

# Electroweak phase transition in the Georgi-Machacek model under constraints of LHC data



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- Intro : why GM model?
- The Georgi-Machacek model
- Constraints
- Simulation
- Electroweak phase transition

# Why GM model?

- Keeping electroweak  $\rho$  parameter at unity
- Enhancement of Higgs signal strength in VV channels
- Generate tiny majorana mass for neutrinos
- More Higgs for new physics

# The GM model

$$\mathcal{L}_{\text{GM}} = \mathcal{L}_{\text{kin}} + \mathcal{L}_Y + \mathcal{L}_\nu - V_H$$

$$\begin{aligned}
 V_H = & m_1^2 \text{tr}(\Phi^\dagger \Phi) + m_2^2 \text{tr}(\Delta^\dagger \Delta) + \lambda_1 \text{tr}(\Phi^\dagger \Phi)^2 + \lambda_2 [\text{tr}(\Delta^\dagger \Delta)]^2 + \lambda_3 \text{tr}[(\Delta^\dagger \Delta)^2] \\
 & + \lambda_4 \text{tr}(\Phi^\dagger \Phi) \text{tr}(\Delta^\dagger \Delta) + \lambda_5 \text{tr} \left( \Phi^\dagger \frac{\tau^a}{2} \Phi \frac{\tau^b}{2} \right) \text{tr}(\Delta^\dagger t^a \Delta t^b) \\
 & + \mu_1 \text{tr} \left( \Phi^\dagger \frac{\tau^a}{2} \Phi \frac{\tau^b}{2} \right) (P^\dagger \Delta P)^{ab} + \mu_2 \text{tr} \left( \Delta^\dagger t^a \Delta t^b \right) (P^\dagger \Delta P)^{ab},
 \end{aligned}$$

$$\Phi = \begin{pmatrix} \phi^{0*} & \phi^+ \\ \phi^- & \phi^0 \end{pmatrix}, \quad \Delta = \begin{pmatrix} \chi^{0*} & \xi^+ & \chi^{++} \\ \chi^- & \xi^0 & \chi^+ \\ \chi^{--} & \xi^- & \chi^0 \end{pmatrix}$$

SU(2)<sub>L</sub>     
 SU(2)<sub>R</sub>

$\text{SU}(2)_L \times \text{SU}(2)_R$

$\downarrow$

Custodial  $\text{SU}(2)_V$

$$\phi^0 = \frac{1}{\sqrt{2}} (h_\phi + i a_\phi),$$

$$\chi^0 = \frac{1}{\sqrt{2}} (h_\chi + i a_\chi),$$

$$\xi^0 = h_\xi,$$

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$$\langle h_\phi \rangle = v_\Phi, \quad \langle h_\chi \rangle = \sqrt{2}v_\Delta, \quad \langle h_\xi \rangle = v_\Delta$$

$$v = \sqrt{v_\phi^2 + 8v_\Delta^2}$$

$SU(2)_L \times SU(2)_R$  $\Delta : 3 \otimes 3$  $\Phi : 2 \otimes 2$ 

↓  
Custodial  $SU(2)_V$

$$H_5 \equiv \begin{bmatrix} H_5^{++} \\ H_5^+ \\ H_5^0 \\ H_5^- \\ H_5^{--} \end{bmatrix}_{m_{H5}}$$

