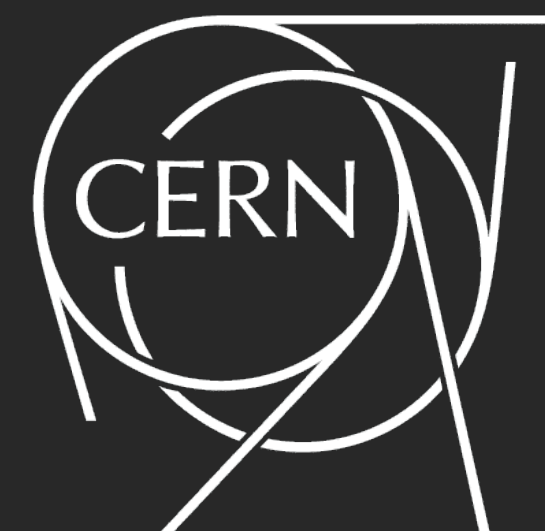


The Supersymmetric Gegenbauer's Twin

Adriana Menkara, Matthew McCullough and Ennio Salvioni

Work in progress

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UTOP MARINA RESORT



**CHUNG-ANG
UNIVERSITY**

Gegenbauer Goldstones



In **QCD** pions are light due to an approx $SU(2)_L \times SU(2)_R$

Quark masses are the only sym breaking parameters

Gegenbauer Goldstones



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Can the higgs also be a pNGb?

Naturalness

Example: Breaking of a U(1) global symmetry

$$\mathcal{L} = \left| \partial_\mu \phi \right|^2 - \frac{\lambda}{4} \left(\left| \phi \right|^2 - \frac{f^2}{2} \right)^2 \quad \text{Associated one pNGb}$$

An explicit sym breaking in the UV

$$V = \epsilon \frac{\lambda}{f^{q-4}} \phi^q + \text{h.c.}$$

Generates the potential

$$V \propto \epsilon m_\rho^2 f^2 \cos \frac{q\Pi}{f} \quad \text{Which is **stable** as a result of symmetry}$$

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UV

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IR

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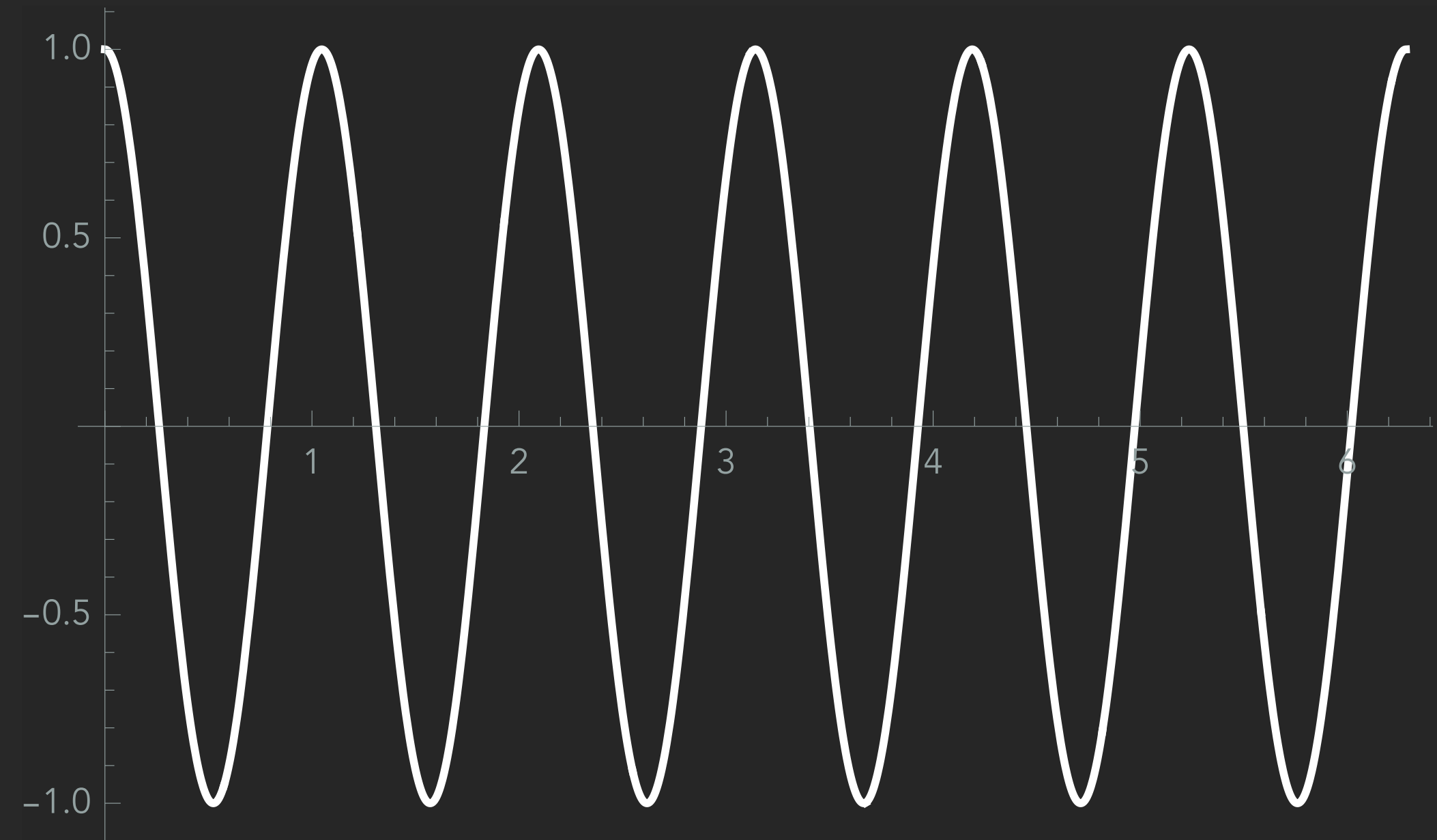
Abelian case

