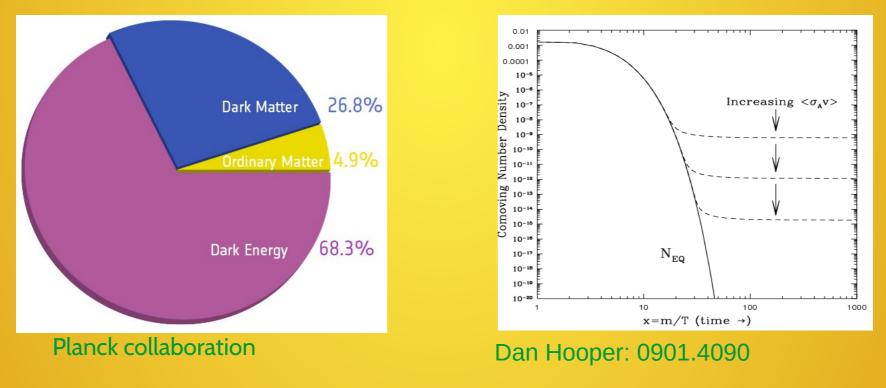
Gravitational wave signals of DM freeze-out

#### Po-Yan Tseng (NTHU) Danny Marfatia (U. of Hawaii)

Reference: JHEP02(2021)022 (arXiv:2006.07313)

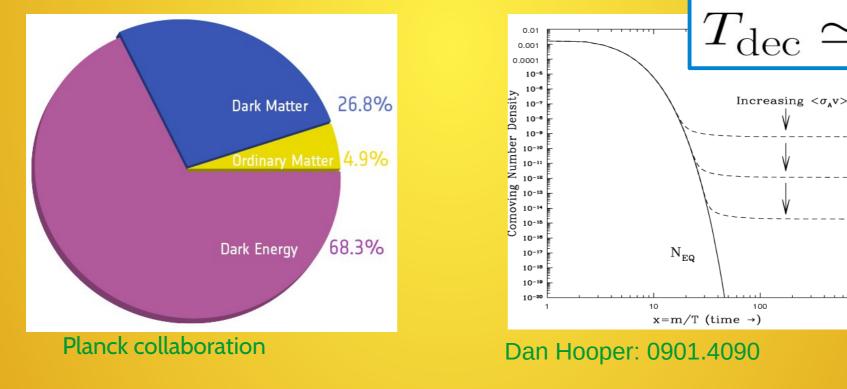
Asia-Pacific Workshop on Particle Physics and Cosmology 2021, Aug. 2<sup>nd</sup> - Aug. 6<sup>th</sup>

- DM production: Thermal freeze-out mechanism.
- Weakly interacting massive DM (WIMP), gives the correct relic density.



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- Weakly interacting massive DM (WIMP), gives the correct relic density.

