 2σ XENON1T allowed regions for $g_{\chi} = \sqrt{4\pi}$ and $g_{\chi} = 0.01$. The constraints from HB stars [13, 14], fixed-target experiments [23] and an approximate constraint from SN 1987A [24] are also shown.

W.Y.Keung,D.Marfatia, P.Y.Tseng, arXiv:2009.04444

Inelastic DM

- If DM relic is mainly composed by χ_1 . When DM wind cross the Sun, χ_1 is excited into χ_2 . Then χ_2 propagate to Earth.

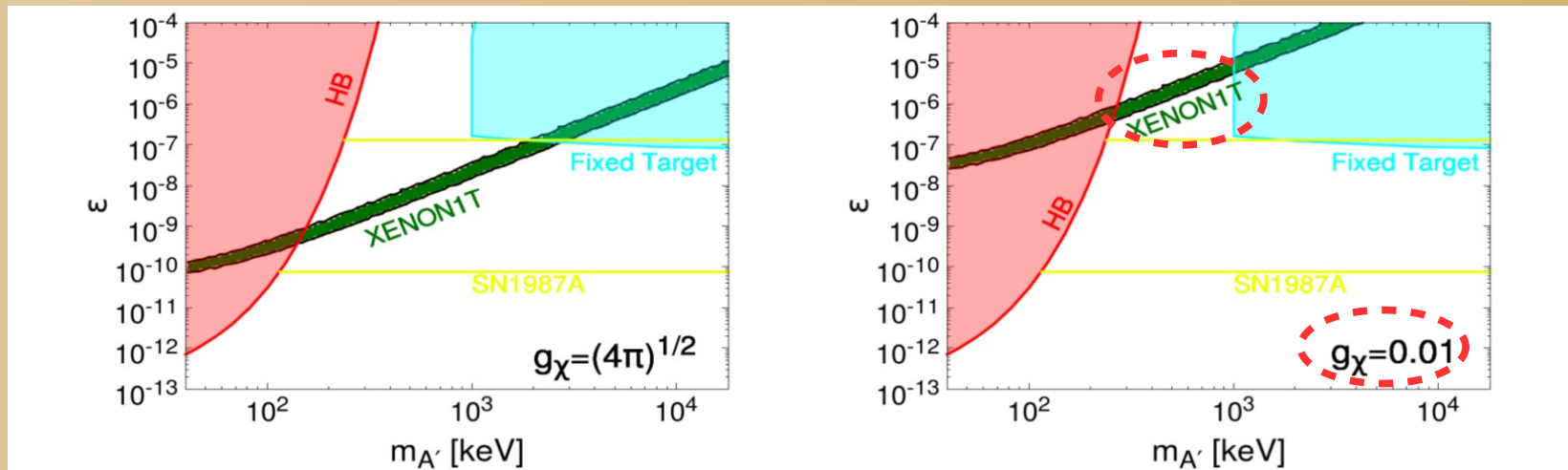
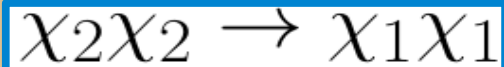


Figure 2. $m_{A'} > \Delta m_\chi$. The 2σ XENON1T allowed regions for $g_\chi = \sqrt{4\pi}$ and $g_\chi = 0.01$. The constraints from HB stars [13, 14], fixed-target experiments [23] and an approximate constraint from SN 1987A [24] are also shown.

W.Y.Keung,D.Marfatia, P.Y.Tseng, arXiv:2009.04444

Inelastic DM

- The χ_2 population after freeze out. The inter-conversion process:



J.Barmante, N.Song: , arXiv:2006.14089

- The ratio of number density is given by:

$$\frac{n_{\chi_2}}{n_{\chi_1}} \sim e^{-\frac{\Delta m_{\chi}}{T_{co}}} \Rightarrow n_{\chi_2} \ll n_{\chi_1}$$

$$T_{co} = \left(\frac{\sqrt{\pi g_*} m_{A'}^4}{\alpha_{\chi}^2 M_{pl}} \right)^{1/3}$$

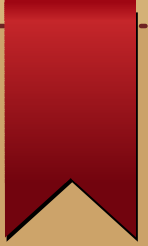
- The χ_1 dominates DM relic.

