

Hubble Selection of the Weak Scale

Possibility from QCD quantum phase transition

TaeHun Kim,
work with Sunghoon Jung

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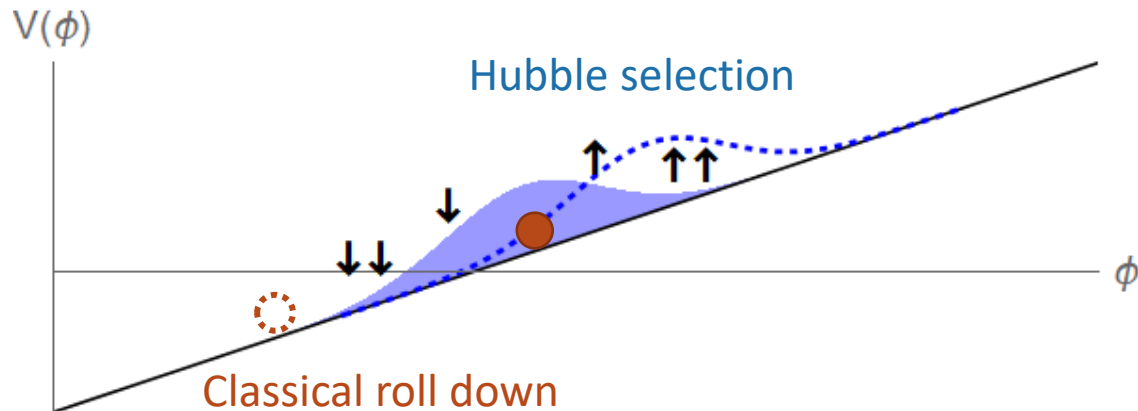
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Hubble selection: Concept

- Inflationary fluctuation of quantum field can “kick” the field value toward higher potentials, against the classical rolling.
- Having higher Hubble rate in higher potentials, a global distribution of field values can behave radically different from the field value of each point;

“Hubble selection”



Hubble selection: History

- Eternal inflation
 - “Self reproduction of inflating universe”
 - “Every local universe reaches reheating, but the inflation never ends globally”
 - “Dominant volume of the universe is always in inflationary stage”
- Recent developments: stochastic axion, SOL, ...

Hubble selection: History

- Self organized localization (SOL) (Giudice et al. 2021)
 - Hubble selection in non-dynamical but equilibrium point of view
 - Driven by Hubble selection (“self organized”), the field value distribution has equilibrium near the boundary of potential (“localization”)
- We give dynamical viewpoint to Hubble selection

Hubble selection: Basic ingredients

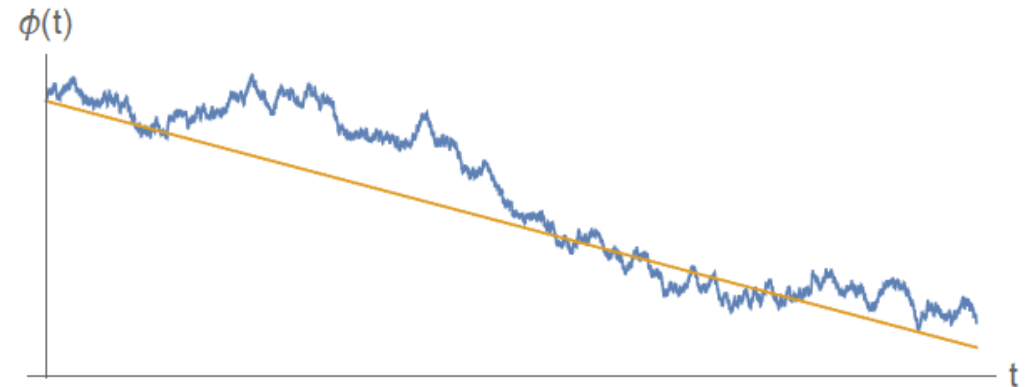
- Quantum fluctuation of fields
 - Inflation \rightarrow de Sitter space time \rightarrow quantum fluctuations are generated
 - Modes exiting the Hubble horizon (when $\lambda \sim \frac{1}{H}$) “freeze”
 - Spatial field profile: each “Hubble patches” of size $\frac{1}{H}$ acquires its own field value coming from the accumulation of horizon crossed modes
 - Different field values for different patches; we look at **patch-by-patch field values**

Hubble selection: Basic ingredients

- Stochastic motion of field value of a patch
 - Continuous accumulation of horizon exiting modes
 - “Continuous random kicks” for field value: the Brownian motion

- $$d\phi = -\frac{V'}{3H} dt + \sqrt{\frac{H^3}{4\pi^2}} dW, \quad \langle dW \rangle = 0, \quad \langle dW^2 \rangle = dt$$

- Classical rolling + **stochastic motion**



Hubble selection: Basic ingredients

- Distribution of field values among different patches?
 - “Is our universe special?” \Leftrightarrow “Where in the distribution corresponds to our universe?”
 - Langevin equation $\phi(t) \Leftrightarrow$ Fokker Planck equation $\rho(\phi, t)$
 - But Hubble expansion comes in...
 - Different $H(\phi)$ for different ϕ 's, due to different $V(\phi)$

Hubble selection: FPV equation

- Volume weighted Fokker-Planck (FPV) equation

$$\frac{\partial \rho(\phi, t)}{\partial t} = \frac{\partial}{\partial \phi} \left(\frac{V'}{3H} \rho \right) + \frac{1}{8\pi^2} \frac{\partial^2 (H^3 \rho)}{\partial \phi^2} + 3\Delta H(\phi) \rho$$

Hubble selection: FPV equation

- Volume weighted Fokker-Planck (FPV) equation

Ordinary Fokker Planck eq.

$$\frac{\partial \rho(\phi, t)}{\partial t} = \underbrace{\frac{\partial}{\partial \phi} \left(\frac{V'}{3H} \rho \right)}_{\text{Drift (classical rolling)}} + \underbrace{\frac{1}{8\pi^2} \frac{\partial^2 (H^3 \rho)}{\partial \phi^2}}_{\text{Diffusion (quantum fluctuation)}} + 3\Delta H(\phi)\rho$$

Drift
(classical rolling)

Diffusion
(quantum fluctuation)

Hubble selection: FPV equation

- Volume weighted Fokker-Planck (FPV) equation

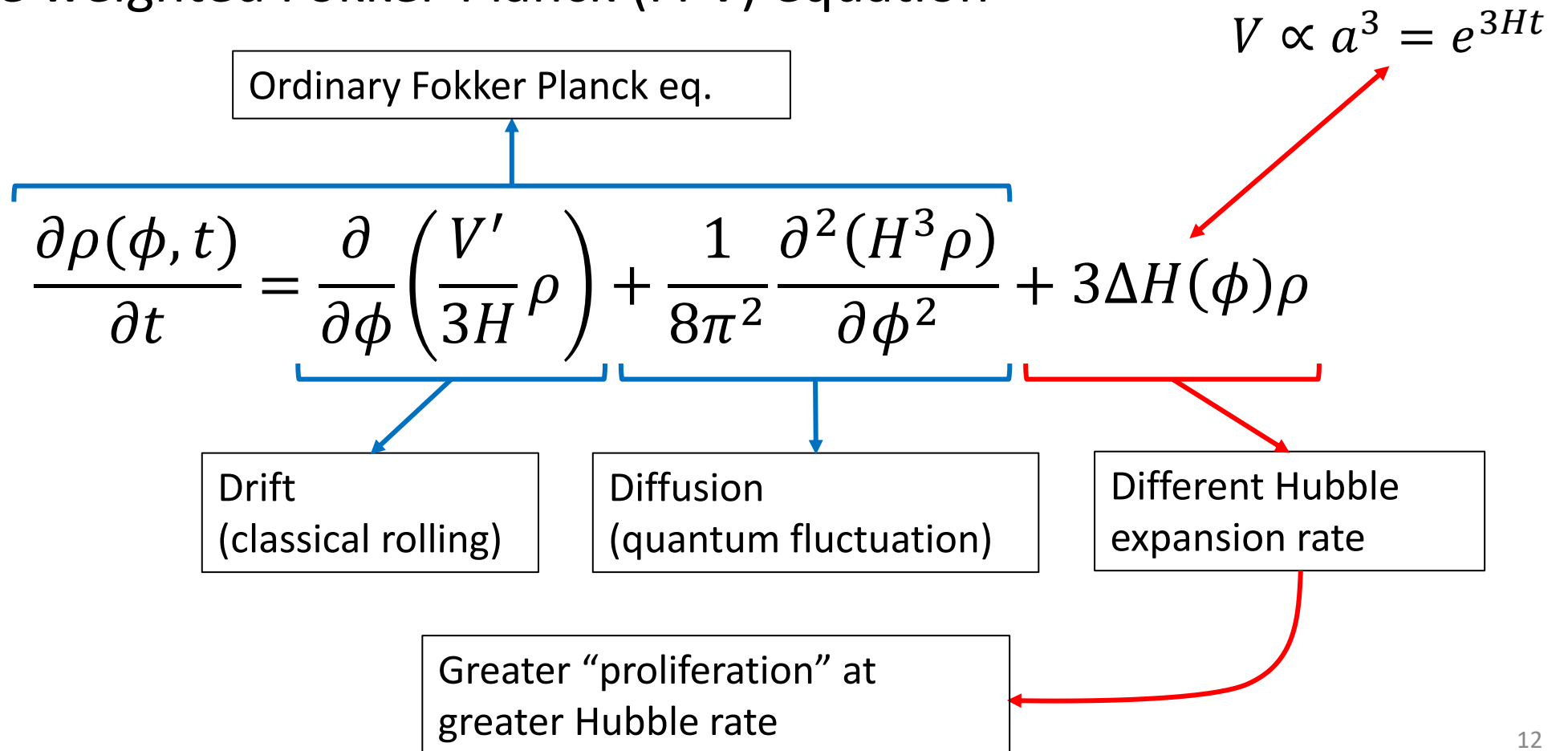
Ordinary Fokker Planck eq.