$$V \propto \epsilon m_\rho^2 f^2 \cos \frac{q\Pi}{f}$$

$$\langle \Pi \rangle = f\pi n/q$$

$$m^2 = \epsilon q^2 m_\rho^2$$

$$\lambda = \epsilon q^4 m_\rho^2 / f^2$$



Natural small vev and mass

Non-abelian case

$$SO(N+1) \longrightarrow SO(N)$$

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{\lambda}{4} \left(\phi \cdot \phi - \frac{f^2}{2} \right)^2$$

N massless pNGbs

$$V = \frac{\lambda}{f^{n-4}} \epsilon_{a_1, a_2, \dots, a_n} \phi^{a_1} \phi^{a_2} \dots \phi^{a_n}$$

Spurion in a symmetric irrep (**Traceless**)

$$V = \epsilon m_\rho^2 f^2 G_n^{(N-1)/2} (\cos \Pi/f)$$

Radiatively stable

Non-abelian case

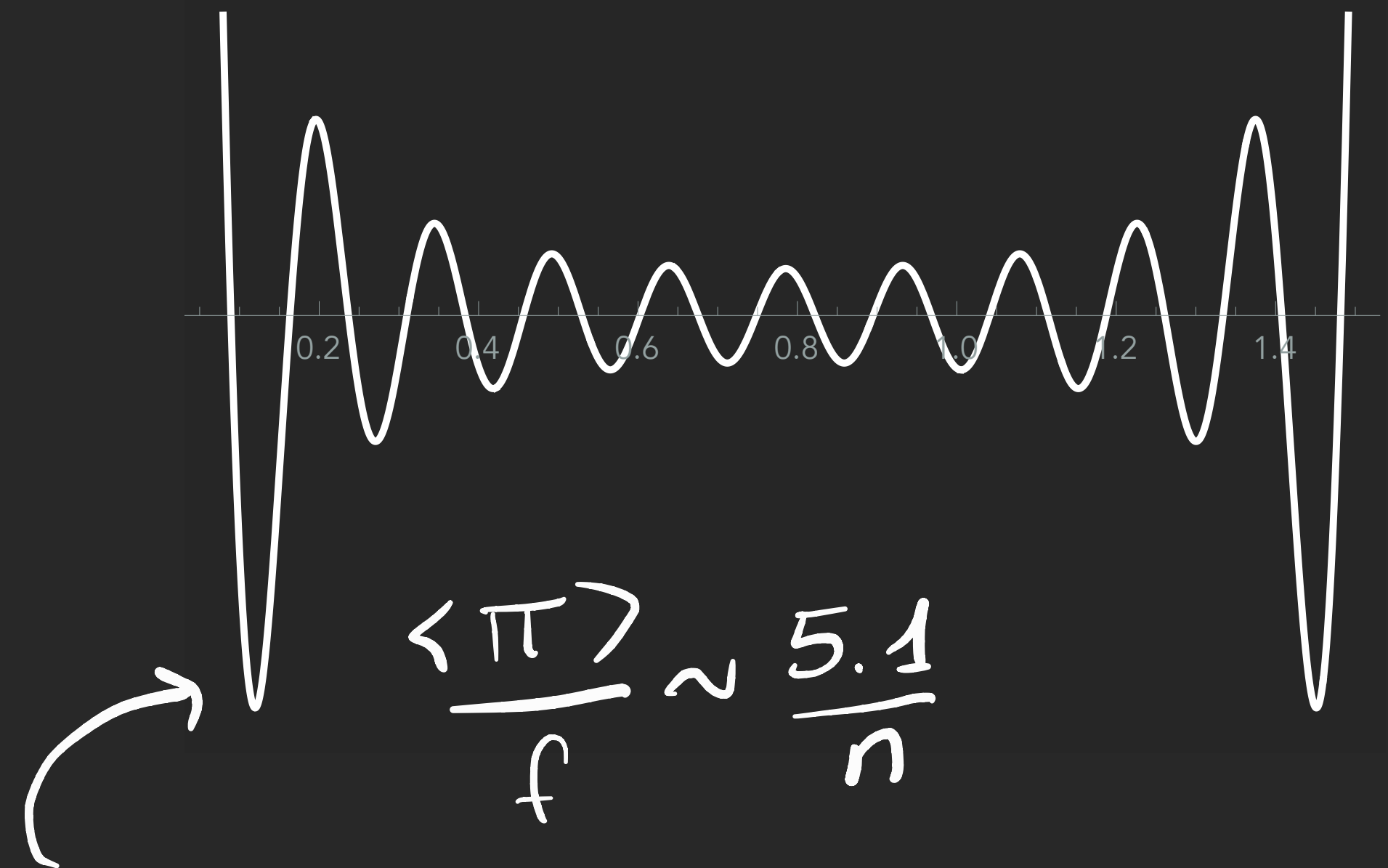
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Gegenbauer polynomial