Summary

- ◆ **Stellar cooling** is sensitive to *microscopy* physics and *particles lighter than keV*.
- ◆ We performed the computation of **keV dark photon** *thermal production from thermal plasma* of stellar systems, and event rate at **XENON1T** detector.
- ◆ The **keV dark photon** could be simultaneous explanation of the **WDLF** cooling anomaly and **XENON1T** excess by including the **Tritium** contribution.

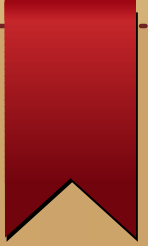
Summary

- For *inelastic DM* with heavy mediator: **thermal electrons** inside the **Sun** provide a power source to excite $\chi_1 + e \rightarrow \chi_2 + e$. After propagate to Earth, it contribute to **XENON1T** signal via down scattering $\chi_2 + e \rightarrow \chi_1 + e$.



Thank You!





Back Up



Stellar Cooling

- Observation of *period decreasing rate* of type DA pulsating WD

WD	class	$\dot{P}_{\text{obs}}[\text{s/s}]$	$\dot{P}_{\text{th}}[\text{s/s}]$
G117 - B15A	DA	$(4.19 \pm 0.73) \times 10^{-15}$	$(1.25 \pm 0.09) \times 10^{-15}$
R548	DA	$(3.33 \pm 1.1) \times 10^{-15}$	$(1.1 \pm 0.09) \times 10^{-15}$
PG 1351+489	DB	$(2.0 \pm 0.9) \times 10^{-13}$	$(0.81 \pm 0.5) \times 10^{-13}$

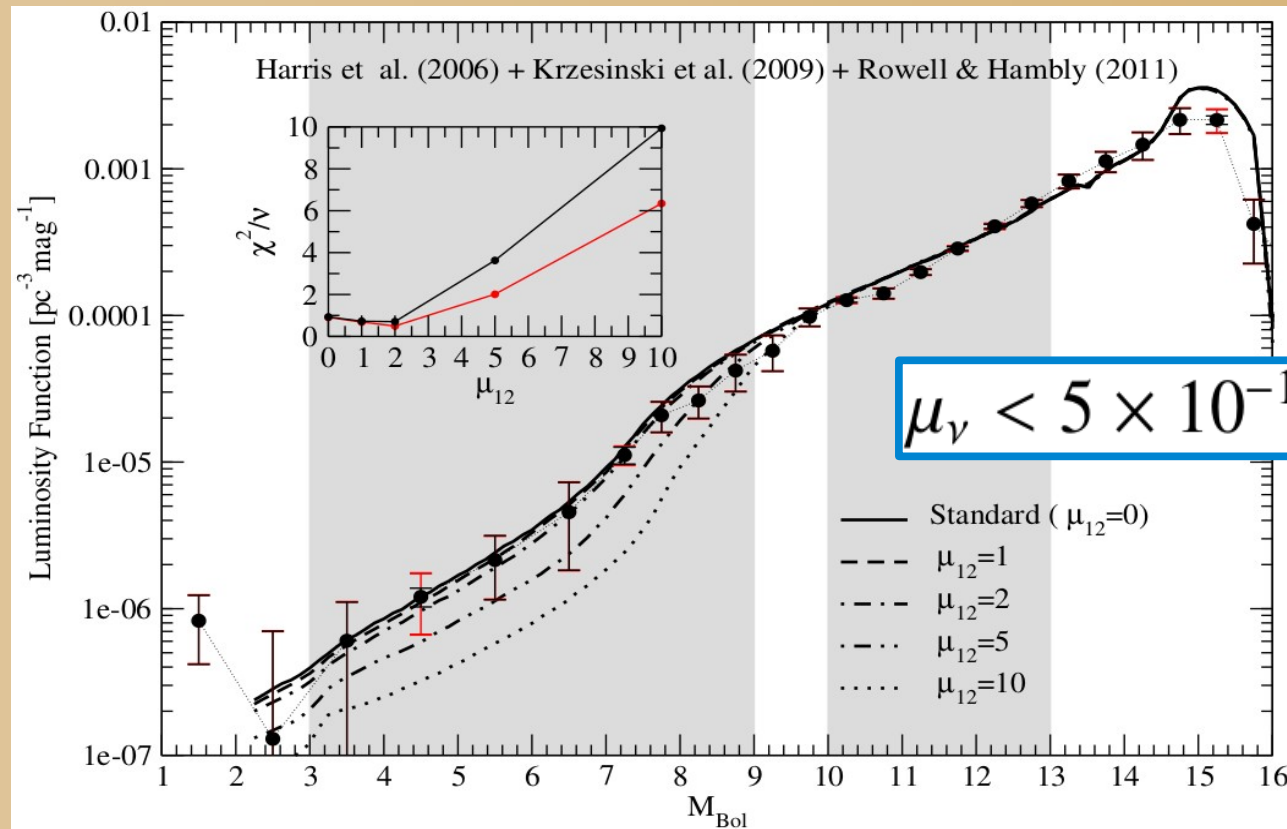
Table 1. Results for \dot{P} for G117 - B15A [4], R548 [6], and PG 1351+489 [7].