$$\frac{\partial \rho(\phi, t)}{\partial t} = \underbrace{\frac{\partial}{\partial \phi} \left(\frac{V'}{3H} \rho \right)}_{\text{Drift (classical rolling)}} + \underbrace{\frac{1}{8\pi^2} \frac{\partial^2 (H^3 \rho)}{\partial \phi^2}}_{\text{Diffusion (quantum fluctuation)}} + \underbrace{3\Delta H(\phi) \rho}_{\text{Different Hubble expansion rate}}$$

$V \propto a^3 = e^{3Ht}$

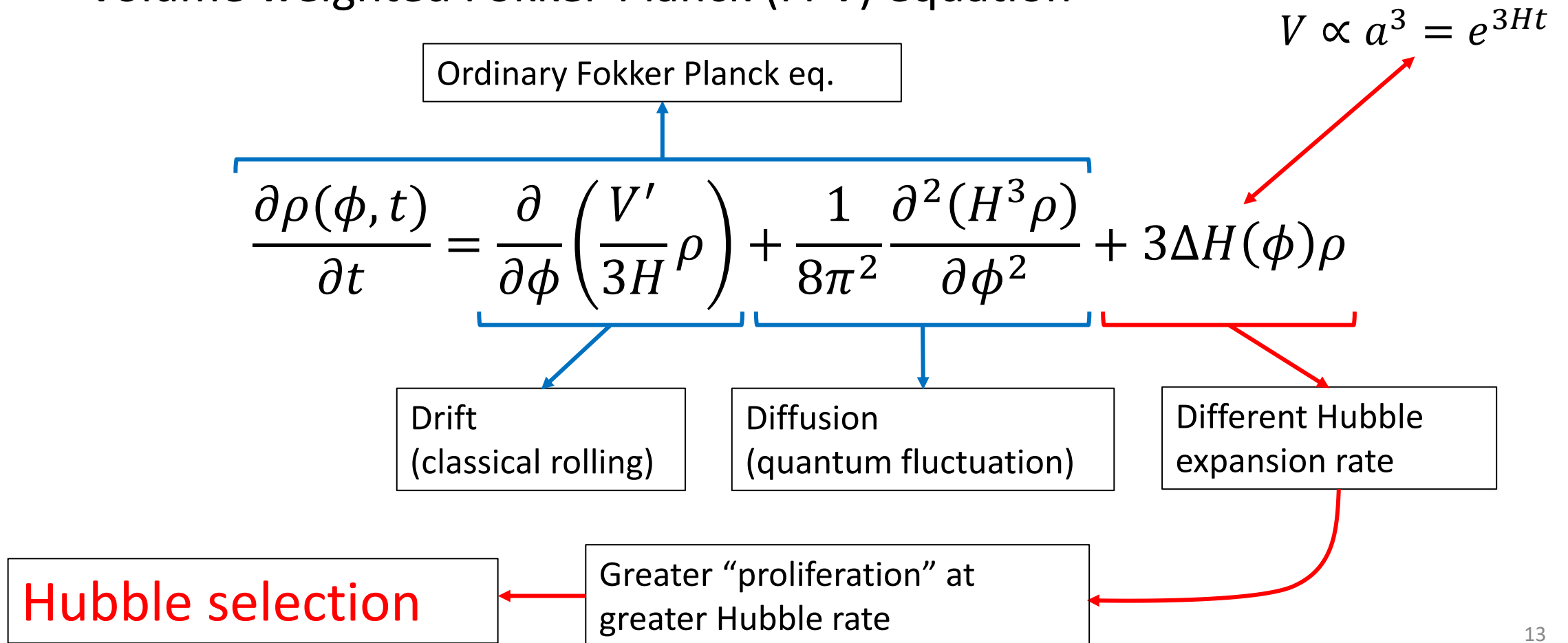
Hubble selection: FPV equation

- Volume weighted Fokker-Planck (FPV) equation



Hubble selection: FPV equation

- Volume weighted Fokker-Planck (FPV) equation



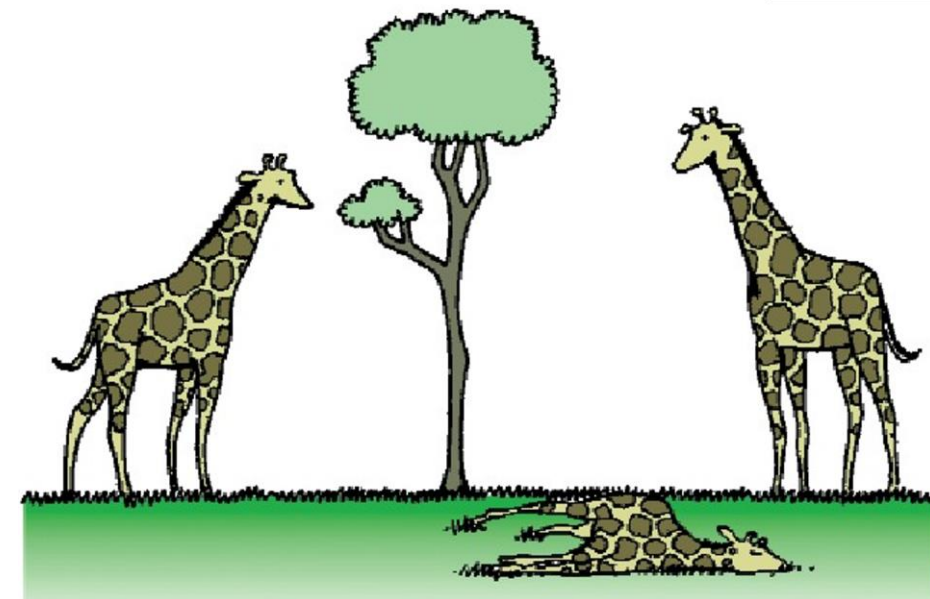
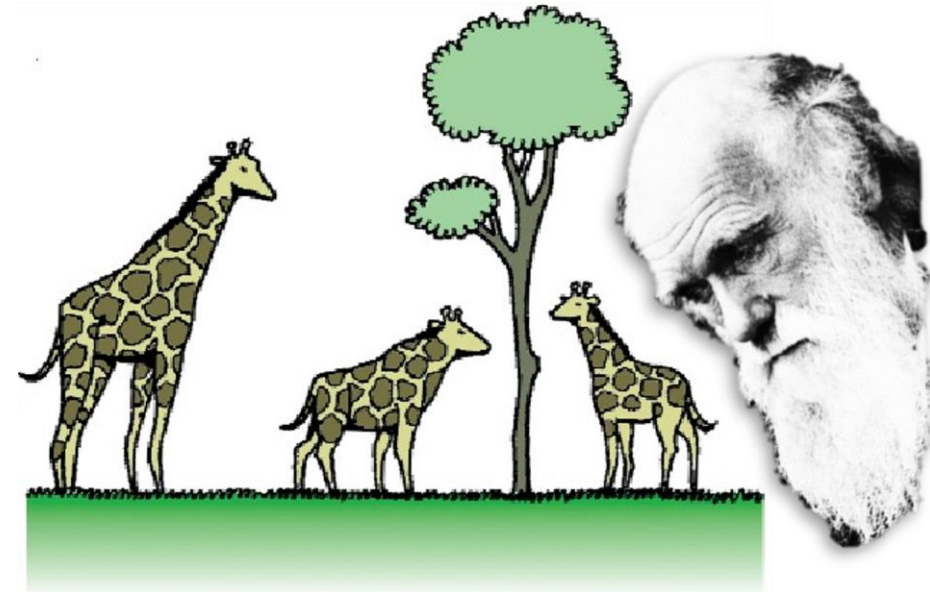
Hubble selection: FPV equation

- $V(\phi) = V_0 + \Delta V(\phi) \quad V_0 \gg \Delta V$
 - Total potential = Sum of inflaton potential (V_0) + ϕ contribution (ΔV)
- $\Delta H(\phi) = \frac{\Delta V(\phi)}{6M_P^2 H_0} \propto \Delta V(\phi) \quad H_0 = \sqrt{\frac{V_0}{3M_P^2}}$

Hubble selection: Analogy

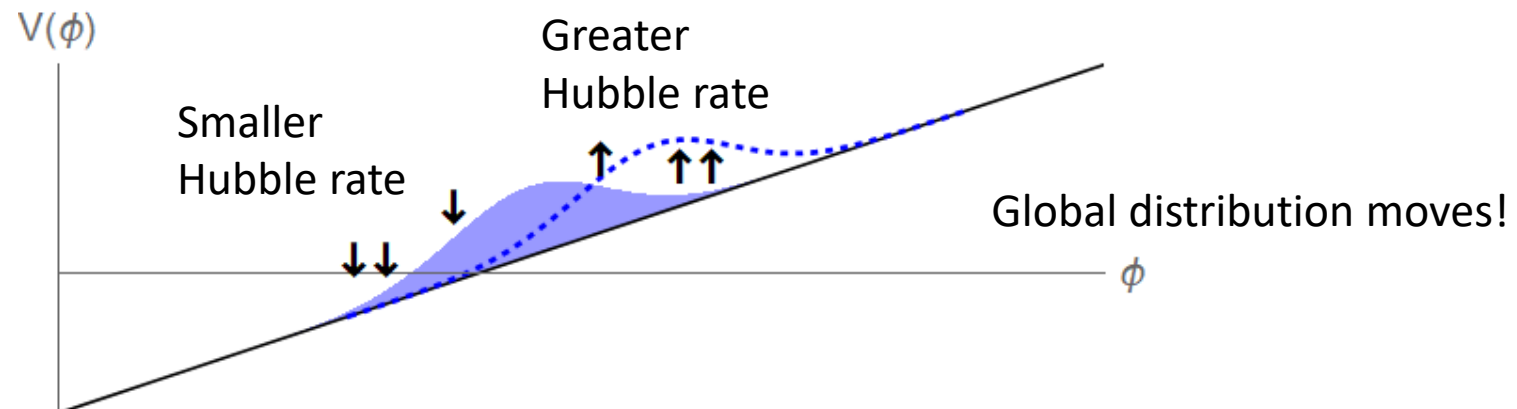
- Naming comes from “natural selection”

| | Natural selection |
|------------------------|---|
| Applied to | Biological creatures |
| “Offspring” production | Reproduction |
| ... rate differs by | Adaptation to environment |
| Diversity comes from | Genetic variation (mutation) |
| Result | Dominance of genotype with higher reproduction rate |