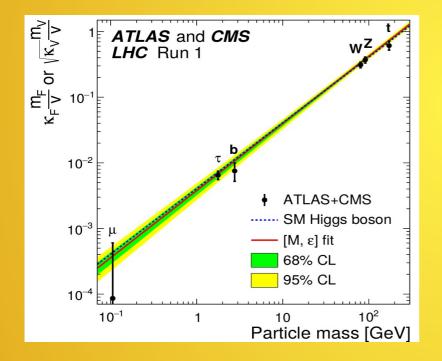


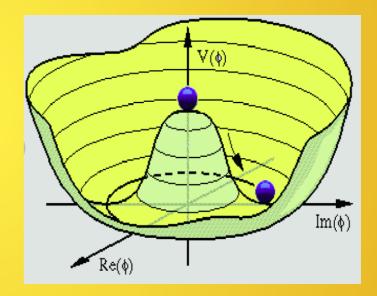
1000

100

 $m_{\hat{}}$ 

 125 GeV Higgs gives the mass to the SM particles through spontaneous symmetry breaking.





Dezso Horvath: Higgs and BSM studies at the LHC

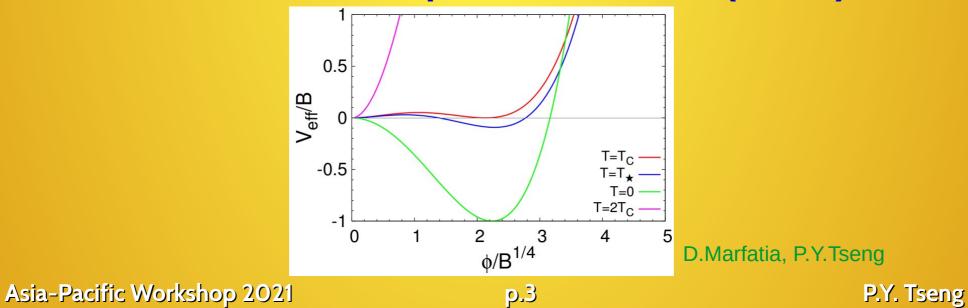


 The origin of DM mass may come from the spontaneous symmetry breaking inducing by another scalar.

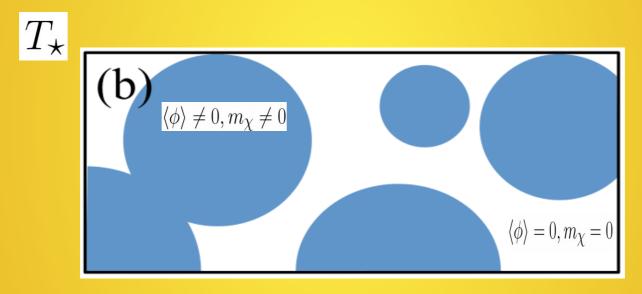
$$\mathcal{L} \supset \bar{\chi} i \partial \!\!\!/ \chi - g_{\chi} \phi \bar{\chi} \chi - V_{\text{eff}}(\phi, T)$$

$$m_{\chi} \simeq g_{\chi} \langle \phi \rangle$$

We consider 1<sup>st</sup> order phase transition (FOPT).



During 1<sup>st</sup> order phase transition (FOPT).



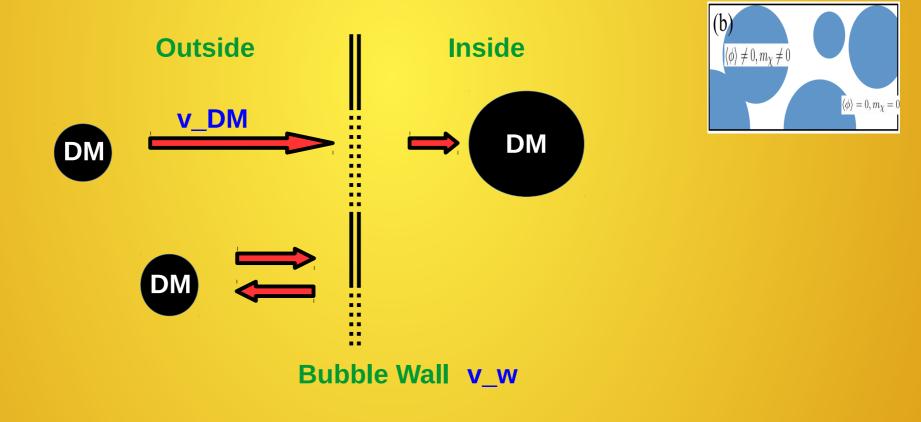
J.P.Hong, S.Jung, K.P.Xie: 2008.04430



## Outline

- Introduction
- Bubble filtering
- Gravitational wave production
- Models
- Summary

 During FOPT, massless (massive) DM particles locate outside (inside) the bubble, and momentum conservation much be satisfied at the bubble wall.