M.Giannotti, I.Irastorza, J.Redondo, A.Ringwald: 1512.08108

- *Period decreasing rate* is proportional to *cooling rate*.

WDLF

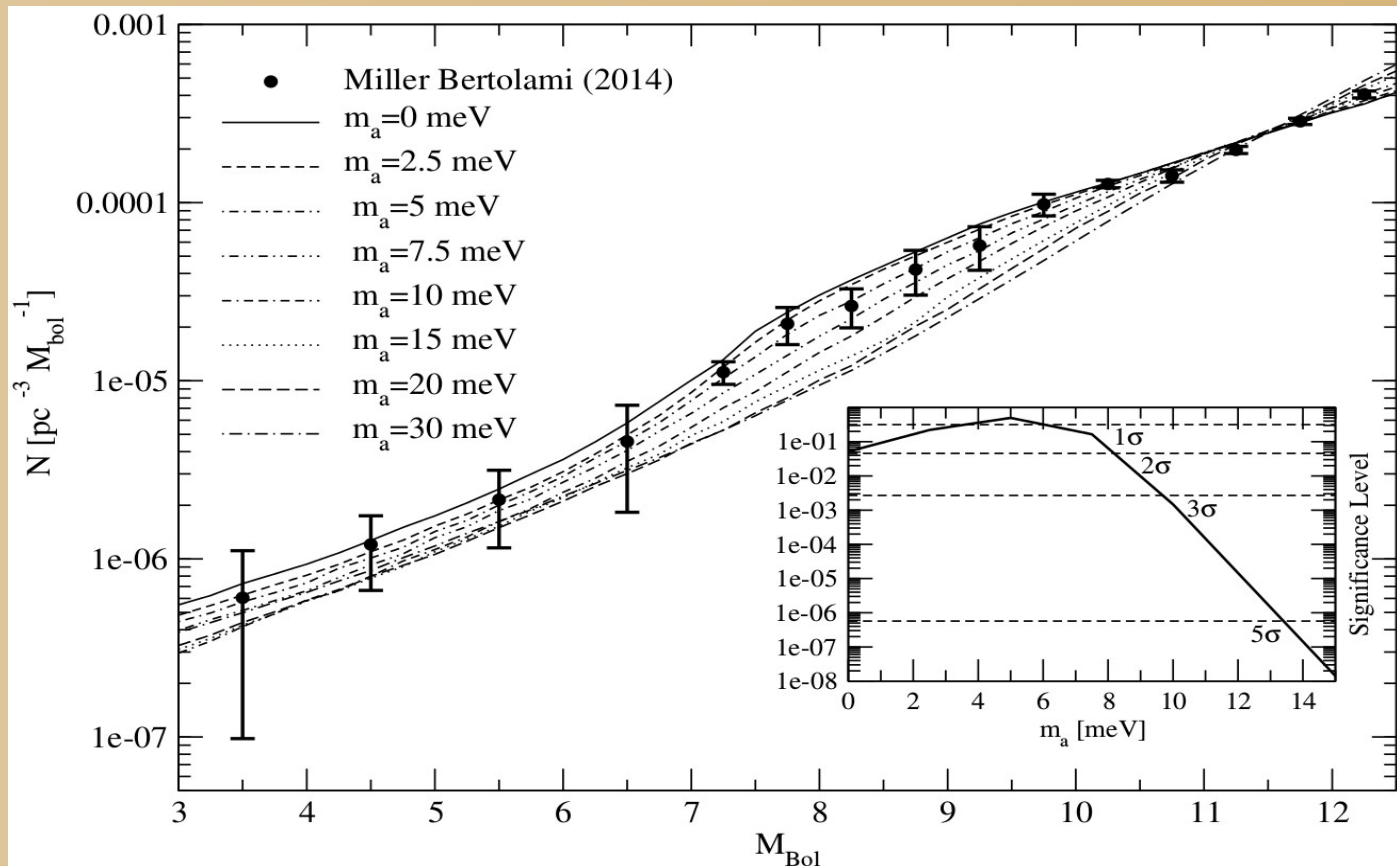
- The additional coolings may due to the **Axion-Like particle(ALP)** or **neutrino magnetic moment**.