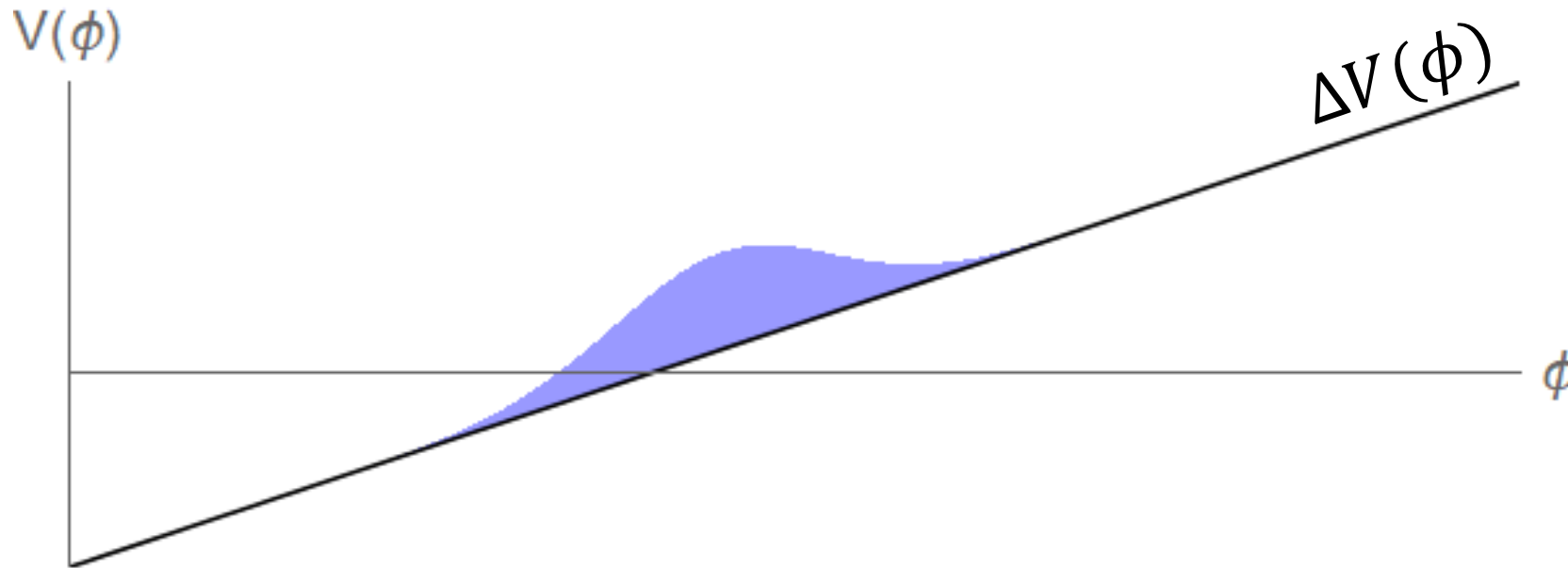
Hubble selection: Analogy

| | Natural selection | Hubble selection |
|------------------------|---|--|
| Applied to | Biological creatures | Hubble patches |
| “Offspring” production | Reproduction | Hubble expansion |
| ... rate differs by | Adaptation to environment | $\Delta H(\phi) = \frac{\Delta V(\phi)}{6M_P^2 H_0}$ |
| Diversity comes from | Genetic variation (mutation) | Quantum fluctuation |
| Result | Dominance of genotype with higher reproduction rate | Dominance of field values with higher Hubble rate |



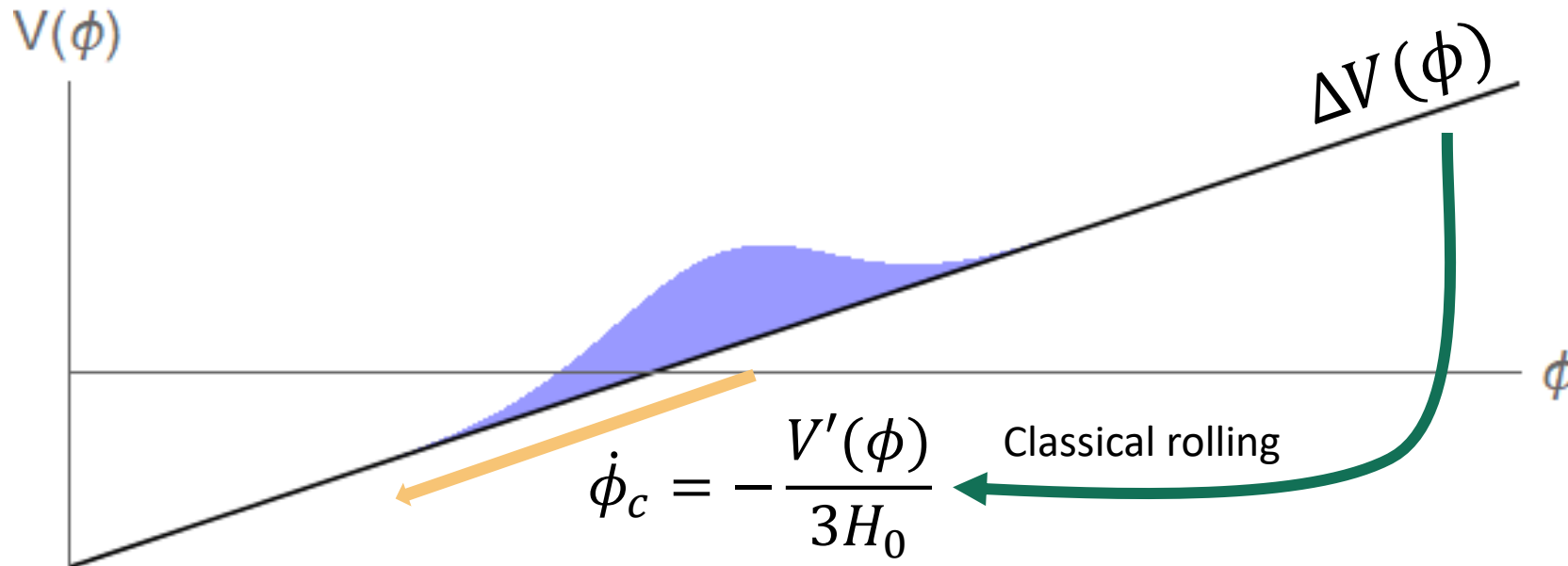
Hubble selection: Quantitative analysis

- Quantitative description
 - Classical rolling vs Hubble selection: inevitable competition



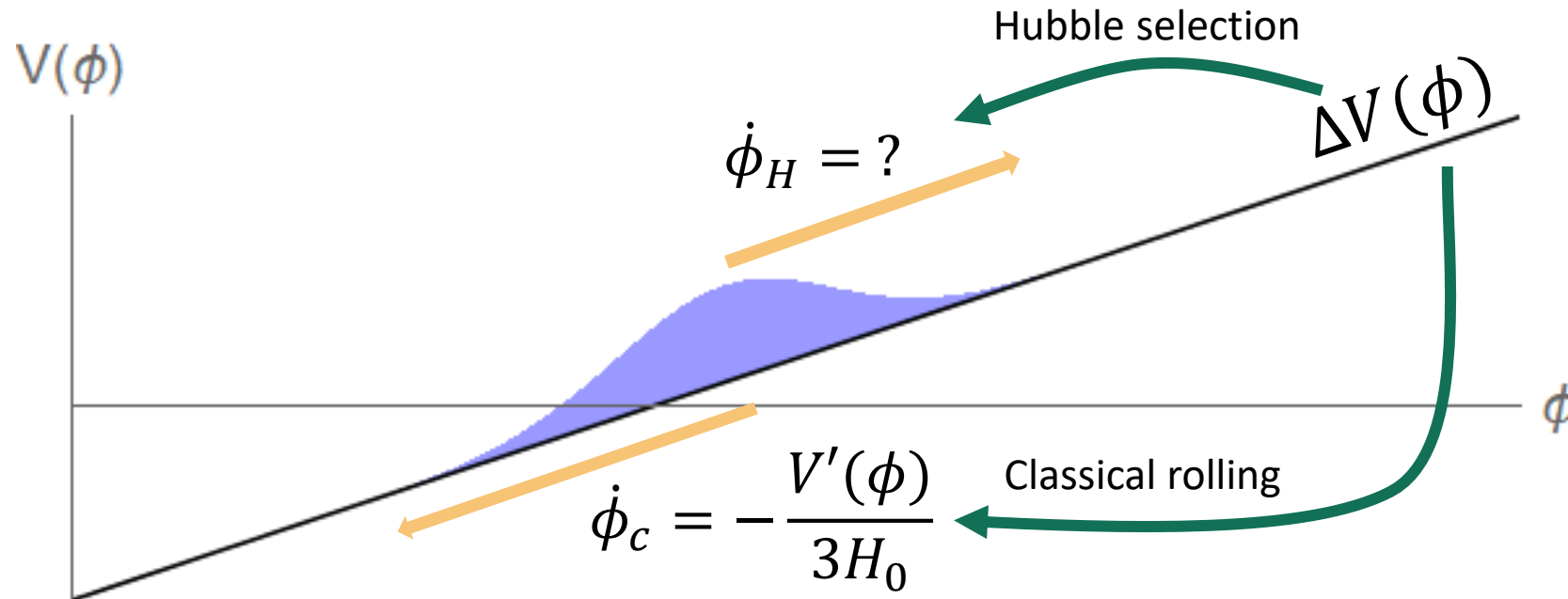
Hubble selection: Quantitative analysis

- Quantitative description
 - Classical rolling vs Hubble selection: inevitable competition



