Electroweak axion string and superconductivity

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QCD Axion

- one of the most popular BSMs
- appears when one extends SM scalar sector and introduces Peccei-Quinn sym. $U(1)_{PQ}$
 - → SSB of $U(1)_{PQ}$ → NG boson = axion a
 - solves strong CP problem

 $|\theta_{QCD}| \lesssim 10^{-10}$

• axion *a* can be DM candidate



Axion string

• string-like soliton (vortex string) associated with $U(1)_{PQ}$ breaking

$$\phi(x) \sim v_{PQ} e^{i\theta}$$

Spatial angle θ is identified with phase of U(1)_{PQ} scalar φ
 → topologically non-trivial

(winding number 1)



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(winding number 1)

• Axion string is produced in the early universe if $U(1)_{PQ}$ is broken after inflation.

[Kibble, '80] [Zurek, '85]

Im ϕ

 $V(\phi)$

Re ϕ

String reconnection

• The most important property characterizing the string dynamics is **reconnection process.**



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 It has been believed that axion strings almost always reconnect when they collide.

Energy-loss mechanism of strings

• Thanks to the reconnection, the axion strings can loose their energy by theirselves.



Question: Do axion strings always reconnect? (It has been believed that this is true.)

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Our answer is NO! Key property is superconductivity!

Eg.) Witten's superconducting string

[Witten '80]

$$\begin{aligned} \mathscr{L} \supset |\partial_{\mu}\phi|^{2} + |D_{\mu}\sigma|^{2} - V(\phi,\sigma) & \phi \sim v_{\phi} e^{i\theta} \\ \swarrow & \swarrow & \swarrow & \swarrow & \downarrow \\ \text{EM neutral} & \text{EM charged} \\ V(\phi,\sigma) &= \lambda_{\phi}(|\phi|^{2} - v_{\phi}^{2})^{2} + \lambda_{mix}(|\phi|^{2} - v_{\phi}^{2})|\sigma|^{2} + m_{\sigma}^{2}|\sigma|^{2} + \lambda_{\sigma}|\sigma|^{4} \\ & \longrightarrow & -(\lambda_{mix}v_{\phi}^{2} - m_{\sigma}^{2})|\sigma|^{2} + \lambda_{\sigma}|\sigma|^{4} \\ \text{(inside the string, } \phi = 0) & \langle \sigma \rangle \neq 0 \quad \clubsuit \quad U(1)_{EM} \end{aligned}$$

- EM-charged scalar boson σ condensates only inside the string.
- → breaks $U(1)_{EM}$ inside the string
- → string becomes **superconducting state**

We found that, in DFSZ axion model, the axion strings can be superconducting string and may not reconnect successfully.

→Cosmological evolution of axion strings is nontrivial than previously thought!

- Introduction
- Axion string with electroweak flux
- Implication of superconductivity

• Summary

Axion string with electroweak flux

two Higgs doublets:
$$H_1, H_2$$
 + singlet scalar: S

	H_1	H_2	S
$SU(2)_W$	2	2	1
$U(1)_Y$	1	1	0

$$V(H_1, H_2, S) \supset \left(\kappa S^2 H_1^{\dagger} H_2 + \text{h.c.}\right) + \kappa_{1S} |H_1|^2 |S|^2 + \kappa_{2S} |H_2|^2 |S|^2$$

• global $U(1)_{PQ}$ symmetry:

$$S \to e^{i\alpha} S \qquad H_1 \to e^{i\alpha} H_1 \qquad H_2 \to e^{-i\alpha} H_2$$

• assume VEVs:

$$\langle S \rangle = v_{PQ} \qquad \langle H_1 \rangle = (0, v_1)^T \qquad \langle H_2 \rangle = (0, v_2)^T$$

 $v_{PQ} \sim 10^{9-12} \,\text{GeV}$ $v_{EW}^2 = 2(v_1^2 + v_2^2) \simeq (246 \,\text{GeV})^2$

In the following, $v_1 = v_2$ for simplicity

Axion string in DFSZ model

High Temperature: $v_{EW} \ll T \leq v_{PQ}$

- $U(1)_{PQ}$ is broken, but electroweak sym. remains \rightarrow ordinary axion string in DFSZ model
- The Higgs doublets acquire **no VEVs.**

$$S \sim v_{PQ} e^{i\theta}$$
$$H_1 = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \qquad H_2 = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

 $x + iy = re^{i\theta}$



Axion string with EW flux (EW exion string) Low temperature: $T \leq v_{EW} \rightarrow EW$ as U(2)• The Higgs doublets also acquire note zero VEVs.

 \rightarrow The doublets can have winding phases \rightarrow EW flux tube



• There are three types of EW axion strings, but I will show one of them. (please see our paper)

EW agion string is superconducting!



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Implication of superconductivity

Primordial magnetic field induce current

- assume existence of primordial magnetic field in early universe
- When a superconducting string moves in the background of PMF, electric current is induced.