If a thermal DM flux is incident on the wall, the number density of DM that enter the bubble is:

$$n_{\chi}^{\rm in} = n_{\bar{\chi}}^{\rm in} \simeq \frac{g_{\rm DM} T_{\star}^3}{\gamma_w v_w} \left( \frac{\gamma_w (1 - v_w) m_{\chi} / T_{\star} + 1}{4\pi^2 \gamma_{\omega}^3 (1 - v_w)^2} \right) e^{-\frac{\gamma_w (1 - v_w) m_{\chi}}{T_{\star}}}$$

D.Chway, T.H.Jung, C.S.Shin: 1912.04238

 DMs are filtered by the non-relativistic and relativistic bubble wall velocity:

$$n_{\chi}^{\rm in} = \begin{cases} \sim e^{-m_{\chi}/T_{\star}} & \text{for } v_w \to 0\\ \sim e^{-m_{\chi}/(2\gamma_w T_{\star})} & \text{for } m_{\chi}/(\gamma_w T_{\star}) \to 0 \end{cases}$$

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- If  $T_{\star} < T_{dec}$ , the DM inside the bubble is decoupled from the thermal bath and become DM relic abundance.
- DM relic abundance today can be calculated by dividing  $n_{\chi}^{in} + n_{\bar{\chi}}^{in}$  by entropy  $s = (2\pi^2/45)g_{\star S}T^3$ :

$$\Omega_{\rm DM} h^2 \simeq 6.29 \times 10^8 \, \frac{m_\chi (n_\chi^{\rm in} + n_{\bar{\chi}}^{\rm in})}{\rm GeV} \frac{1}{g_{\star S} T_\star^3}$$

$$\Omega_{\rm DM} h^2 \simeq \begin{cases} 1.27 \times 10^8 \left(\frac{m_{\chi}}{\rm GeV}\right) \left(\frac{g_{\rm DM}}{g_{\star S}}\right) \left(\frac{m_{\chi}}{2\gamma_w T_{\star}} + 1\right) e^{-\frac{m_{\chi}}{2\gamma_w T_{\star}}}, & \text{for } v_w \to 1\\ 3.19 \times 10^7 \left(\frac{m_{\chi}}{\rm GeV}\right) \left(\frac{g_{\rm DM}}{g_{\star S}}\right) \left(\frac{1}{v_w}\right) \left(\frac{m_{\chi}}{T_{\star}} + 1\right) e^{-\frac{m_{\chi}}{T_{\star}}}, & \text{for } v_w \to 0. \end{cases}$$

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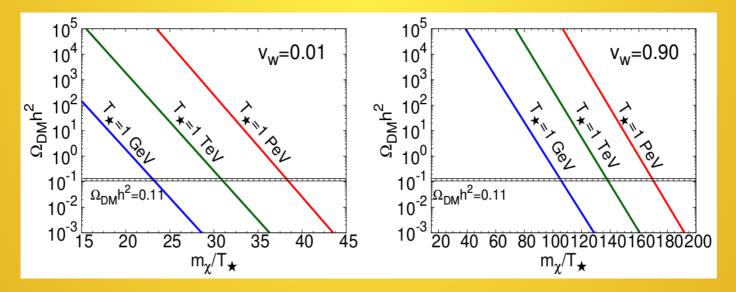


- If  $T_{\star} < T_{dec}$ , the DM inside the bubble is decoupled from the thermal bath and become DM relic abundance.
- DM relic abundance today can be calculated by dividing  $n_{\chi}^{in} + n_{\bar{\chi}}^{in}$  by entropy  $s = (2\pi^2/45)g_{\star S}T^3$ :
- For example:  $m_{\chi} \simeq 1$  TeV,  $v_w \to 1$  requires

$$\frac{m_{\chi}}{2\gamma_w T_{\star}} \simeq 27$$



- If  $T_{\star} < T_{dec}$ , the DM inside the bubble is decoupled from the thermal bath and become DM relic abundance.
- DM relic abundance today can be calculated by dividing  $n_{\chi}^{in} + n_{\bar{\chi}}^{in}$  by entropy  $s = (2\pi^2/45)g_{\star S}T^3$ :



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- Sudden DM freeze-out induced by a FOPT can easily accommodate DM mass above a PeV, which is beyond the current DM direct detection and LHC searches.
- We focus on the Gravitational Wave (GW) signals of Sudden DM freeze-out with a FOPT.