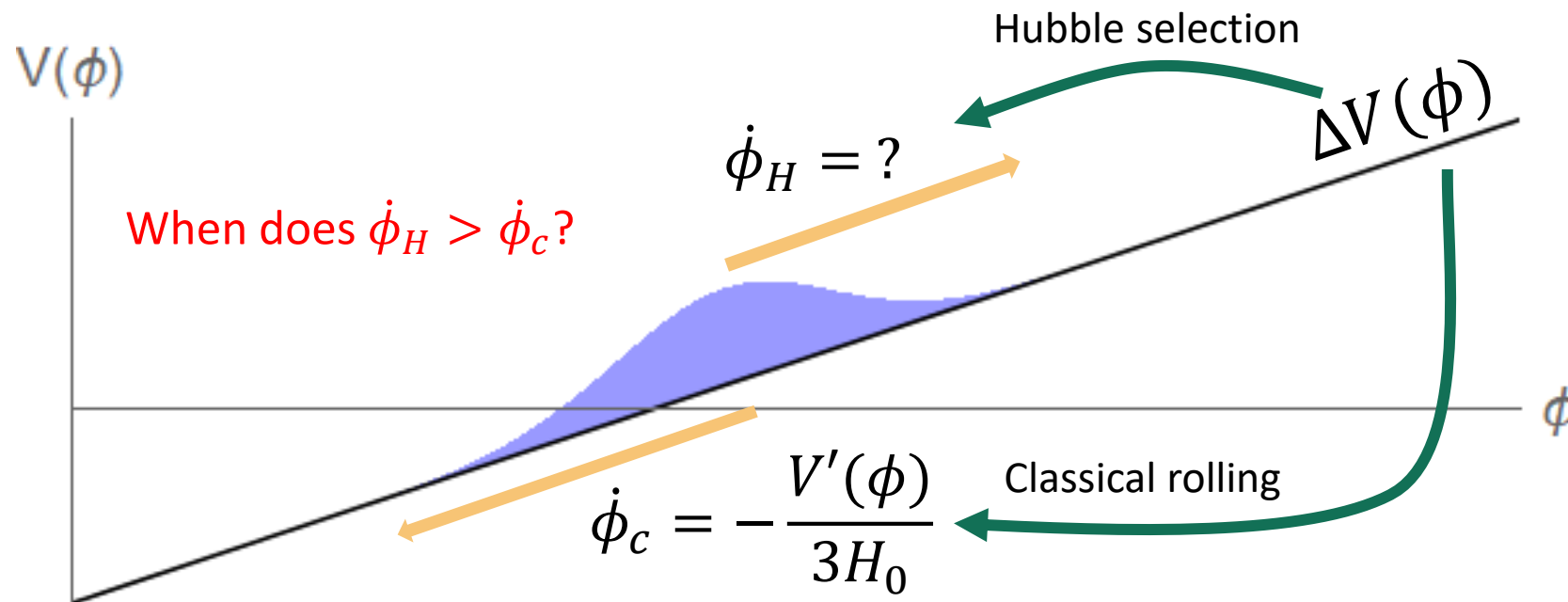
Hubble selection: Quantitative analysis

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Hubble selection: Quantitative analysis

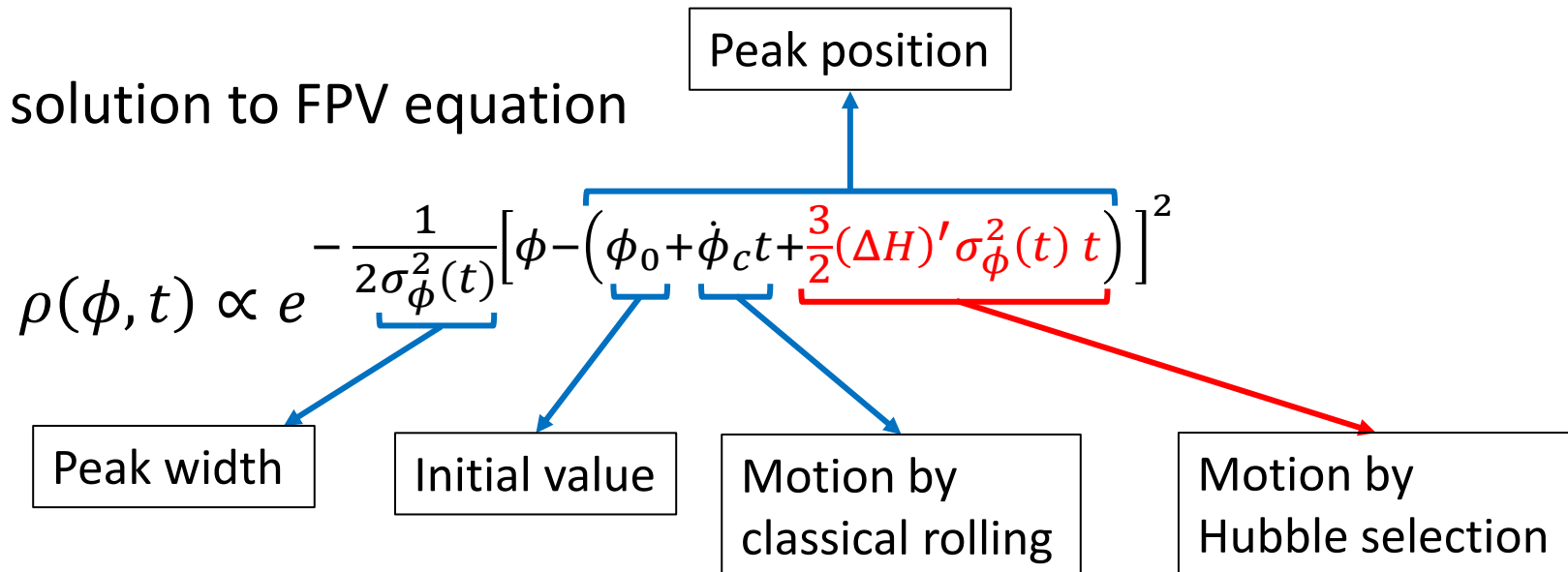
- Quantitative description
 - Classical rolling vs Hubble selection: inevitable competition



Hubble selection: Quantitative analysis

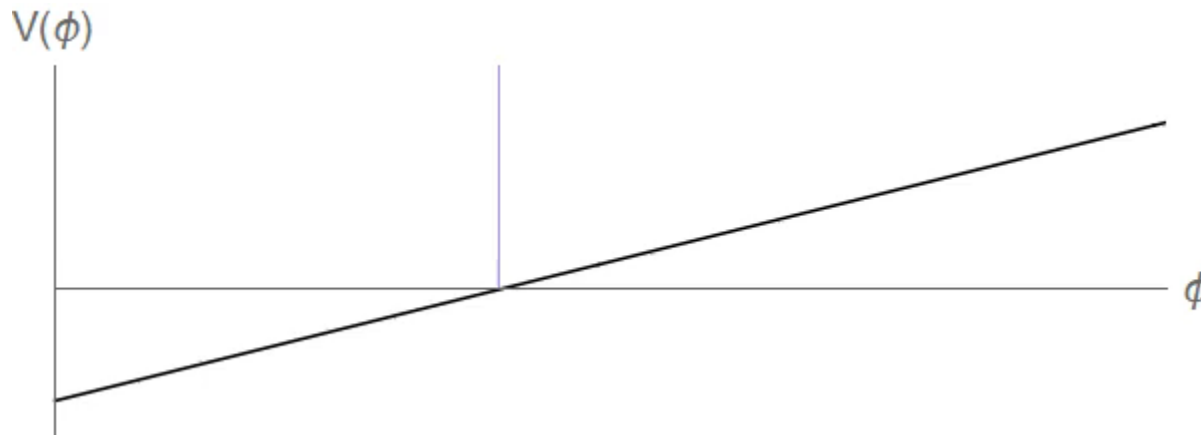
- Recall: necessary ingredient for selection: “diversity”
 - ΔH between the “head” and the “tail” drives the distribution
 - Efficient for broadened distributions

- Linear potential solution to FPV equation



Hubble selection: Quantitative analysis

- $\sigma_\phi^2(t) = \frac{H^3}{4\pi^2} t$: variance (diffusion by stochastic motion)
- $\dot{\phi}_H = 3(\Delta H)' \sigma_\phi^2 = \frac{V' \sigma_\phi^2}{2M_P^2 H_0}$: motion induced by Hubble selection
 - $\propto \sigma_\phi^2$: “need for diversity”, $\propto V'$: “strength of selection”



Hubble selection: Quantitative analysis

- Turning point: $\dot{\phi}_c + \dot{\phi}_H = 0$

- Planckian width: $\sigma_\phi \sim M_P$

- Maximum roll down excursion: $\Delta\phi > M_P \sim \sigma_\phi$

- Eternal inflation: $\Delta N \simeq \frac{8\pi^2 M_P^2}{3H^2} > \frac{2\pi^2 M_P^2}{3H^2} = \text{“de Sitter entropy bound”}$

- **Super Planckian field range** and **eternal inflation** are necessary

- No Hubble selection for usual well-known quantum fields

Condition to climb up:
Full field range accommodates $\Delta\phi$

Hubble selection: Equilibrium

- **Equilibrium:** localization near the boundary
 - out of EFT / potential drop (phase transition) / anything else...
 - Distribution cannot get over and stop near the boundary
- Why equilibrium is important?
 - Eternal inflation: reheating universe is continuously dominated by “young population”
 - After finite time, newly generated (and dominating) young population will follow the equilibrium distribution.