CP-even

 $5 \oplus 3 \oplus 1$ Mixing angle  $\beta$ 

$$H_3 \equiv \begin{bmatrix} H_3^+ \\ H_3^0 \\ H_3^- \end{bmatrix}_{m_{H3}}$$

$$G \equiv \begin{bmatrix} G^+ \\ G^0 \\ G^- \end{bmatrix}$$

CP-odd

 $3 \oplus 1$ Mixing angle  $\alpha$ 

$$H_1 \equiv \begin{bmatrix} H_1^0 \end{bmatrix}_{m_{H1}}$$

$$h \quad (\text{SM-like})_{m_h}$$

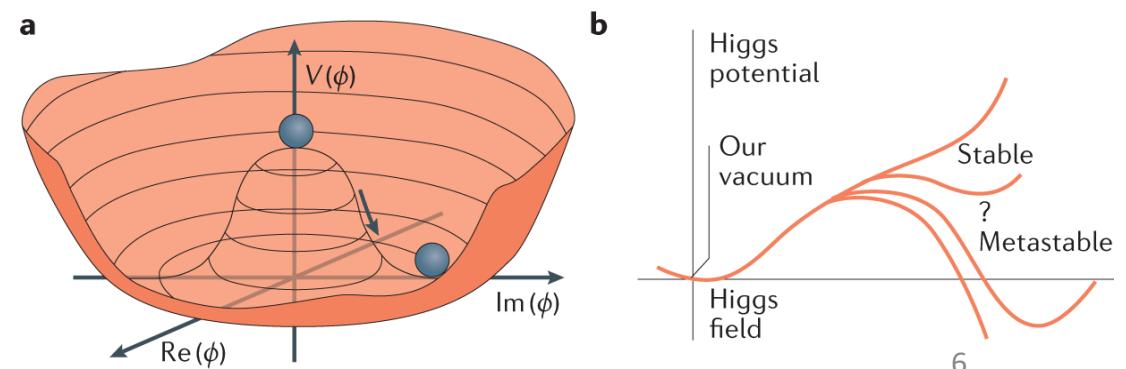
CP-even

# Constraints: theoretical

- Perturbative unitarity conditions
- Bounded from below conditions
- Tadpole conditions:

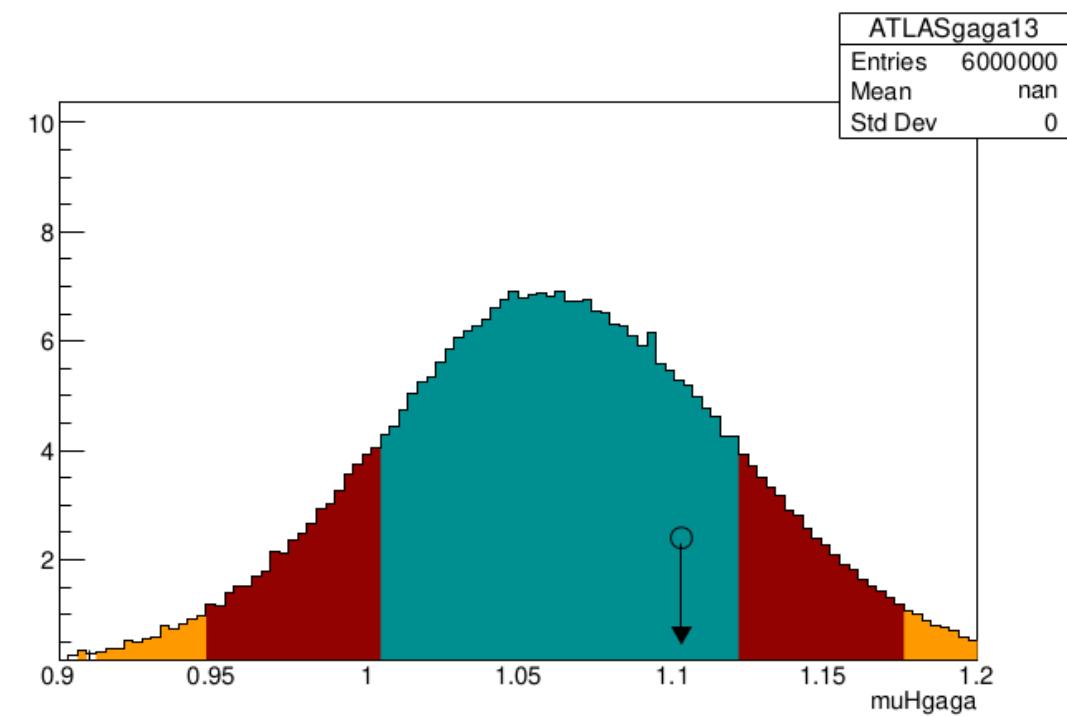
$$\frac{\partial V(\Phi, \Delta)}{\partial h_\phi} = \frac{\partial V(\Phi, \Delta)}{\partial h_\chi} = \frac{\partial V(\Phi, \Delta)}{\partial h_\xi} = 0$$