- A FOPT generates GWs from three processes: I).
   Bubble collisions, II). Sound wave in the plasma, III) Magnetohydrodynamic (MHD) turbulence.
- The relevant parameters are required to calculate the GW signals:

$$\begin{cases} T_{\star}, \\ \alpha \equiv \frac{\left(1 - T\frac{\partial}{\partial T}\right) \Delta V_{\text{eff}}|_{T_{\star}}}{\rho(T_{\star})}, \quad \rho \equiv \pi^2 g_{\star} T^4 / 30 \\ \frac{\beta}{H_{\star}} \simeq T_{\star} \frac{d(S_3/T)}{dT} \Big|_{T_{\star}} \\ v_W \end{cases}$$

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A FOPT generates GWs from: I). Bubble collisions

$$h^{2}\Omega_{\rm env}(f) = 1.67 \times 10^{-5} \left(\frac{H_{*}}{\beta}\right)^{2} \left(\frac{\kappa\alpha}{1+\alpha}\right)^{2} \left(\frac{100}{g_{*}}\right)^{\frac{1}{3}} \left(\frac{0.11 \, v_{w}^{3}}{0.42 + v_{w}^{2}}\right) \, S_{\rm env}(f)$$

C.Caprini et. al: 1512.06239

$$S_{\rm env}(f) = \frac{3.8 \ (f/f_{\rm env})^{2.8}}{1 + 2.8 \ (f/f_{\rm env})^{3.8}}$$

• The peak frequency is determined by the time scale of **FOPT**  $1/\beta$ :

$$\frac{f_*}{\beta} = \left(\frac{0.62}{1.8 - 0.1v_w + v_w^2}\right)$$



A FOPT generates GWs from: I). Bubble collisions

$$h^{2}\Omega_{\rm env}(f) = 1.67 \times 10^{-5} \left(\frac{H_{*}}{\beta}\right)^{2} \left(\frac{\kappa\alpha}{1+\alpha}\right)^{2} \left(\frac{100}{g_{*}}\right)^{\frac{1}{3}} \left(\frac{0.11 \, v_{w}^{3}}{0.42 + v_{w}^{2}}\right) \, S_{\rm env}(f)$$

C.Caprini et. al: 1512.06239

$$S_{\rm env}(f) = \frac{3.8 \ (f/f_{\rm env})^{2.8}}{1 + 2.8 \ (f/f_{\rm env})^{3.8}}$$

 The peak frequency is determined by the time scale of FOPT. Then red-shift to present epoch

$$f_{\rm env} = 16.5 \times 10^{-3} \,\mathrm{mHz} \,\left(\frac{f_*}{\beta}\right) \,\left(\frac{\beta}{H_*}\right) \left(\frac{T_*}{100 \,\mathrm{GeV}}\right) \left(\frac{g_*}{100}\right)^{\frac{1}{6}}$$



## Model

## Model: Scalar quartic Model

The finite-temperature quartic effective scalar potential is:

$$V_{\rm eff}(\eta, T) = \frac{\mu^2 + DT^2}{2}\eta^2 - \xi T\eta^3 + \frac{\lambda}{4}\eta^4$$

F.C.Adams: hep-ph/9302321

 Including one-loop Coleman-Weinberg and finitetemperature contributions, potentials of this form are commonly found in *inert singlet, inert doublet, MSSM,* and Majoron models.



#### Models: Scalar quartic Model

The finite-temperature quartic effective scalar potential is:

$$V_{\rm eff}(\eta, T) = \frac{\mu^2 + DT^2}{2}\eta^2 - \xi T\eta^3 + \frac{\lambda}{4}\eta^4$$

#### Benchmark points:

<b>Table 1</b> . Benchmark points (with $\lambda = 0.1$ ) for the Scalar Quartic Model that give $\Omega_{\rm DM}h^2 = 0.11$ .				
	P1	P2	$\mathbf{P3}$	P4
ξ	0.943	0.863	0.796	0.901
D	19.7	16.5	14.0	18.0
$g_\chi$	2.97	3.22	3.48	3.31
α	0.089	0.082	0.076	0.121
$eta/H_{\star}$	1116	1062	1015	1085
$v_\eta/T_\star$	25.71	23.41	21.49	24.51
$v_w$	0.768	0.763	0.760	0.791
$T_{\star}/{ m GeV}$	21.5	23.8	26.1	22.7
$m_{\chi}/{ m GeV}$	1642	1799	1953	1838