Hubble selection: Equilibrium

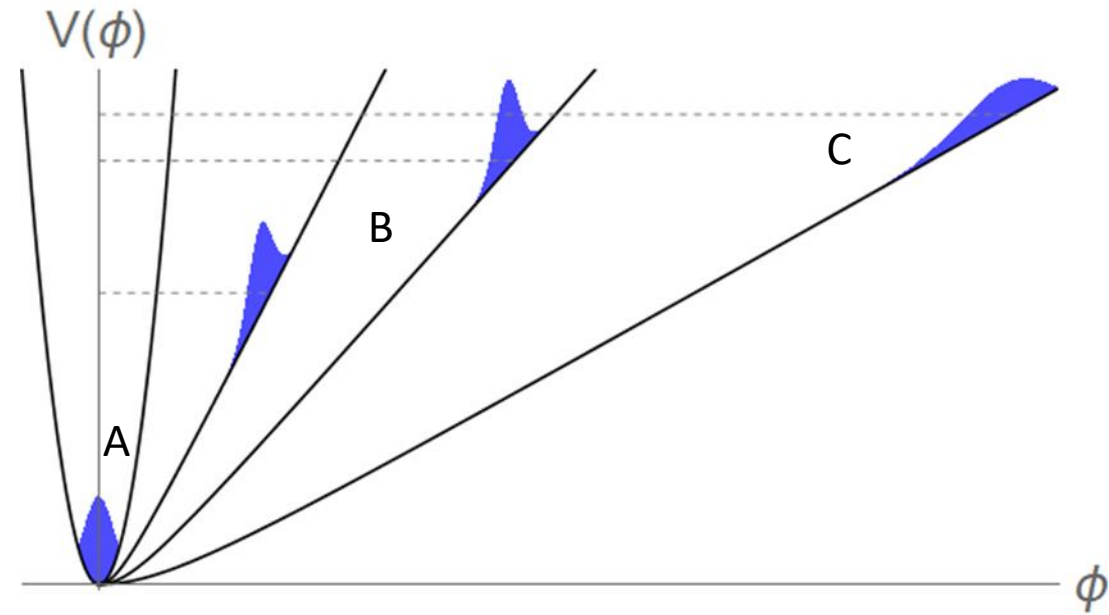
- **Width** and **position** of equilibrium distribution
 - How are they determined? We give qualitative description
 - Full quantitative description: Giudice et al. 2021

Hubble selection: Equilibrium

- Rule of thumb:

less steep potential gives equilibrium

distribution **closer to the upper boundary**



Hubble selection: Equilibrium

- A: Classically dominated

- $\dot{\phi}_c \gg \dot{\phi}_H$

- B: Quantum dominated

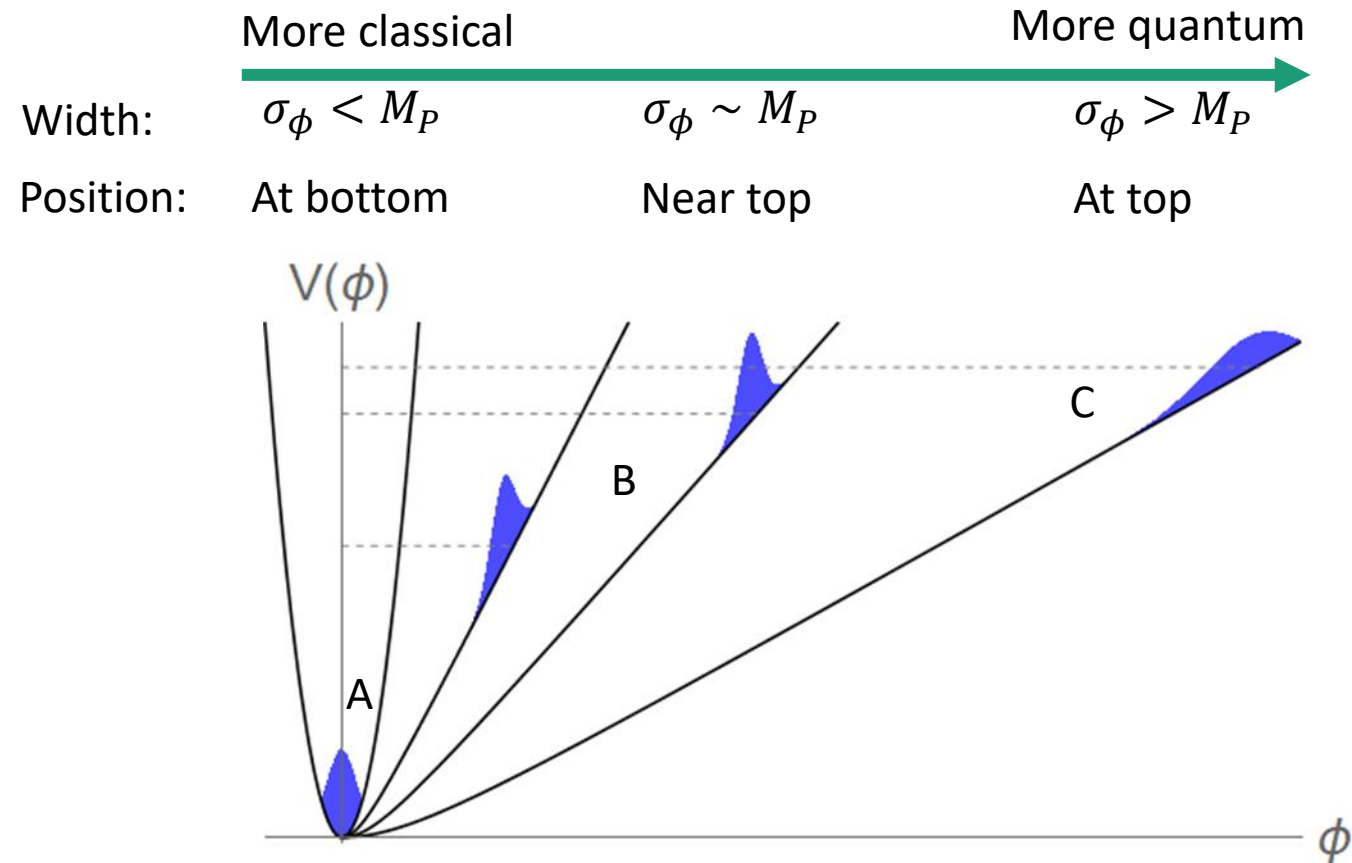
- $\dot{\phi}_c + \dot{\phi}_H = 0$

- C: Extremely quantum dominated

- $\dot{\phi}_c + \dot{\phi}_H + \dot{\phi}_b = 0$

- C/QV/Q2V regimes in SOL

- Quantitatively consistent



Possibility from QCD

- Can we use Hubble selection to select the weak scale?
 - “Localization near the upper boundary”
- 1st order quantum phase transition gives the boundary
 - Vacuum energy is peaked at the critical point
- QCD chiral phase transition could be 1st order, at QCD scale
 - Not far from the current weak scale
 - Subject to strange quark mass
 - Not firmly established yet... thus a “possibility”

Possibility from QCD

- Need three sectors
 - ϕ : relaxion; scans the Higgs mass (or VEV v_h)
 - h : Higgs; v_h determines quark mass; triggers Σ 's phase transition at $v_h^* = v_h(\phi^*)$
 - Σ : meson field; QCD d.o.f. below $\Lambda_{QCD} \simeq 200$ MeV; undergoes 1st order phase transition

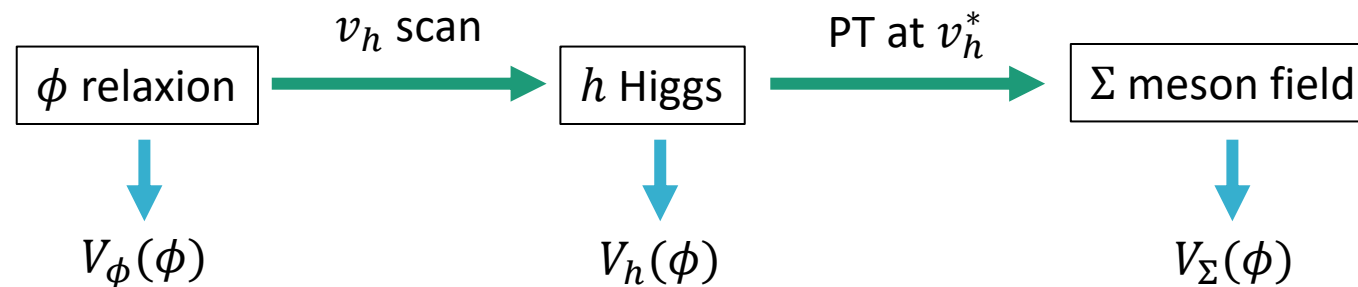
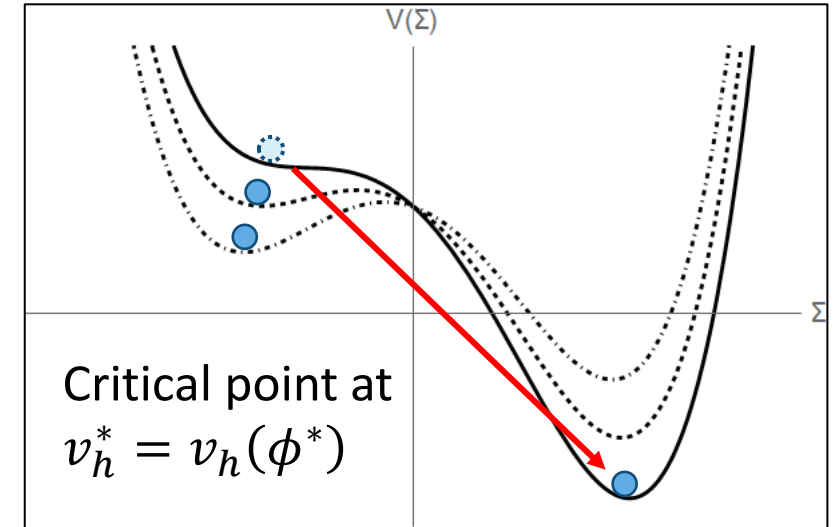


- h and Σ always follow their minimum; no Hubble selection for them

Possibility from QCD

- $V = V_\phi(\phi) + V_h(\phi) + V_\Sigma(h(\phi))$
- $V_\phi = \Lambda_\phi^4 \cos\left(\frac{\phi}{f_\phi}\right), \quad V_h = \frac{1}{2}(M^2 - g\phi)h^2 + \frac{\lambda_h}{4}h^4$
- $V_\Sigma = \mu^2 \text{Tr}(\Sigma\Sigma^\dagger) + \dots - \text{Tr}(\mathcal{H}(\Sigma + \Sigma^\dagger))$