- Custodial symmetry preservation



# Constraints: Higgs signal strength

- ATLAS & CMS
  - Energy threshold: 7, 8, 13 TeV
  - Process: ggF, VBF, Vh, tth
  - Channels: bb, WW,  $\tau\tau$ , ZZ,  $\gamma\gamma$ , Z $\gamma$ ,  $\mu\mu$
- CDF & D0
  - Energy threshold: 2 TeV
  - Process: Vh, tth
  - Channels: bb



# Constraints: direct searches

## Neutral heavy Higgs bosons

- $\phi_0$ :  $t\bar{t}$ ,  $b\bar{b}$ ,  $\tau\tau$ ,  $\gamma\gamma$ ,  $Z\gamma$ ,  $WW$ ,  
 $VV$ ,  $\phi_0'Z$
- $H_1$ :  $hh$ ,  $hZ$
- $H_3$ :  $hZ$ ,  $H_1Z$

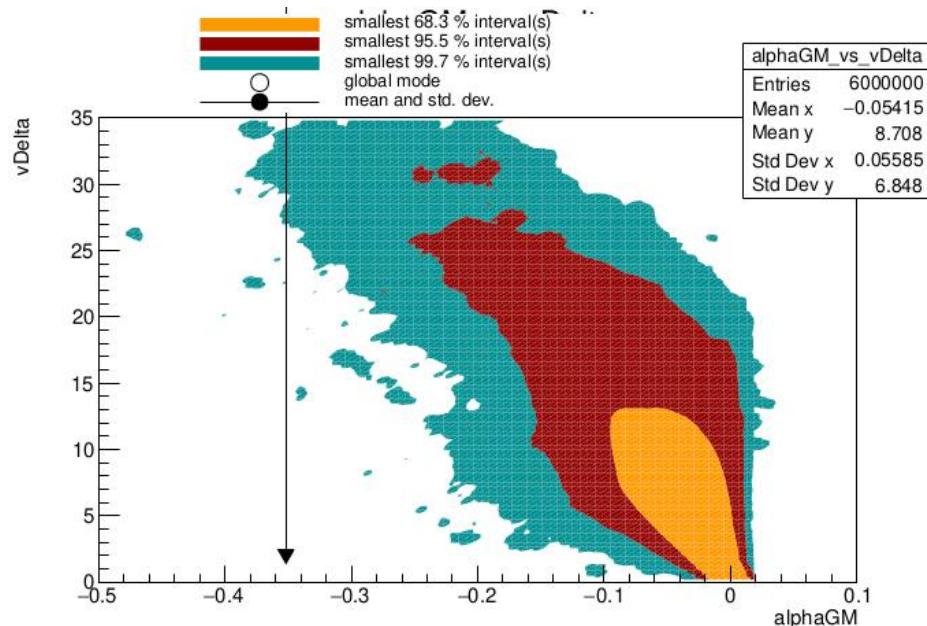
## Charged heavy Higgs bosons

- $H_3^\pm$  :  $\tau^\pm\nu$ ,  $tb$
- $H_5^\pm$  :  $W^\pm Z^\pm$
- $H_5^{\pm\pm} H_5^{\mp\mp}$  :  $(W^\pm W^\pm)(W^\mp W^\mp)$
- $H_5^{\pm\pm}$  :  $(W^\pm W^\pm)$

# HEPfit

Bayesian MCMC:

$$P(\vec{x}|D) = \frac{P(D|\vec{x})P_0(\vec{x})}{\int P(D|\vec{x})P_0(\vec{x})d\vec{x}}$$