M.M.M. Bertolami: 1407.1404

WDLF

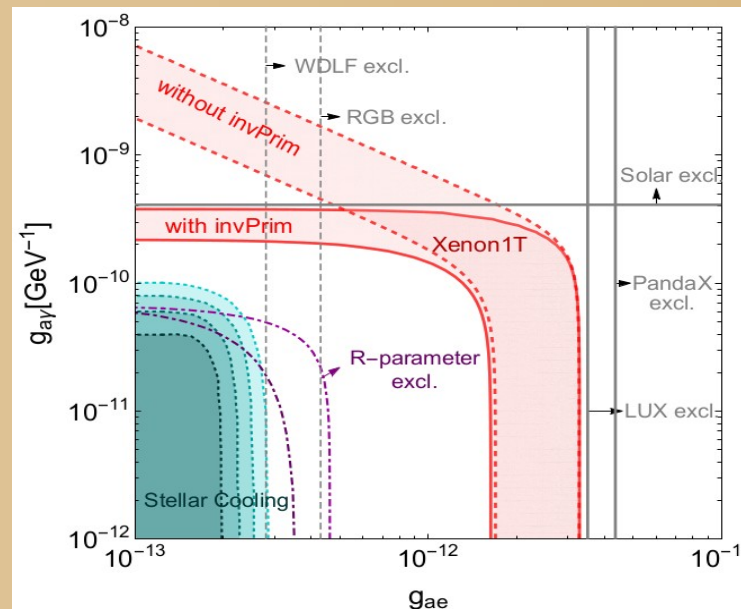
- The additional coolings may due to the **Axion-Like particle(ALP)** or **neutrino magnetic moment**.



M.M.M. Bertolami, B.E.Melendez, L.G.Althaus, J.Isern:
1406.7712
P.Y. Tseng,

XENON1T Anomaly

- XENON1T experiment reported 3-sigma excess in the electron recoil spectrum between 2-3 keV.
- Axion interpretation is not well consistent with constraints from stellar cooling