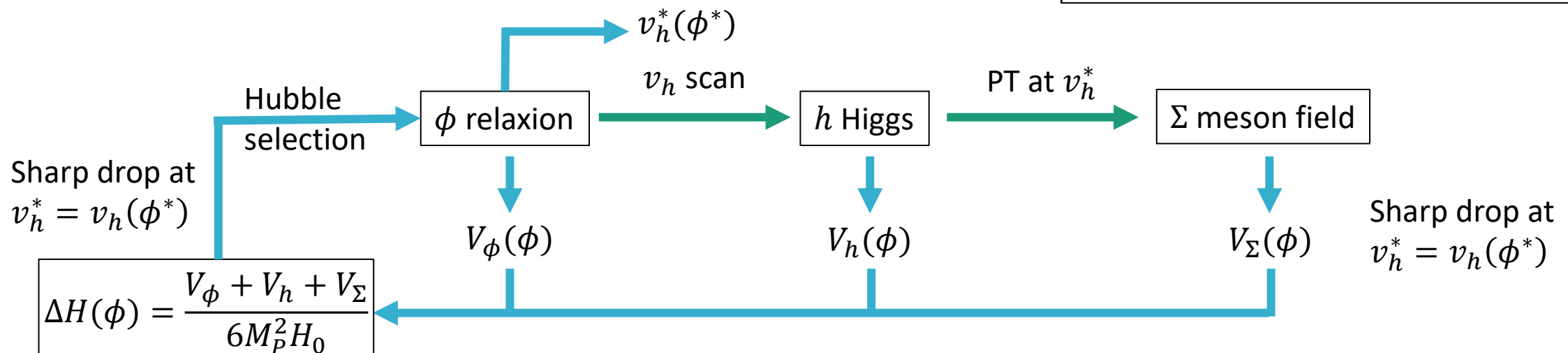
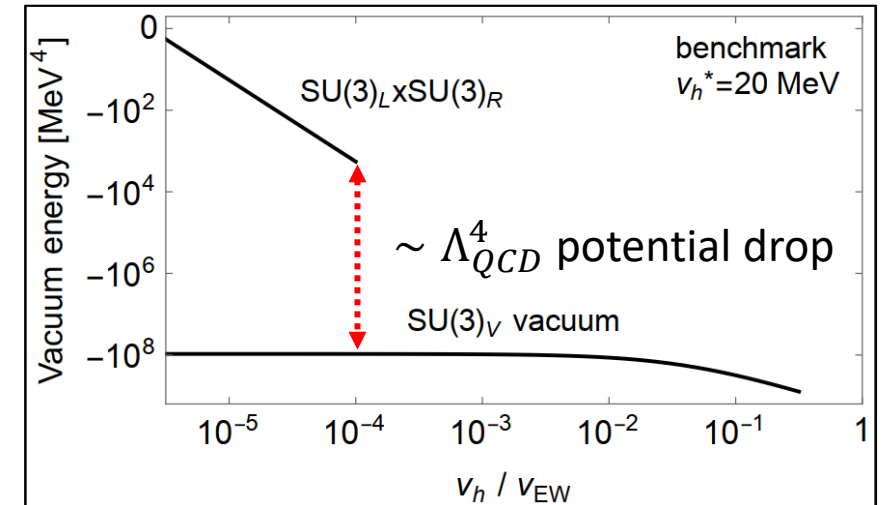
$\mathcal{H}(v_h)$ plays role of external magnetic field in magnetization



Possibility from QCD

- $V = V_\phi(\phi) + V_h(\phi) + V_\Sigma(h(\phi))$
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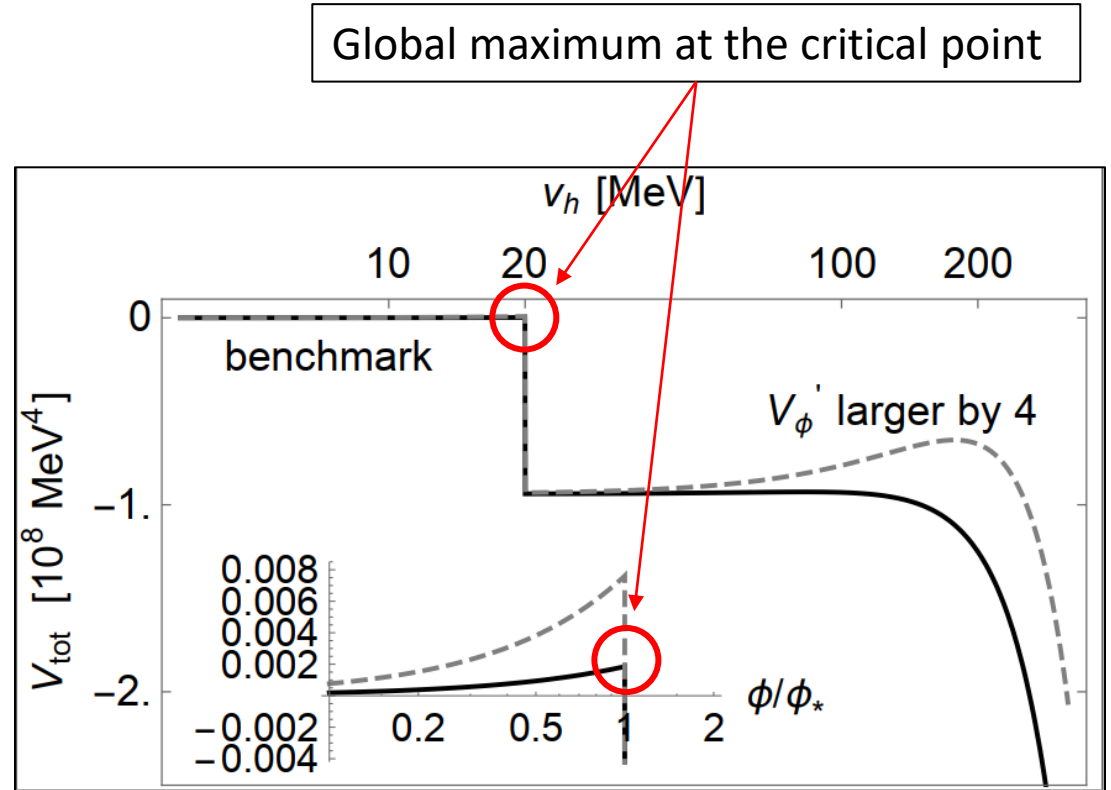
Σ vacuum structure



Possibility from QCD

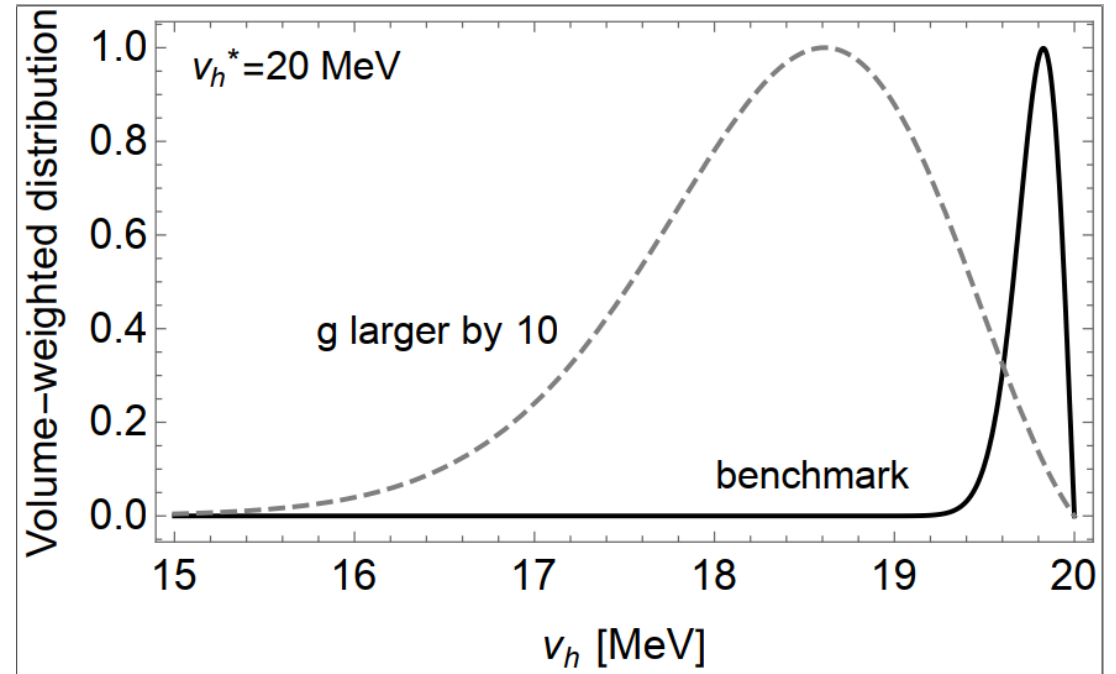
- Required conditions
 - ...
 - $V(\phi^*)$ is global maximum (greatest Hubble rate):
$$v_h^* \lesssim \Lambda_\phi^2/M \lesssim \Lambda_{QCD}$$

“Slope of V_ϕ should not be too small nor too large”
- $\Lambda_\phi \ll M$ is need to have high cutoff



Possibility from QCD

- Successful benchmark
 - $H = v_h^* \simeq 20 \text{ MeV}$
 $M = 3 \times 10^{-3} M_P$
 $\Lambda_\phi^2 = 10^{-2} H M_P$
 $g = 10^{-3} H^2 / M_P$
- Well localized v_h near $v_h^* \simeq 20 \text{ MeV}$



Possibility from QCD

- High cutoff $M = 3 \times 10^{-3} M_P$, well localized $v_h \dots$

Are we successfully solved the Higgs naturalness problem? Unfortunately, **no**.