D.Marfatia, P.Y. Tseng: 2006.07313

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### Models: Scalar quartic Model

The finite-temperature quartic effective scalar potential is:

$$V_{\rm eff}(\eta, T) = \frac{\mu^2 + DT^2}{2}\eta^2 - \xi T\eta^3 + \frac{\lambda}{4}\eta^4$$

GW signals: 10<sup>-6</sup> LIGO 02 10<sup>-8</sup> 05 41 10<sup>-10</sup> Ω<sub>GW</sub>h² 10<sup>-12</sup> 10<sup>-14</sup> BBO 10<sup>-16</sup> 10<sup>-18</sup>  $0.3 \le v_w \le 1.0$ 10<sup>-20</sup> 10<sup>-4</sup> 10<sup>-2</sup> 10<sup>0</sup> 10<sup>-6</sup>  $10^{2}$  $10^{4}$ f [Hz] D.Marfatia, P.Y. Tseng: 2006.07313

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P.Y. Tseng

## Summary

- We studied the sudden freeze-out DM as an alternative to the continuous thermal freeze-out.
- A necessary ingredient is a FOPT generates DM mass.
- The DM relic abundance may be determined by bubble filtering.
- Because FOPT triggers sudden DM freeze-out, GW offers a signature.



## Thank you!

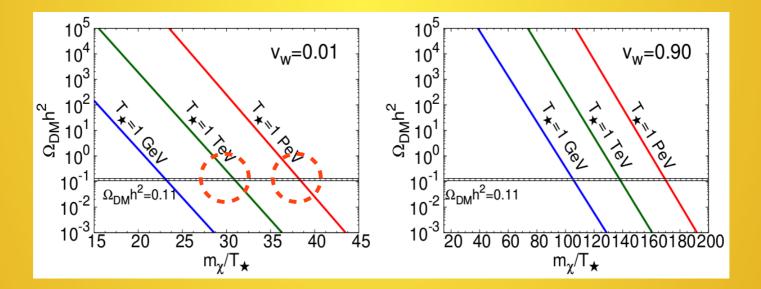
# Back up

- Only massless **DM** particles carry kinetic energy larger than  $m_{\chi}$  can penetrate the bubble walls and become massive.
- DM inside the bubbles abruptly decouples from the thermal bath if  $T_{\star} < T_{\rm dec}$  .
- The bubbles *filter out* certain amount of DM and determine the DM relic abundance.



• The  $m_{\chi}/T_{\star}$  needed to produce the DM relic abundance depends on the velocity of bubble wall  $v_w$ .

$$T_{\star} = m_{\chi}/30$$
 for  $m_{\chi} = 1$  TeV,  $v_w = 0.01$ 

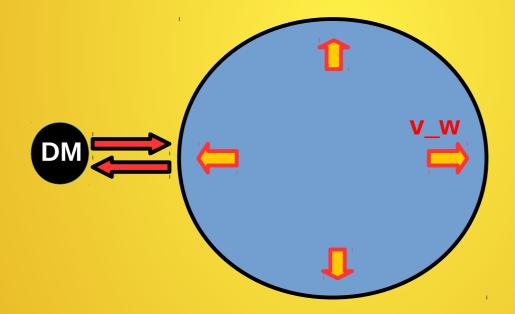


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## **Bubble wall velocity**

 Particles reflected by the bubble wall exert pressure on it, and slow down the bubble wall velocity.