Natural small vev

Gegenbauer as irreps

$$|t\phi - \tilde{\phi}|^{1-N} = \sum_{n=0}^{\infty} t^n K_n^{i_1 i_2 \dots i_n}(\tilde{\phi}) \phi^{i_1} \phi^{i_2} \dots \phi^{i_n}$$

Taylor series

Similar to angular momentum decomposition

$\tilde{\phi}$ gets a vev

$$\tilde{\phi} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ \vdots \\ f \end{pmatrix}$$

Gegenbauer as irreps

$$|t\phi - \tilde{\phi}|^{1-N} = \sum_{n=0}^{\infty} t^n K_n^{i_1 i_2 \dots i_n}(\tilde{\phi}) \phi^{i_1} \phi^{i_2} \dots \phi^{i_n}$$

$$\phi = \frac{1}{\pi} \sin\left(\frac{\pi}{f}\right) \begin{pmatrix} \pi_1 \\ \pi_2 \\ \vdots \\ \pi_N \\ \cot \frac{\pi}{f} \end{pmatrix}$$

Taylor expansion on t

$$K_n^{i_1 i_2 \dots i_n}(\phi) = \frac{1}{n!} \frac{\partial^n \phi^{1-N}}{\partial \phi^{i_1} \partial \phi^{i_2} \dots \partial \phi^{i_n}}$$

$$|t\phi - \tilde{\phi}|^{1-N} = \sqrt{1 - 2t \cos \Pi/f + t^2}$$

$$= \sum_{n=0}^{\infty} t^n G_n^{(N-1)/2}(\cos \Pi/f)$$

Generating function

$$(1 - 2tx + t^2)^{-\lambda} = \sum_{n=0}^{\infty} t^n G_n^\lambda(x)$$

Gegenbauer as irreps

$$|t\phi - \tilde{\phi}|^{1-N} = \sum_{n=0}^{\infty} t^n K_n^{i_1 i_2 \dots i_n}(\tilde{\phi}) \phi^{i_1} \phi^{i_2} \dots \phi^{i_n}$$