- $\Lambda_\phi \ll M$ is unstable from quantum correction originating from relaxion-Higgs interaction; at least $\Lambda_\phi \sim M$

- Recall: $V_h = \frac{1}{2}(M^2 - g\phi)h^2 + \frac{\lambda_h}{4}h^4$

Possibility from QCD

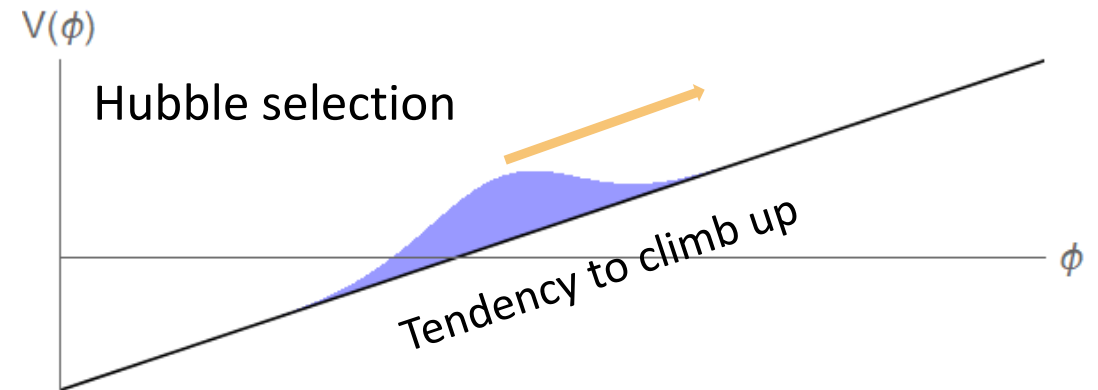
- $\Lambda_\phi \ll M$ = “The tail wagging the dog” is not allowed by naturalness
- Several different translation of original naturalness problems.

At least some of them are quantum stable fine tunings, but I’d omit here.

- We saw some possibility in QCD, but naïve model building was doomed.
- Then, should Hubble selection be discarded? **NO!**

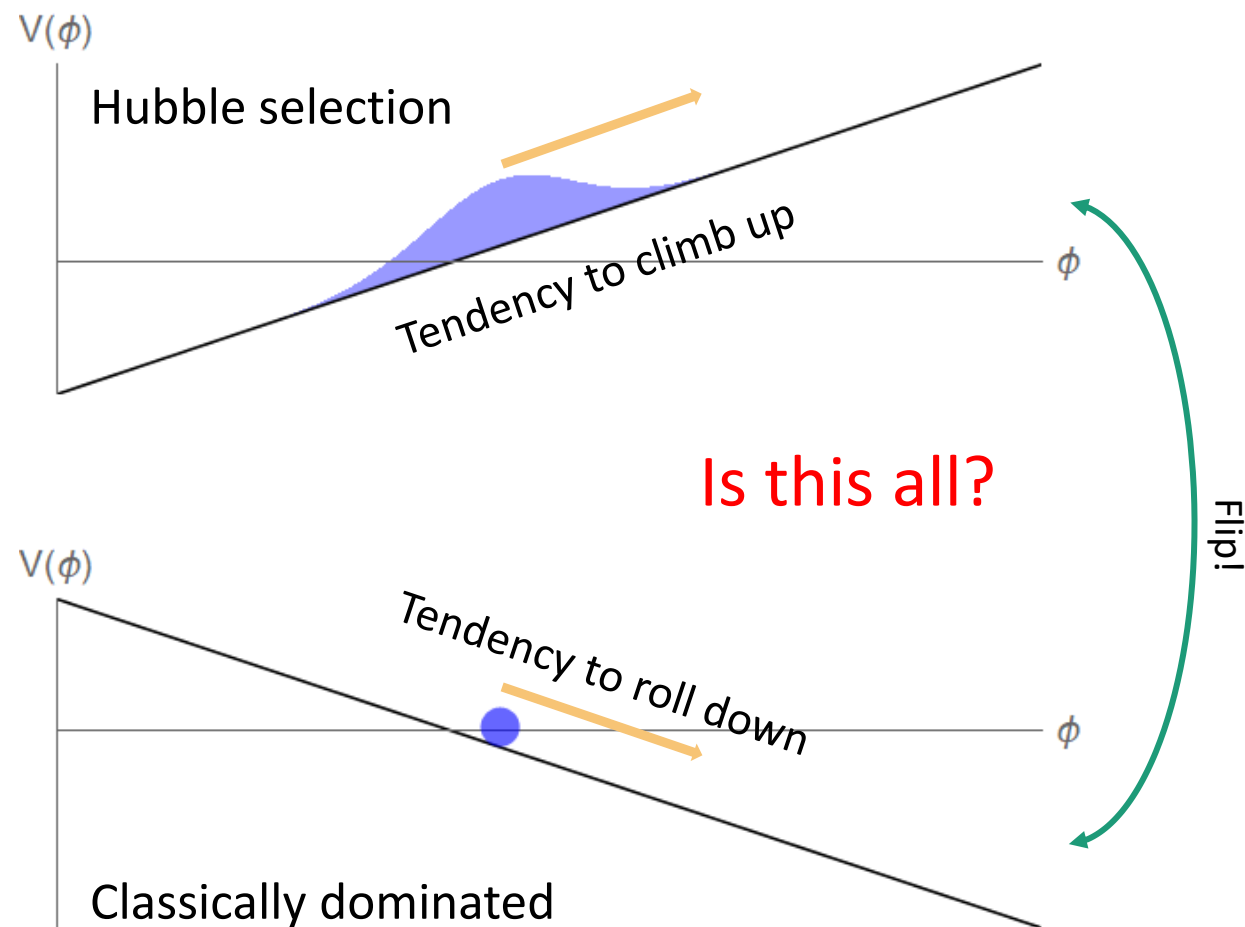
Discussion

- One step back... let us see the big picture
- “Climbing up”



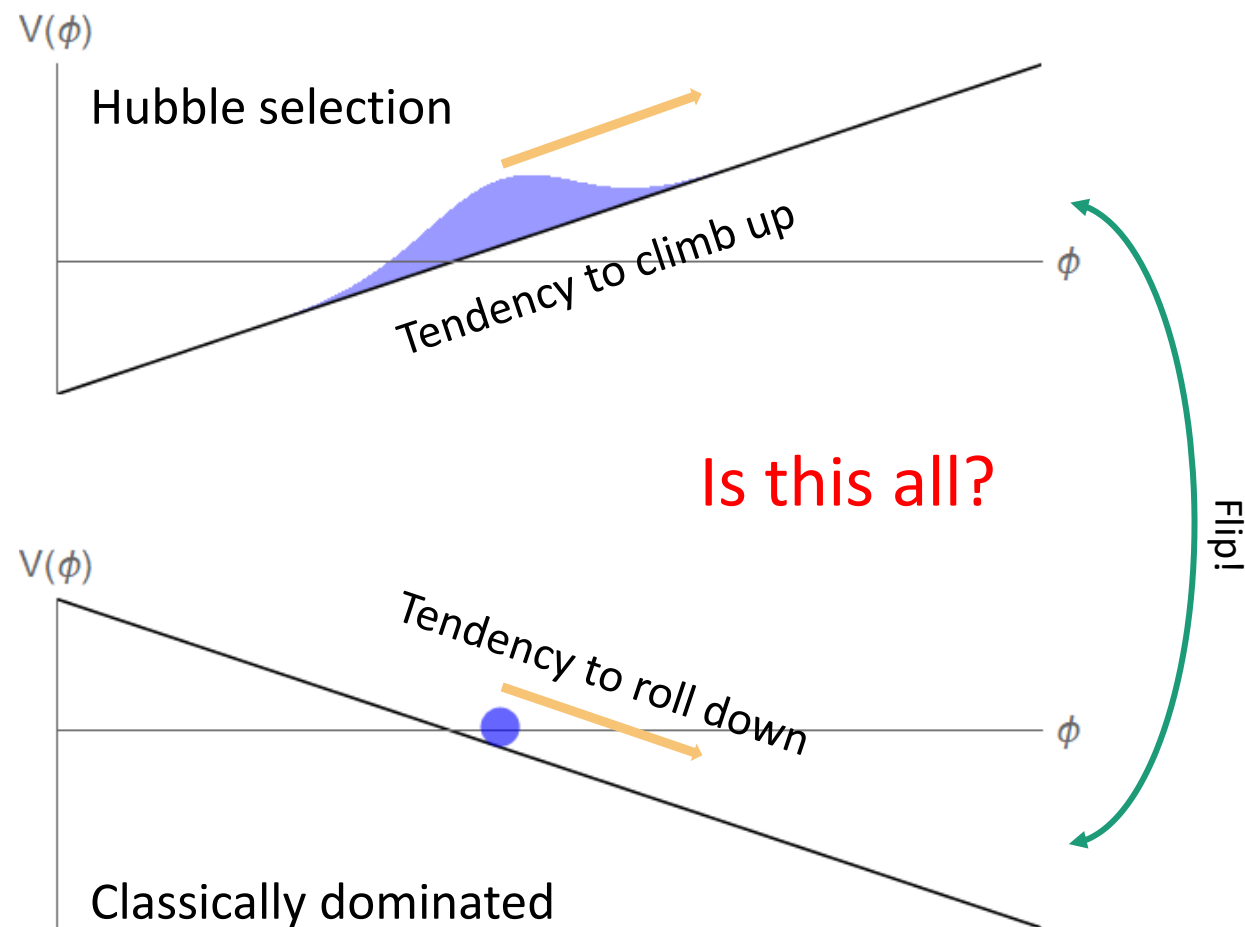
Discussion

- One step back... let us see the big picture
- “Climbing up”
 - Just different tendency, is this all?



Discussion

- One step back... let us see the big picture
- “Climbing up”
 - Just different tendency, is this all?
 - Why should we keep an eye on Hubble selection?