Parameters	Scale	Shape	Error/Range
<b>Input</b>	Priors		
$m_2^2 / \text{GeV}^2$	log	Gaussian	$(10^{-6}, 10^6)$
$\lambda_2$	linear	Uniform	$(-3, 3)$
$\lambda_3$	linear	Uniform	$(-3, 3)$
$\lambda_4$	linear	Uniform	$(-3, 3)$
$\lambda_5$	linear	Uniform	$(-12, 12)$
$\mu_1 / \text{GeV}$	linear	Gaussian	$(-10^3, 10^3)$
$\mu_2 / \text{GeV}$	linear	Gaussian	$(-10^3, 10^3)$
<b>Output</b>	Posteriors		
$m_{H_{1,3,5}}/\text{GeV}$	reals	Flat	$(50, 200)$

# Electroweak phase transition

high-temperature expansion:

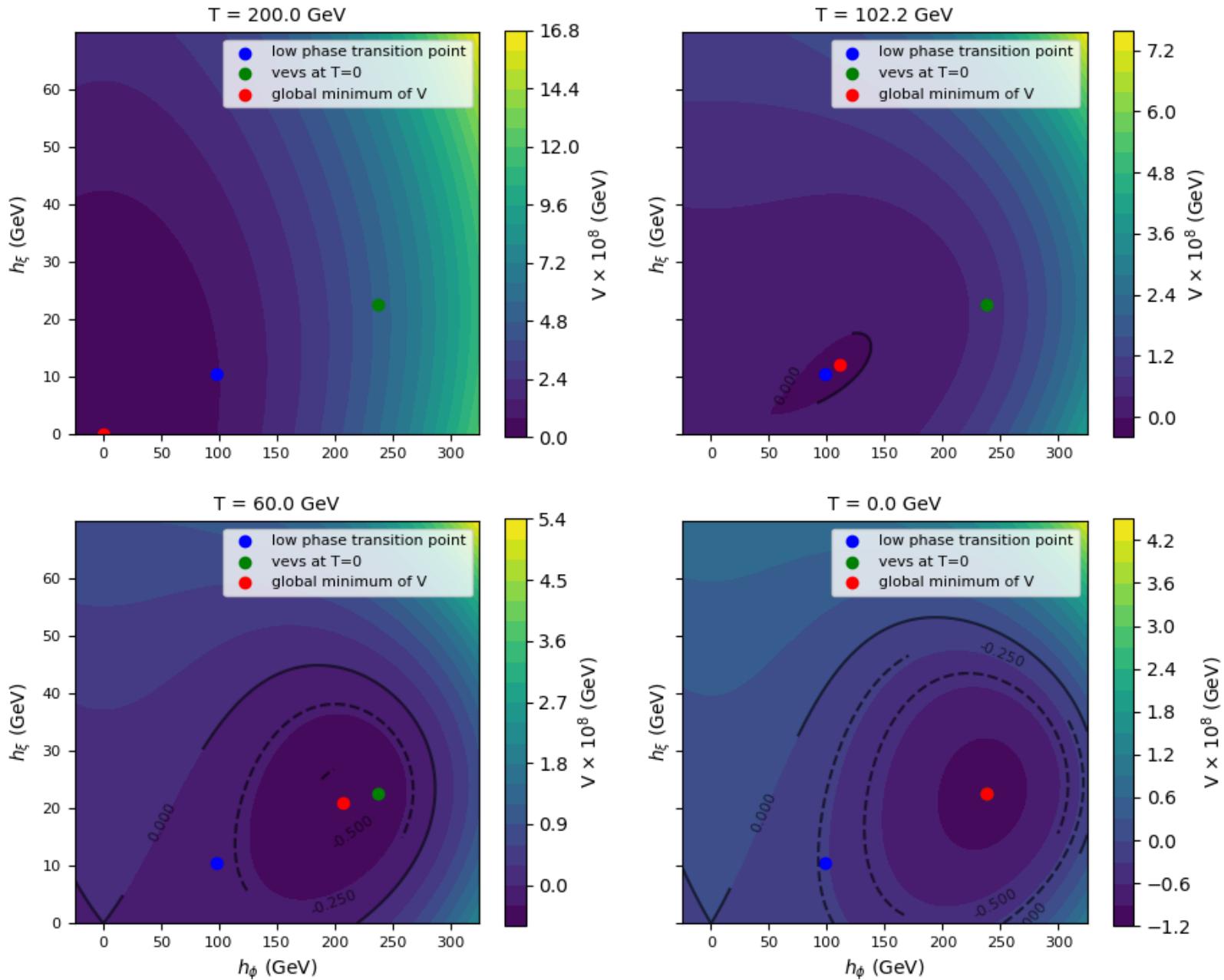
$$V_T = V_0 + \frac{c_\phi T^2}{2} h_\phi^2 + \frac{c_\xi T^2}{2} h_\xi^2 + \frac{c_\chi T^2}{2} h_\chi^2$$