• Using observational upper bound on the magnetic field: $B_0 \lesssim 10^{-11}$ (Gauss),

$$J_{em} \sim e^2 B \beta t \lesssim 10^{12} \, ({\rm GeV})$$

Large electric current can be induced.

Interaction between the strings

• Current-induced magnetic interaction between the strings



$$F_{axion} \sim v_{PQ}^2 / R$$

- The magnetic interaction can dominate the string interaction.
- → affect the reconnection process of the strings

• Due to such a long-range attractive force, the strings can form **a bound state, called Y-junction.**



→ Reconnection doesn't work

Superconductivity may spoil the reconnection property!

Probability of formation of Y-junctions

• The probability of the formation of Y-junctions is roughly estimated to be $\sim 1/2$.



Probability of formation of Y-junctions

• When an EW axion string intersects with itself, it creates an exotic object with probability 1/2.



Cosmological evolution is non-trivial any more!

- Reconnection of axion strings is a very important process to avoid dominating the energy of the universe.
- In DFSZ model, axion strings obtain electroweak gauge fluxes after the EW sym. breaking. (EW axion string)
- EW axion string breaks $U(1)_{EM}$ inside it, leading to superconductivity.
- Assuming background of PMF, the strings do not reconnect successfully, but form a bound state (Y-junction).
- → affect cosmology?

Backup

Three types of EW axion strings

• type-A string :

$$S \sim v_{PQ} e^{i\theta}$$
 $H_1 \sim e^{i\theta} \begin{pmatrix} 0 \\ v \end{pmatrix}$ $H_2 \sim e^{-i\theta} \begin{pmatrix} 0 \\ v \end{pmatrix}$

• type-B string :

$$S \sim v_{PQ} e^{i\theta} \qquad H_1 \sim e^{i\theta} e^{-i\theta\sigma_3} \begin{pmatrix} 0 \\ v \end{pmatrix} \qquad H_2 \sim e^{-i\theta} e^{-i\theta\sigma_3} \begin{pmatrix} 0 \\ v \end{pmatrix}$$

type-C string :

$$S \sim v_{PQ} e^{i\theta} \qquad H_1 \sim e^{i\theta} e^{i\theta\sigma_1} \begin{pmatrix} 0 \\ v \end{pmatrix}$$

$$H_2 \sim e^{-i\theta} e^{i\theta\sigma_1} \begin{pmatrix} 0 \\ v \end{pmatrix}$$
$$U(1)_{PQ} \quad U(1)_{W^1}$$

 $U(1)_{PQ}$ $U(1)_Z$

 $U(1)_{PQ}$ winding

String tensions

• Compare the tensions of the three strings



- Type-C string can be the most stable string
 → not exotic, but energetically favored in some parameter range!
- In the range, the axion string necessarily becomes type-C string after the EW sym. breaking.

Numerical solution of type-C string

Again we take $\tan \beta = 1 \rightarrow Z$ flux vanishes.

$$H_{1} = \frac{1}{2}v_{1} \begin{pmatrix} f(r)e^{2i\theta} - h(r) \\ f(r)e^{2i\theta} + h(r) \end{pmatrix} \qquad H_{2} = \frac{1}{2}v_{2} \begin{pmatrix} h(r) - f(r)e^{-2i\theta} \\ h(r) + f(r)e^{-2i\theta} \end{pmatrix}$$

$$S = v_{PQ} \phi(r) e^{i\theta}$$

$$W_{\theta}^{1} = \frac{-4\pi}{g}(1 - w(r))$$



 $v_{PQ} = 10 v_1 \qquad m_h^2 = (125 \text{ GeV})^2 \qquad \tan \beta = 1$ $\kappa_{1S} = \kappa_{2S} = 0.4 \qquad \kappa = -2(v/v_{PQ})^2 \quad \lambda_S = 1$ length unit: $v_1 = 0.2$

$U(1)_{EM}$ is broken inside the string!

• To see the inside of the string, we should use the "smeared" ansatz:

$$H_{1} = \frac{1}{2}v_{1} \begin{pmatrix} f(r)e^{2i\theta} - h(r) \\ f(r)e^{2i\theta} + h(r) \end{pmatrix} \qquad H_{2} = \frac{1}{2}v_{2} \begin{pmatrix} h(r) - f(r)e^{-2i\theta} \\ h(r) + f(r)e^{-2i\theta} \end{pmatrix}$$

$$\begin{split} \hat{Q}_{EM} &: U(1)_{EM} \text{ generator} \\ \hat{Q}_{EM} H_i &\propto f - h \begin{cases} \rightarrow 0 & (r \rightarrow \infty) \\ \neq 0 & (r \lesssim v_{\text{EW}}^{-1}) \end{cases} \end{split}$$

 $U(1)_{EM}$ is spontaneously broken!

→ superconducting state!



Probability of formation of Y-junctions

• Reduce the 3D collision to two 2D collisions



• vortex-vortex collision forms a bound sate for $\alpha \gtrsim 10^{-3}$