Discussion

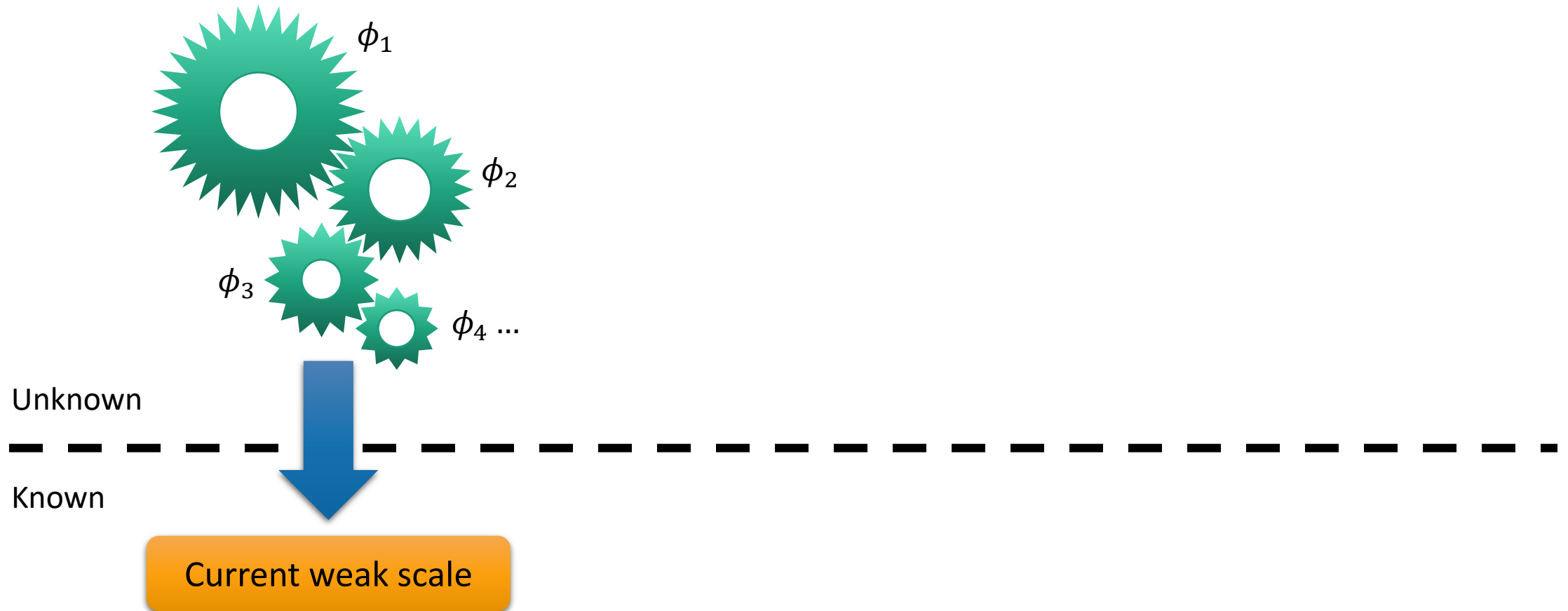
- Answer:

“Because we now have two **competing tendencies**”

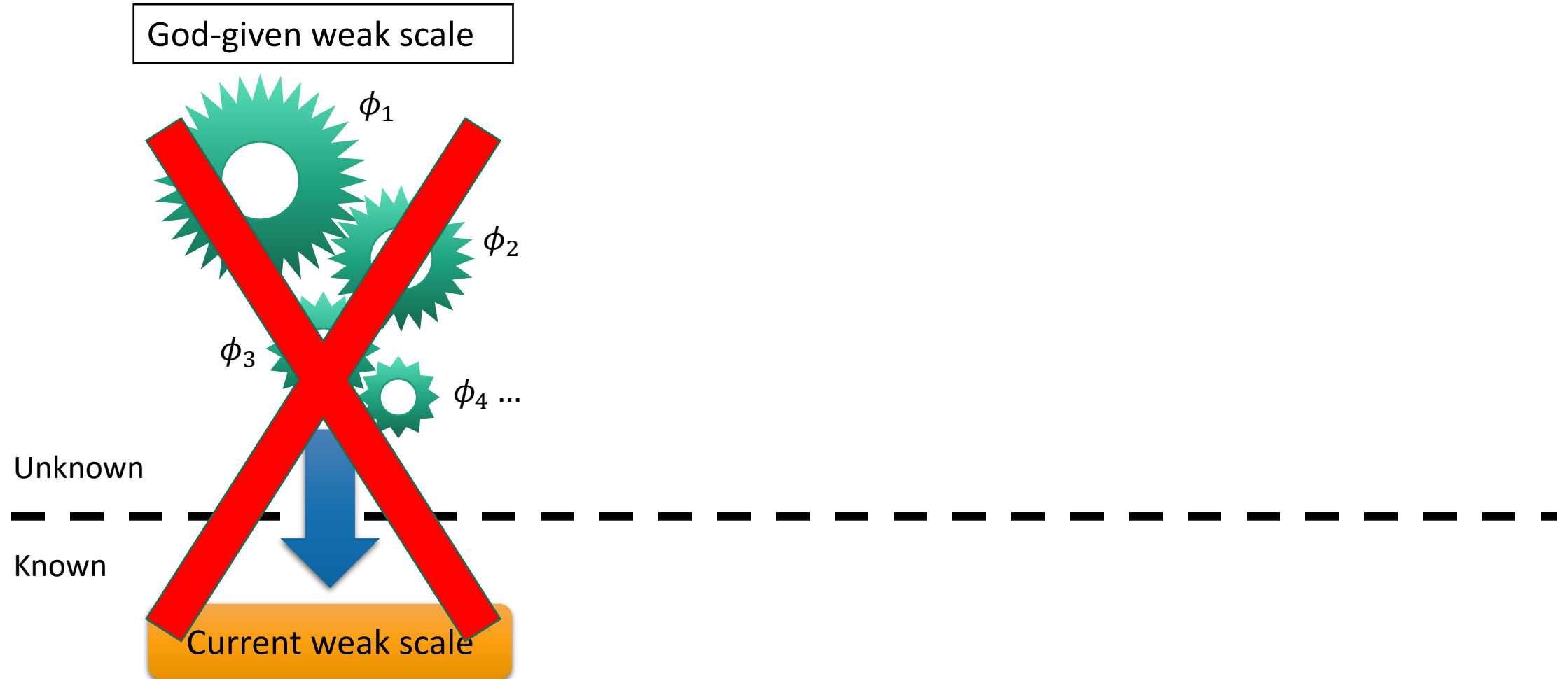
- Let us ask first...

“What kinds of model we are looking for?”

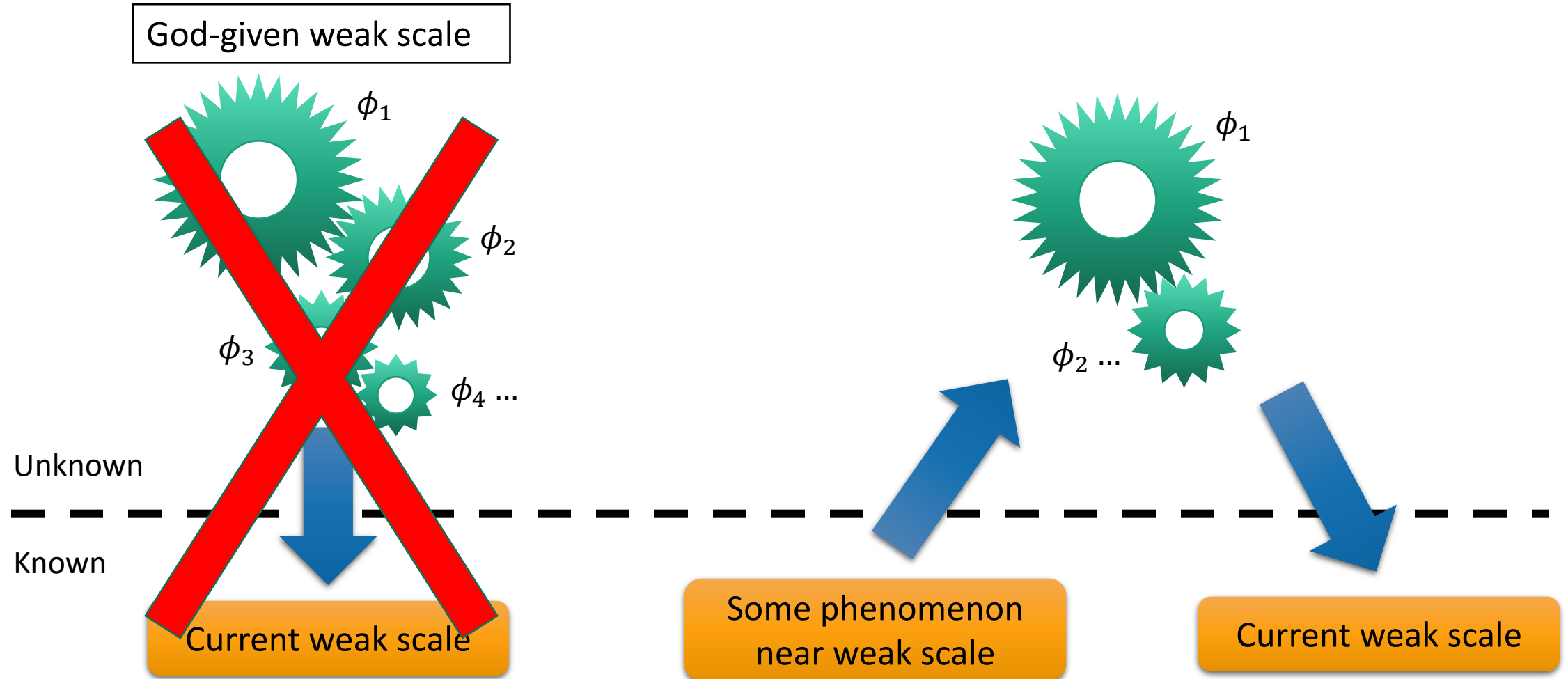
Discussion



Discussion



Discussion



Discussion

- Known sector \rightarrow sub Planckian field range
 - No Hubble selection; always rolls down; minimizes potential energy
 - We want this to act as some trigger or brake of unknown sector to select the current weak scale
- But a brake can work only for **competing tendency**!
 - Rolling down cannot act as a brake for rolling down...
- Thus, Hubble selection is still noteworthy for mechanism in unknown sector

Concluding remarks

- Hubble selection: global field value distribution climbs up the potential
 - Higher potential, higher Hubble rate
 - Super Planckian field range & eternal inflation are required
- Possibility from QCD: weak scale might be selected from QCD phase transition
- Naïve model building was not successful, but Hubble selection is still an attractive mechanism.
 - competing tendency against classical rolling
 - Please keep an eye on it!