## **Bubble wall velocity**

In the ultrarelativistic limit, the pressure on bubble wall can be obtain from the light degree of freedom inside and outside the bubble:

$$P = \frac{d_n g_\star \pi^2}{90} (1 + v_w)^3 \gamma_\omega^2 T_\star^4$$

D.Chway et.al : 1912.04238 J.R.Espinosa et.al: 1004.4187 D.Bodeker et.al : 0903.4099

$$d_n \equiv \frac{1}{g_{\star}} \left[ \sum_{0.2M_i > \gamma_w T_{\star}} \left( g_i^b + \frac{7}{8} g_i^f \right) \right]$$

• The  $v_w$  can be obtained by solving the eq.  $P = \Delta V_{\text{eff}}$ :

$$\alpha = \frac{d_n}{3}(1+v_w)^3\gamma_\omega^2$$

$$\alpha \equiv \frac{\left(1 - T\frac{\partial}{\partial T}\right) \Delta V_{\text{eff}}|_{T_{\star}}}{\rho(T_{\star})}, \quad \rho \equiv \pi^2 g_{\star} T^4 / 30$$

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## **Bubble wall velocity**

• For bubble wall velocity  $v_w$  faster than the sound speed in plasma, but not ultrarelativistic, we use the approximation:

P.J.Steinhardt, Phys. Rev. D. 25, 2074 (1982)

$$v_w = \frac{\frac{1}{\sqrt{3}} + \sqrt{\alpha^2 + \frac{2}{3}\alpha}}{1 + \alpha}$$



- A FOPT generates GWs from three processes: I).
   Bubble collisions, II). Sound wave in the plasma, III) Magnetohydrodynamic (MHD) turbulence.
- The Euclidean action:

$$S_3(T) = 4\pi \int_0^\infty r^2 dr \left[\frac{1}{2} \left(\frac{d\phi}{dr}\right)^2 + V_{\text{eff}}(\phi, T)\right]$$

Bubble nucleation rate per unit volume:

$$\Gamma(T) = T^4 \left(\frac{S_3}{2\pi T}\right)^{3/2} e^{-\frac{S_3}{T}}$$



- A FOPT generates GWs from three processes: I).
   Bubble collisions, II). Sound wave in the plasma, III) Magnetohydrodynamic (MHD) turbulence.
- The fraction of space in the false vacuum:

$$F(t) = \exp\left[-\frac{4\pi}{3}v_w^3 \int_{t_c}^t dt'(t-t')^3 \Gamma(t')\right]$$

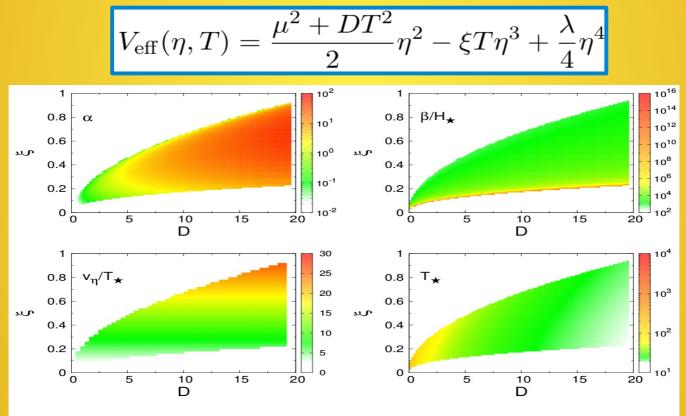
• The percolation temperature  $T_{\star}$  of **FOPT** is determined by :

$$F(t_{\star}) = 1/e \simeq 0.37$$



#### Models: Scalar quartic Model

The finite-temperature quartic effective scalar potential is:





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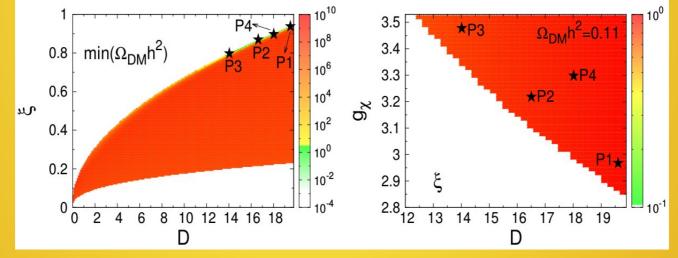
#### Models: Scalar quartic Model

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Correct DM relic:

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## Models: SU(2)x model

 In this dimensionless model, the SM gauge group is extended by as SU(2)x