$$c_\phi = \frac{3g^2}{16} + \frac{g'^2}{16} + \lambda_1 + \frac{3\lambda_4}{4} + \frac{1}{4}y_t^2 \csc^2 \beta$$

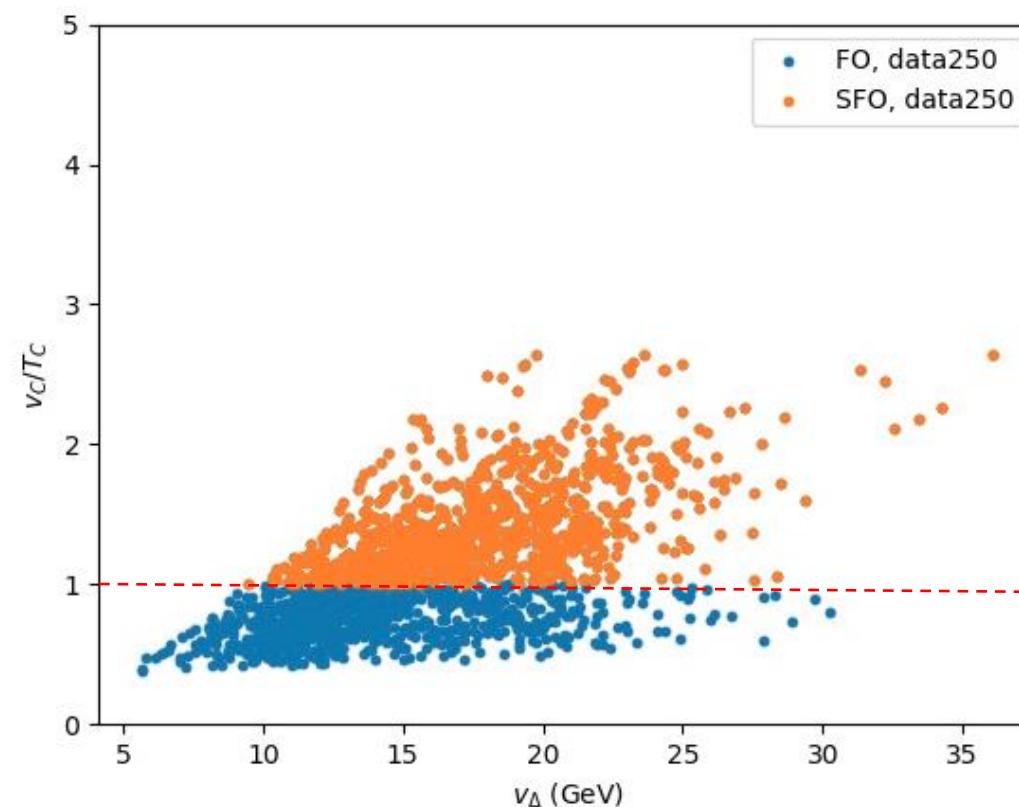
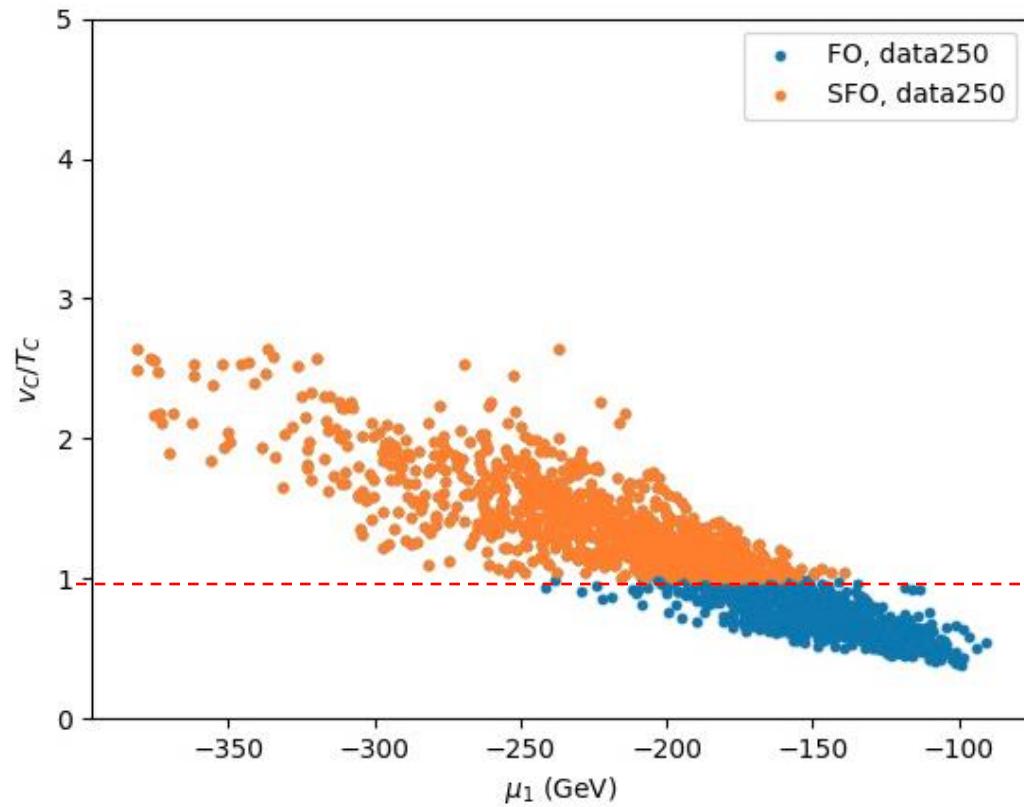
$$c_\xi = \frac{g^2}{2} + \frac{11\lambda_2}{6} + \frac{7\lambda_3}{6} + \frac{\lambda_4}{3}$$

$$c_\chi = \frac{g^2}{2} + \frac{g'^2}{4} + \frac{11\lambda_2}{6} + \frac{7\lambda_3}{6} + \frac{\lambda_4}{3}$$

$\lambda_1 = 0.03521$   
 $\lambda_2 = -0.0106$   
 $\lambda_3 = 0.122212$   
 $\lambda_4 = 0.176963$   
 $\lambda_5 = 0.937933$   
 $\mu_1 = -137.967 \text{ (GeV)}$   
 $\mu_2 = -22.5348 \text{ (GeV)}$   
 $\alpha = -0.06368$   
 $v\Delta = 22.6364 \text{ (GeV)}$   
 $mH1 = 289.333 \text{ (GeV)}$   
 $mH3 = 252.872 \text{ (GeV)}$   
 $mH5 = 115.005 \text{ (GeV)}$   
 $T_c = 102.2332 \text{ (GeV)}$



# Finding the strong first-order phase transition



# Conclusion

- Within LHC data, the strong first-order phase transition is possible.
- No two-step first-order phase transition was found.

# Outlook

- Explore more valid phase space for strong first-order PT.
- Research on determinant/phenomenon for SFOPT.
- Investigate the gravitational wave produced by SFOPT.