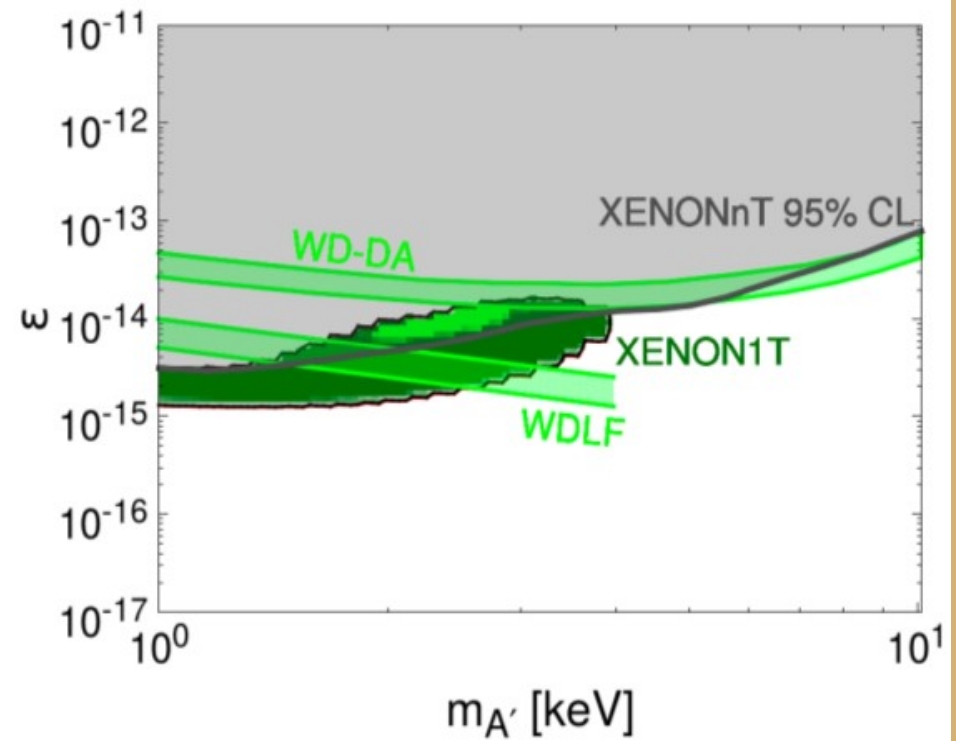
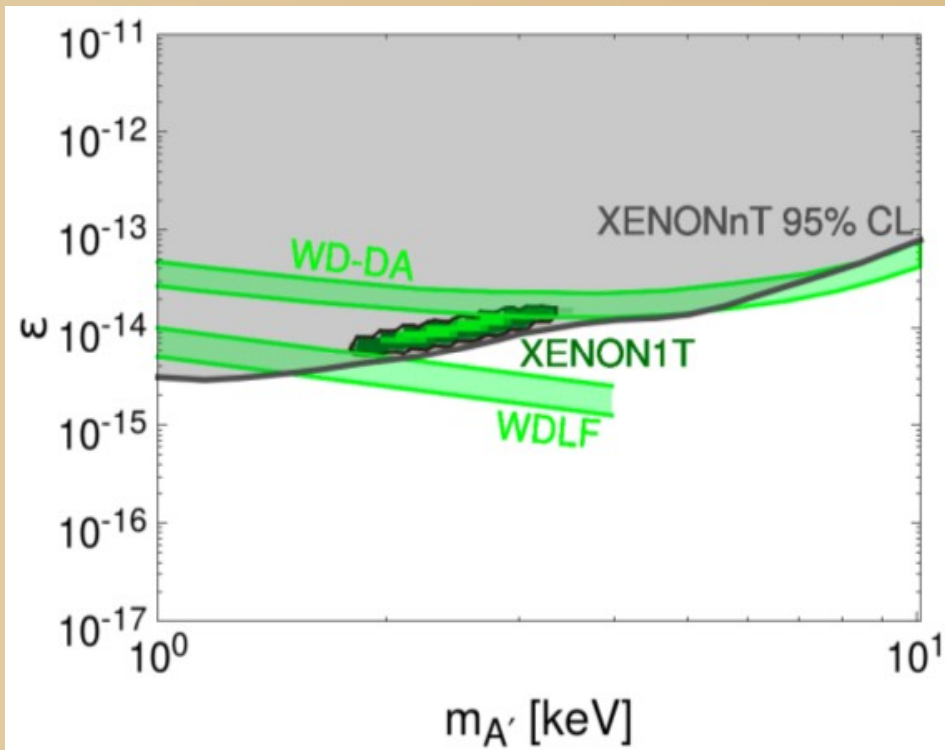
C.Gao, J.Liu, L.T.Wang, X.P.Wang, W.Xue, Y.M. Zhong :2006.14598

P.Y. Tseng,

p.28

XENONnT

- XENONnT project limit:



W.Y.Keung, D.Marfatia, P.Y.Tseng, arXiv:2009.04444