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Tuning

i) Higgs vev

$$V_T \approx \frac{3y_t^2}{16\pi^2} f^2 M_T^2 \left(a_1 \cos \frac{\Pi}{f} + a_2 \cos \frac{2\Pi}{f} \right)$$

$$\langle \Pi \rangle = 0$$

$$\langle \Pi \rangle \approx n\pi f \quad n \in \mathbb{Z}$$

Phenomenologically not viable

$$V_G = \epsilon f^2 M^2 G_n^{3/2} \left(\cos \frac{\Pi}{f} \right) \quad \text{Large } n \text{ allows for natural } v^2 \ll f^2$$

ii) Higgs mass

$$V_T \approx \frac{3y_t^2}{16\pi^2} f^2 M_T^2 \left(a_1 \cos \frac{\Pi}{f} + a_2 \cos \frac{2\Pi}{f} \right)$$

$$m_h^2 \sim \frac{3y_t^2}{16\pi^2} M_T^2$$

$M_T \gtrsim 900 \text{ GeV}$, adds additional fine tuning



Gegenbauer's twin

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Gegenbauer's twin

Recap

Gegenbauer's Goldstones → Natural EW scale

Gegenbauer's Twin → Natural Higgs mass

iii) Couplings

Radiative corrections in the Gegenbauer's Twin

$$V_t \approx \frac{3y_t^4}{64\pi^2} f^4 \left[\sin^4 \frac{h}{f} \log \frac{a}{\sin^2 \frac{h}{f}} + \cos^4 \frac{h}{f} \log \frac{a}{\cos^2 \frac{h}{f}} \right]$$

a is a UV parameter

$$a \approx \frac{g_\star^2}{y_t^2}$$

$$M_\star \approx g_\star f$$

For a pNGb model, $g_\star \sim 4\pi$, thus, the UV is **strongly coupled**

For the Gegenbauer Higgs, $a \lesssim 2$, the UV is perturbative



SUSY Gegenbauer's twin

SUSY Gegenbauer's Twin

$$W = \mu \Phi_u \Phi_d + \lambda S \Phi_u \Phi_d + MS^2$$

$$V = (m_{Hu}^2 + \mu^2) |\phi_u|^2 + (m_{Hd}^2 + \mu^2) |\phi_d|^2 - b (\phi_u \phi_d + h.c.) + \lambda^2 |\phi_u \phi_d|^2$$

Z_2 between MSSM and its mirror \longrightarrow approximate global $U(4)$ symmetry

The Gegenbauer explicitly breaks SUSY

$$V_G^{(n)} = \epsilon f^2 M^2 F(\beta) G_n^{3/2} \left(\cos \left(\frac{2h}{f} \right) \right)$$

$$\tan \beta \equiv \frac{\mu^2 + m_{Hd}^2}{\mu^2 + m_{Hu}^2}$$

SUSY Gegenbauer's Twin

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Sources of symmetry breaking

$$V_G^{(n)} = \epsilon f^2 M^2 F(\beta) G_n^{3/2} \left(\cos \left(\frac{2h}{f} \right) \right) \text{ Explicit}$$

$$V_{g',g} = \frac{g^2 + g'^2}{8} \left[\left(|\phi_u^A|^2 - |\phi_d^A|^2 \right)^2 + \left(|\phi_u^B|^2 - |\phi_d^B|^2 \right)^2 \right] \text{ D-terms}$$

$$V_t = \frac{3y_t^4}{16\pi^2} \left(|\phi_u^A|^4 \log \frac{m_{\tilde{t}}^2}{m_t^{A2}} + |\phi_u^B|^4 \log \frac{m_{\tilde{t}}^2}{m_t^{B2}} \right) \text{ Top sector}$$

SUSY contribution to the pNGb potential

$$V_{U(4)} = \frac{3(y_t f \sin \beta)^4}{64\pi^2} \left[\sin^4 \left(\frac{h}{f} \right) \log \frac{a}{\sin^2 \left(\frac{h}{f} \right)} + \cos^4 \left(\frac{h}{f} \right) \log \frac{a}{\cos^2 \left(\frac{h}{f} \right)} \right]$$

$$a = \frac{2m_{\tilde{t}}^2}{(y_t f \sin \beta)^2} e^{\frac{2\pi^2 \cos^2(2\beta)(g'^2 + g^2)}{3(y_t \cos \beta)^4}}$$

$m_{\tilde{t}} \sim \mathcal{O}(f)$: we have a model that is a hybrid between a **SUSY** model and a **pNGb** model

Summary

- ✓ Gegenbauer polynomials are a class of radiatively stable potentials
- ✓ The twin Gegenbauer points to a perturbative coupled UV
- ✓ We provided a SUSY UV completion, valid up to very small scales
- ✓ Since SUSY is perturbative, IR parameters can be calculated

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On-going work: LHC-HL. Precision measurements vs direct searches

Further questions: Origin of the explicit sym. breaking?

Cosmology (baryogenesis, phase transitions...)